GENERALIZED TRAVELING WAVES IN DISORDERED MEDIA: EXISTENCE, UNIQUENESS, AND STABILITY

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ABSTRACT. We prove existence, uniqueness, and stability of transition fronts (generalized traveling waves) for reaction-diffusion equations in cylindrical domains with general inhomogeneous ignition reactions. We also show uniform convergence of solutions with exponentially decaying initial data to time translates of the front. In the case of stationary ergodic reactions the fronts are proved to propagate with a deterministic positive speed. Our results extend to reaction-advection-diffusion equations with periodic advection and diffusion.

1. Introduction and Results

In this paper we study time-global solutions, called transition fronts or generalized traveling waves, of reaction-diffusion equations on infinite cylinders. We consider the PDE

$$u_t = \Delta u + f(x, u) \tag{1.1}$$

which is used in modeling of processes such as autocatalytic chemical reactions, propagation of advantageous genes in a population, and combustion. The function $u(t,x) \in [0,1]$ is the (normalized) concentration of a reactant or allele, or temperature of a combusting solid or gaseous medium. The non-negative reaction term f accounts for increase of concentration/temperature due to a chemical reaction or burning and satisfies f(x,0) = f(x,1) = 0. We will be particularly interested in *ignition reactions*, which vanish for u smaller than some ignition temperature $\theta(x) > 0$ and are used in the modeling of combustion, but we will also treat general non-negative reactions. The function f will satisfy some uniform bounds but will otherwise be an arbitrary non-periodic function of x.

We will also consider the more general equation

$$u_t + q(x) \cdot \nabla u = \operatorname{div}(A(x)\nabla u) + f(x, u), \tag{1.2}$$

with an incompressible mean-zero vector field q representing advection and a uniformly elliptic diffusion operator $\operatorname{div}(A\nabla)$ representing inhomogeneous diffusion. Unlike f, both q and A will be assumed periodic (with the same period).

Our main goal is the proof of existence and uniqueness of transition fronts under very general conditions on f. Moreover, we also want to prove uniform convergence of arbitrary solutions of (1.1)/(1.2) with exponentially decaying initial data to these fronts, thus describing the behavior of very general solutions of the PDEs. We will do this for (1.1)/(1.2) with ignition reactions, and also prove existence of fronts for some non-ignition reactions, on the cylindrical domain $D \equiv \mathbb{R} \times \mathbb{T}^{d-1}$ (i.e., $\mathbb{R} \times [0,1]^{d-1}$ with periodic boundary conditions). However, all our results can be extended to open connected domains $D \subseteq \mathbb{R}^d$ with a smooth

boundary which are periodic in the first variable and bounded in the others, with either periodic or Neumann boundary conditions on ∂D (the latter being $\nu \cdot \nabla u = 0$ for (1.1) or $\nu \cdot A \nabla u = 0$ and $q \cdot \nu = 0$ for (1.2), with ν the outward unit normal to ∂D). These include cylinders with periodically ondulating boundaries and a periodic array of holes.

Such domains, unbounded in arbitrarily many variables, have been considered in [1] where transition fronts for periodic f (as well as q and A) have been studied. We restrict ourselves here to domains unbounded in only one variable because this is essentially the only case when transition fronts for ignition reactions can be unique even in homogeneous media (one moving right and one moving left). Moreover, we will show elsewhere [23] that there are examples of ignition reactions on $D = \mathbb{R}^2$ where no transition fronts exist! Nevertheless, some questions about solutions of (1.1)/(1.2) on domains unbounded in several variables can also be treated by our methods [23].

The following is the definition of a transition front from [2], adapted to our domain D.

Definition 1.1. A transition front (moving to the right) is a solution $w : \mathbb{R} \times D \to [0,1]$ of (1.1) or (1.2) that is global in time and satisfies for each $t \in \mathbb{R}$,

$$\lim_{x_1 \to -\infty} w(t, x) = 1 \quad \text{and} \quad \lim_{x_1 \to +\infty} w(t, x) = 0$$
 (1.3)

uniformly in $(x_2, \ldots, x_d) \in \mathbb{T}^{d-1}$. In addition, the front must have a bounded width (uniformly in time). That is, if $I_{\varepsilon}(t) \subset \mathbb{R}$ is the smallest interval such that $w(t, x) \in [0, \varepsilon] \cup [1 - \varepsilon, 1]$ for $x \in D \setminus (I_{\varepsilon}(t) \times \mathbb{T}^{d-1})$, then $L_{w,\varepsilon} \equiv \sup_{t \in \mathbb{R}} |I_{\varepsilon}(t)| < \infty$ for each $\varepsilon > 0$. The domain $I_{\varepsilon_0}(t) \times \mathbb{T}^{d-1}$ for some small $\varepsilon_0 > 0$ will be referred to as the reaction zone.

We define a transition front moving to the left as above but with (1.3) replaced by

$$\lim_{x_1 \to -\infty} w(t, x) = 0 \quad \text{and} \quad \lim_{x_1 \to +\infty} w(t, x) = 1.$$

Remark. In cylindrical domains fronts moving both right and left can exist. Since the transformation $x_1 \mapsto -x_1$ interchanges the two directions, we will mostly consider only fronts moving to the right.

The simplest case of transition fronts are traveling fronts whose shape is time-independent. Their study goes back to the works of Kolmogorov, Petrovskii, Piskunov [11] and Fisher [8] in 1937. They considered (1.1) in one spatial dimension $D=\mathbb{R}$ and with x-independent KPP reaction f(u) (such that $0 < f(u) \le f'(0)u$ for $u \in (0,1)$). In this case traveling fronts of the form w(t,x)=W(x-ct) exist precisely when the front speed $c \ge c^*$, with the minimal speed being $c^*=2\sqrt{f'(0)}$. The front profile W is time independent and satisfies the ODE $W_{xx}+cW_x+f(W)=0$ with $W(\infty)=0$ and $W(-\infty)=1$. The situation is the same for general positive reactions (such that f(u)>0 for $u \in (0,1)$) but the formula for $c^*>0$ is more complicated. In contrast, the front and its speed $c^*>0$ are unique for ignition reactions (such that f(u)=0 for $u \in [0,\theta]$ and f(u)>0 for $u \in (\theta,1)$, with $\theta>0$).

The ansatz $w(t, x) = W(x - cte_1)$ also works in more dimensions when q, A, f are independent of x_1 , and the answers are the same as above. In particular, the case of mean-zero shear flows q and x-independent A, f has been treated by Berestycki, Larrouturou, Lions [5], and Berestycki and Nirenberg [6]. On the other hand, for periodic q, A, f (with the same period p), the front profile can only be expected to be time-periodic in a moving frame in

the sense $w(t+p/c,x+pe_1)=w(t,x)$ for some speed c>0. Such pulsating fronts are of the form $w(t,x)=W(x_1-ct,x)$ with the profile W decreasing in the first variable, periodic in the second, and satisfying $\lim_{s\to-\infty}W(s,x)=1$ and $\lim_{s\to+\infty}W(s,x)=0$ uniformly in $x\in D$. Existence and uniqueness of pulsating fronts (with mean-zero q) was proved by Xin [20] for x-independent ignition reactions in \mathbb{R}^d , and by Berestycki and Hamel [1] for x-periodic ignition reactions (such that if $\tilde{\theta}\equiv\inf_x\theta(x)$, with $\theta(x)\equiv\inf\{u>0\,|\,f(x,u)>0\}$ the ignition temperature, then $\tilde{\theta}>0$ and $\sup_{x\in D}f(x,u)>0$ for each $u\in(\tilde{\theta},1)$) in general periodic domains. [1] also treats x-periodic positive reactions (as above but with $\tilde{\theta}=0$) and again proves existence of fronts with precisely the speeds $c\geq c^*$ for some $c^*>0$.

The situation is different for disordered media, when no such ansatz exists and one has to work directly with the original PDE. Constant or periodic front profiles cannot be expected and fronts need not have a well defined speed. The definition of a transition front above has been given by Berestycki and Hamel [2] in a more general setting and on arbitrary domains. An alternative definition has been given by Shen [18], who studied fronts in time-random one-dimensional media and established some sufficient conditions on their existence. This formalizes an earlier definition by Matano which essentially requires the profile of the front to be a continuous function of the medium near the reaction zone.

Due to the above difficulties, existence and uniqueness of transition fronts in general disordered media has so far only been proved for (1.1) in one dimension, for some ignition reactions with x-independent ignition temperatures. Namely, Nolen and Ryzhik [15] and independently Mellet, Roquejoffre, Sire [14] have proved that such fronts exist on $D = \mathbb{R}$ when the reaction satisfies $b_0 f_0(u) \leq f(x, u) \leq b_1 f_0(u)$ and $f'_0(1) < 0$. Here $0 < b_0 \leq b_1 < \infty$ and f_0 is of ignition type with $f_0(u) > 0$ if and only if $u \in (\theta, 1)$, $\theta > 0$. Moreover, Mellet, Nolen, Roquejoffre, Ryzhik [13] proved that if $f(x, u) = b(x) f_0(u)$ with $b(x) \in [b_0, b_1]$, then the (right-moving) front is unique and exponentially stable with respect to exponentially in space decaying perturbations. In addition, Vakulenko and Volpert [19] proved existence of fronts for (1.1) in 1D with f some small perturbation of an x-independent bistable reaction $f_0(u)$ (such that $f'_0(0), f'_0(1) < 0$, $f_0(u) < 0$ for $u \in (0, \theta)$, and $f_0(u) > 0$ for $u \in (\theta, 1)$). We also mention here results for x-independent bistable reactions in some special domains — non-existence of moving fronts (while stationary fronts exist) in quickly widening cylinders by Chapuisat and Grenier [7], and later also Berestycki, Hamel, Matano [3], and existence of fronts in \mathbb{R}^d with compact star-shaped obstacles [3].

The usage of 1-dimensional techniques plays an important role in [13, 14, 15], as does the requirement of an x-independent ignition temperature θ . We present here a new method that can handle more general reactions, works in several dimensions, and in the presence of (periodic) q and A. In particular, we prove existence of a unique transition front when f lies between two arbitrary x-independent ignition reactions with different ignition temperatures. We also prove existence of fronts in the more general case when the upper bound is a positive reaction, only requiring a bound on its derivative at zero (these fronts are not unique in general, as is the case for homogeneous media). We note that the requirement of a bound of this type is necessary to guarantee existence and cannot be improved except possibly by a constant (see Remark 1 after Theorem 1.3).

Let us now state our main results. We will start with the special case (1.1). We will assume the following hypotheses on f.

(H1): The reaction f is uniformly Lipschitz with constant $K \geq 1$ and lies between two x-independent reactions, one of ignition type and the other positive or ignition. More specifically, there are Lipschitz functions f_0, f_1 , decreasing on $[1 - \varepsilon, 1]$ for some $\varepsilon > 0$, such that $f_0(u) \leq f(x, u) \leq f_1(u)$ for $(x, u) \in D \times [0, 1]$. In addition, $f_0(0) = f_0(1) = f_1(0) = f_1(1) = 0$, there is $\theta \in (0, 1)$ such that $f_0(u) = 0$ for $u \in [0, \theta]$ and $f_0(u) > 0$ for $u \in (\theta, 1)$, and there is $\theta' \in [0, 1)$ such that $f_1(u) = 0$ for $u \in [0, \theta']$ and $f_1(u) > 0$ for $u \in (\theta', 1)$.

Assume that $c_0 > 0$ is the speed of the unique (right-moving) traveling front for (1.1) with f replaced by the x-independent reaction f_0 (this front is of the form $w(t,x) = W(x_1 - c_0 t)$). We then obtain existence of a transition front for f provided $f'_1(0) < c_0^2/4$. Note that this condition is automatically satisfied when $f'_1(0) = 0$ (e.g., if f_1 is ignition). For instance, a front exists in the case $f(x,u) = b(x)f_1(u)$ with $f'_1(0) = 0$ and b Lipschitz and uniformly bounded away from 0 and ∞ . Moreover, if f_1 is ignition, then we also prove uniqueness of the (right-moving) front and that this unique front is a global attractor of general exponentially decaying initial data.

Our reaction f can have an x-dependent ignition temperature $\theta(x) \in [\theta', \theta]$ and we do not require $\sup_{x \in D} f(x, u) > 0$ for each $u \in (\inf_{x \in D} \theta(x), 1)$. Nevertheless, we will need to impose a similar but weaker natural hypothesis (which is automatically satisfied when $\theta' = \theta$) that if f(x, u) is large enough (rather than just positive) for some (x, u) and $v \in (u, 1)$, then $f(\cdot, v)$ cannot become uniformly arbitrarily small on some large neighborhood of $x \in D$.

Definition 1.2. Given any $\zeta > 0$, define

$$\alpha_f(x) \equiv \inf (\{u \in (0,1) \mid f(x,u) \ge \zeta u\} \cup \{1\}).$$
 (1.4)

Assume $g \in C([0,1])$ with g(0) = g(1) = 0 and g(u) > 0 for $u \in (0,1)$. If for all $x \in D$ and all $u \in [\alpha_f(x), 1]$ we have $\sup_{|\tilde{x}_1 - x_1| \le ||g||_{\infty}^{-1}} f(\tilde{x}, u) \ge g(u)$, then we say that $f \subseteq D$ and g(u) = 0 and g(u) = 0 are g(u) = 0.

Remarks. 1. If $\zeta > f_1'(0)$, then $\alpha_f(x) = \min(\{u \in (0,1) \mid f(x,u) = \zeta u\} \cup \{1\}) > \theta'$.

2. We could replace $||g||_{\infty}^{-1}$ by some large M, but then the notation would have to include it, thus being too cumbersome. Of course, we will consider $f \zeta$ -majorizing some small g (with a large $||g||_{\infty}^{-1}$), so this simplification can be made without loss.

The figure below gives an example of such an f at a fixed x (satisfying f(x, u) > 0 for $u \in [\alpha_f(x), 1]$, while the definition requires non-vanishing of f only at such u and some u-dependent point \tilde{x} whose distance from x is uniformly bounded).

Here is our main result for (1.1).

Theorem 1.3. Let f satisfy the hypotheses (H1). Assume that f ζ -majorizes g for some $\zeta < c_0^2/4$ and some g as in Definition 1.2.

- (i) If $f'_1(0) < c_0^2/4$, then there exists a transition front w for (1.1) moving to the right with $w_t > 0$ (and another moving to the left).
- (ii) If f_1 is an ignition reaction (i.e., $\theta' > 0$) and f is non-increasing in u on $[\theta'', 1]$ for some $\theta'' < 1$, then there is a unique (up to time shifts) transition front w_+ for (1.1) moving to the right (and another w_- moving to the left).

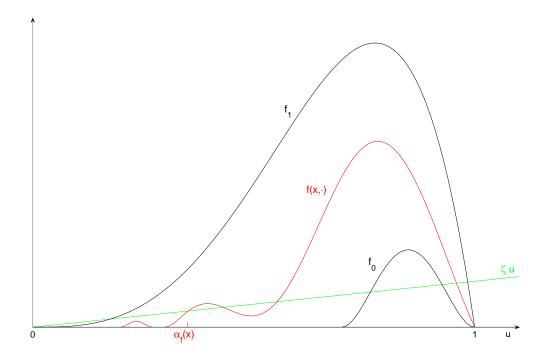


FIGURE 1. Example of a reaction (at a fixed $x \in D$) satisfying our hypotheses.

(iii) In the setting of (ii) we have convergence of solutions with exponentially decaying initial data to time shifts of w_{\pm} in the sense of Definition 1.4 below. Moreover, this convergence is uniform in f, a, u — the s_{ε} in the definition only depends on $f_0, f_1, \zeta, g, K, \theta'', Y, \mu, \nu, \varepsilon$ and the L_{ν} on f_0, ν .

Definition 1.4. Let w_{\pm} be some right- and left- moving transition fronts for (1.1) (or (1.2)) on D. We say that solutions with exponentially decaying initial data converge to time shifts of w_{\pm} if the following hold for any $Y, \mu, \nu > 0$ and $a \in \mathbb{R}$.

(a) If u solves (1.1) (or (1.2)) with initial datum

$$(\theta + \nu)\chi_{(-\infty,a]}(x_1) \le u_0(x) \le e^{-\mu(x_1 - a - Y)},$$

then there is τ_u such that for every $\varepsilon > 0$ there is $s_{\varepsilon} > 0$ such that for each $t \geq s_{\varepsilon}$,

$$||u(t,x) - w_+(t+\tau_u,x)||_{L_x^{\infty}} < \varepsilon \tag{1.5}$$

(and similarly for solutions exponentially decaying as $x_1 \to -\infty$ and for w_-).

(b) There is $L_{\nu} < \infty$ such that if $L \geq L_{\nu}$ and u solves (1.1) (or (1.2)) with initial datum

$$(\theta + \nu)\chi_{[a-L,a+L]}(x_1) \le u_0(x) \le \min\{e^{-\mu(x_1-a-L-Y)}, e^{\mu(x_1-a+L+Y)}\},$$

then there are $\tau_{u,\pm}$ such that for every $\varepsilon > 0$ there is $s_{\varepsilon} > 0$ such that for each $t \geq s_{\varepsilon}$,

$$||u(t,x) - w_{+}(t + \tau_{u,+}, x) - w_{-}(t + \tau_{u,-}, x) + 1||_{L_{\infty}^{\infty}} < \varepsilon.$$
(1.6)

As mentioned before, transition fronts need not exist in general (and are typically not unique) when the domain D is unbounded in more than one variable, even for ignition reactions. We next make several remarks which illuminate the necessity of the assumptions in Theorem 1.3 on cylindrical domains D, and thus show that our result is qualitatively sharp.

- Remarks. 1. The main condition here is $f'_1(0) < c_0^2/4$. Some condition of this type is necessary for existence of fronts, as can be seen from a result of Nolen, Roquejoffre, Ryzhik, and the author in [16]. It shows that there are examples with $D = \mathbb{R}$ and $f'_1(0)$ arbitrarily close to c_0^2 where no transition fronts exist! In these examples, in fact, each global in time solution 0 < u < 1 is a spatially extended pulse with $||u(t, x)||_{L_x^{\infty}} \to 0$ as $t \to -\infty$.
- 2. The condition $f'_1(0) < c_0^2/4$ is equivalent to $2\sqrt{f'_1(0)} < c_0$, meaning that minimal-speed fronts for KPP reactions f(u) with $f'(0) = f'_1(0)$ are slower than the front for $f_0(u)$. It then follows that the graph of $f'_1(0)u$ must intersect that of $f_0(u)$ at some u > 0 because the minimal front speed is monotone with respect to the reaction. Thus we cannot treat the case when f_1 is a KPP reaction, which is not surprising in the light of the previous remark.
- 3. Some condition of non-vanishing of f after it has become large (such as in Definition 1.2) is also needed, otherwise no transition front connecting 0 and 1 might exist. An example of such a situation can be obtained by taking f(x,u) = f(u) with $f(u+\frac{1}{2}) \ge f(u)$ for $u \in [0,\frac{1}{2}]$ (so $f(\frac{1}{2}) = 0$) and strict inequality somewhere. It can be then showed that if there is a front connecting 0 and $a \in [0,1]$, then $a \le \frac{1}{2}$.
- 4. As mentioned earlier, in the general positive reaction case (i) transition fronts are not unique even in homogeneous media.
- 5. We also note that some decay assumption on u_0 in (iii) needs to be made. Indeed, it is not hard to show that if $f(x, u) = f_0(u)$ and u_0 decays slowly enough, then u will "overtake" all time shifts of w_+ .
- 6. $L_{\nu} < \infty$ in Definition 1.4(b) is such that solutions of (1.1) with f_0 in place of f and $u_0(x) \ge (\theta + \nu)\chi_{[-L_{\nu},L_{\nu}]}$ are guaranteed to spread, that is, $u(t,x) \to 1$ locally uniformly in $x \in D$ as $t \to \infty$. This can be taken to be the L_{ν} from Lemma 5.1.

Theorem 1.3 extends to the more general case of (1.2) with periodic q and A which satisfy the following hypotheses.

(H2): The flow $q \in C^{\eta}(D)$ (for some $\eta > 0$) is incompressible $\nabla \cdot q \equiv 0$, p-periodic in x_1 , and with mean-zero first coordinate $\int_{\mathcal{C}} q_1(x) dx = 0$ (where $\mathcal{C} = [0, p] \times \mathbb{T}^{d-1}$). The matrix $A \in C^{1,\eta}(D)$ is symmetric, p-periodic, and with $\underline{A}I \leq A(x) \leq \bar{A}I$ for some $0 < \underline{A} \leq \bar{A} < \infty$ and all $x \in D$.

Again let $c_0 > 0$ be the speed of the unique (right-moving) pulsating front for (1.2) with f replaced by f_0 [1]. We also let $\zeta_0 > 0$ be such that the minimal pulsating front speed for (1.2) with f replaced by $\zeta_0 u(1-u)$ is c_0 [4]. Equivalently, ζ_0 is the unique positive number such that the right hand side of (2.9) below with ζ replaced by ζ_0 equals c_0 . For the left-moving front we have a possibly different speed $c_0^- > 0$, and we define $\zeta_0^- > 0$ accordingly, with e_1, q_1

replaced by $-e_1$, $-q_1$ in (2.10). The condition $f'_1(0) < c_0^2/4$ is now replaced by $f'_1(0) < \zeta_0$ and our main result for (1.2) is as follows.

Theorem 1.5. Let q, A, f satisfy the hypotheses (H1), (H2). Assume that $f \ \zeta$ -majorizes g for some $\zeta < \zeta_0$ and some g as in Definition 1.2.

- (i) If $f'_1(0) < \zeta_0$, then there exists a transition front w for (1.2) moving to the right with $w_t > 0$ (and another moving to the left when ζ_0 is replaced by ζ_0^- in the hypotheses).
- (ii) If f_1 is an ignition reaction (i.e., $\theta' > 0$) and f is non-increasing in u on $[\theta'', 1]$ for some $\theta'' < 1$, then there is a unique (up to time shifts) transition front w_+ for (1.2) moving to the right (and another w_- moving to the left).
- (iii) In the setting of (ii) we have convergence of solutions with exponentially decaying initial data to time shifts of w_{\pm} in the sense of Definition 1.4. Moreover, this convergence is uniform in f, a, u the s_{ε} in the definition depends only on $q, A, f_0, f_1, \zeta, g, K, \theta'', Y, \mu, \nu, \varepsilon$ and the L_{ν} on q, A, f_0, ν .
- Remarks. 1. Although the front speed is not well defined in general disordered media, it is easy to see from our proof that the reaction zone of w_{\pm} moves with speed $\geq c_0$ resp. $\geq c_0^-$. We also prove an f-independent bound on the width of w_{\pm} (see the remark after the proof of Lemma 3.1 and the last paragraph of Section 3). In addition, w_{+} decays at an exponential rate $\geq \lambda_{\zeta}$ as $x_1 \to \infty$ (and similarly w_{-} as $x_1 \to -\infty$) determined from (2.9) (see (3.1)). In fact, λ_{ζ} can be replaced here by any λ such that the fraction in (2.9) is smaller than c_0 .
- 2. The proof of Theorem 1.5 can, in fact, be made independent of previous results on transition fronts. In particular, we can prove that a unique pulsating front with a unique speed c_0 exists for (1.2) with the x-independent reaction f_0 (which we need in order to state Theorem 1.5). This is done at the end of the proof of Lemma 5.1. The only result we will need is convexity of the function $\kappa(\lambda)$ in (2.10) and $\kappa'(0) = 0$, from [1, Proposition 5.7(iii)].

We next note that (at least for fixed periodic A, q) the ignition fronts in (ii) satisfy the earlier mentioned definition of Matano where the shape of the front depends continuously on the reaction f in the neighborhood of the reaction zone. This follows from the uniform in time convergence of general solutions to the front in (iii). We provide an application of this principle to periodic and random media.

Corollary 1.6. Assume the hypotheses of Theorem 1.5(ii) and that f is also p-periodic in x_1 . Then there is $c_{\pm} > 0$ such that the unique transition front w_{\pm} from that theorem satisfies $w_{\pm}(t+p/c_{\pm}, x \pm pe_1) = w_{\pm}(t, x)$ (i.e., w_{\pm} is time-periodic in the frame moving right/left at the speed c_{\pm} , which is thus the speed of w_{\pm}).

Remark. Existence, uniqueness, and periodicity of the transition front in this setting also follows from results in [1, 2], but only within the class of transition fronts which have a constant mean speed (i.e., there is $c \geq 0$ such that for \tilde{X}_w as below, $\sup_{(t,s)} \left| |\tilde{X}_w(t) - \tilde{X}_w(s)| - c|t-s| \right| < \infty$), with the extra assumption $\sup_{x \in D} f(x,u) > 0$ for each $u \in (\inf_{x \in D} \theta(x), 1)$, and with somewhat stronger regularity assumptions on q, A, f (the last two in [1]). Hence, our result of existence, uniqueness, and periodicity of pulsating fronts in periodic media with ignition reactions is new in this generality.

Proof. Consider only w_+ . Let u, u' be solutions of (1.2) with initial data $u_0(x), u_0(x+p)$, where u_0 is as in the first part of (iii). Then u, u' are p-translates of each other (in x_1), so the same is true about $w_+(t+\tau_u,x)$ and $w_+(t+\tau_{u'},x)$. The result follows with $c_+\equiv p/(\tau_u-\tau_{u'})$. \square

Although fronts in disordered media do not have constant (mean) speeds in general, our results can be used to show that there is a deterministic asymptotic speed of fronts for (1.1)/(1.2) when f is random (stationary and ergodic with respect to translations in x_1). The asymptotic front speed of a transition front w is defined as

$$c \equiv \lim_{|t| \to \infty} \frac{|\tilde{X}_w(t)|}{|t|},\tag{1.7}$$

provided the limit exists. Here $\tilde{X}_w(t)$ is the first coordinate of some point in the reaction zone of the front, for instance, one could take $\tilde{X}_w(t)$ such that $w(t,x) = \frac{1}{2}$ for some $x = (\tilde{X}_w(t), x_2, \ldots, x_d)$. Such $\tilde{X}_w(t)$ may not be unique but the requirement of a bounded width of the front w shows that the limit in (1.7) is independent of the choice.

Corollary 1.7. Consider a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ and assume that a measurable function $f: \Omega \to L^{\infty}_{loc}(D \times [0,1])$ satisfies the hypotheses of Theorem 1.5(ii) (with ω -independent q, A) uniformly in $\omega \in \Omega$. In addition, assume that f is stationary and ergodic. That is, there is a group $\{\pi_k\}_{k\in\mathbb{Z}}$ of measure preserving transformations acting ergodically on Ω such that $f(\pi_k\omega; x, u) = f(\omega; x - kpe_1, u)$. Then there are constants $c_{\pm} > 0$ such that the ω -dependent transition front $w_{\pm,\omega}$ from that theorem has asymptotic speed c_{\pm} for almost all $\omega \in \Omega$.

Remarks. 1. Theorem 1.5(iii) then shows that solutions with (large enough) exponentially decaying initial data almost surely spread with asymptotic speed c_+ to the right and c_- to the left.

2. Corollary 1.7 was proved in [15] for (1.1) in the one-dimensional setting $D = \mathbb{R}$ and with the random reaction function $f(\omega; x, u) = b(x, \omega) f_0(u)$, where b is bounded below and above by positive constants and f_0 is of ignition type.

Proof. Consider only w_+ and the case of p-periodic q, A. Let v(x) be the function from Lemma 2.1 below with $\tilde{\theta}$ chosen as at the beginning of Section 3. Let u_m solve (1.2) with initial condition $u_m(0,x) = v(x - mpe_1)$ (so that $(u_m)_t > 0$). For integers $n \geq m$ define

$$\tau_{m,n}(\omega) \equiv \inf \{ t \ge 0 \mid u_m(t,x) \ge v(x - npe_1) \text{ for all } x \in D \}.$$

Then the proof of Theorem 1.5 (more precisely, (3.6) and (3.9) below) shows that $\tau_{m,n}(\omega) \in [C_0(n-m), C_1(n-m)]$ for some $0 < C_0 < C_1 < \infty$ and all ω . Moreover, $\tau_{m,n}$ is measurable because $f: \Omega \to L^{\infty}_{loc}(D \times [0,1])$ is measurable and $\tau_{m,n}: L^{\infty}_{loc}(D \times [0,1]) \to \mathbb{R}$ is a lower semi-continuous function of the reaction due to continuity of solutions of (1.2) with respect to $f \in L^{\infty}_{loc}(D \times [0,1])$ and the properties of v.

The comparison principle shows that $\tau_{m,n}(\omega) \leq \tau_{m,k}(\omega) + \tau_{k,n}(\omega)$ when $m \leq k \leq n$. We also have $\tau_{m+k,n+k}(\pi_k\omega) = \tau_{m,n}(\omega)$ for $k \in \mathbb{Z}$. Since the group $\{\pi_k\}_{k\in\mathbb{Z}}$ acts ergodically on Ω , the subadditive ergodic theorem [10, 12] shows that there is $\tau_+ \in [C_0, C_1]$ such that

$$\tau_{+} = \lim_{n \to \infty} \frac{\tau_{0,n}(\omega)}{n} = \lim_{n \to \infty} \frac{\tau_{-n,0}(\omega)}{n}$$

for almost all ω . Uniform convergence (in ω) of the solution u_0 to the front $w_{+,\omega}$ and the proof of Theorem 1.5 (more precisely, (3.1) below) then show that $c_+ = p/\tau_+$ is the asymptotic speed of $w_{+,\omega}$ for almost all ω .

Let us finish this introduction with a brief description of the proof of Theorem 1.5. In Section 2 we construct the front as a limit of a (sub)sequence of special solutions u_n of (1.2), increasing in time and initially (at a sequence of times $\tau_n \to -\infty$) supported increasingly farther to the left. The sequence τ_n is chosen so that the reaction zone for u_n arrives at the origin at t = 0. The main issue is to show that the u_n have a uniformly (in n) bounded width in the sense of Definition 1.1. This will boil down to showing that the leftmost point $Y_n(t)$ such that to the right of it u_n decays no slower than at a fixed exponential rate, cannot escape too far to the right from the leftmost point $X_n(t)$ such that to the right of it $u_n(t,x) < \alpha_f(x)$. When going from left to right, the former point can be thought of as the beginning of the region where u_n is 'small' (for obvious reasons), while the latter as the end of the region where u_n is 'large enough' (because the end of the region where u_n is close to 1 will be within a uniform distance from $X_n(t)$, thanks to $f \zeta$ -majorizing g and thus u_n growing close to 1 within a uniformly bounded time after it becomes at least $\alpha_f(x)$ — see Lemma 2.6(i)). The uniform boundedness of $Y_n(t) - X_n(t)$ will be proved in Lemma 2.5 using the crucial hypothesis $f_1'(0) < \zeta_0$, which guarantees that the reaction is so small when $u_n(t,x) < \alpha_f(x)$ that it can make $Y_n(t)$ propagate to the right only with a speed strictly smaller than c_0 when $Y_n(t) - X_n(t)$ is large. At the same time $X_n(t)$ always propagates to the right with speed at least c_0 because that is the case for the region where u_n is close to 1 (see Lemma 2.6(i)). Thus $Y_n(t) - X_n(t)$ cannot grow unbounded, as desired.

We note that this argument requires q and A to be periodic because the upper bound on the speed of propagation of $Y_n(t)$ is the speed of propagation of an exponentially decreasing super-solution (at small u) of (1.2), which is a solution if f(x, u) is replaced by ζu . We were able to construct this solution when the latter PDE is x-periodic by solving the cell problem (2.10) below but not for general non-periodic q, A. Similarly, the lower bound c_0 on the speed of propagation of $X_n(t)$ is the minimal front speed for the x-periodic PDE with f(x, u) replaced by $f_0(u)$. It is possible that if one could show that these bounds depend on some local norms of q, A but not on q, A themselves, and primarily not on the period p, then one could obtain a front for non-periodic q, A as a limit of fronts for a sequence of periodic q_k , A_k with growing periods. We leave this as an open problem.

In Section 3 we show that if w is any transition front, then for each u_n one obtains L_x^{∞} convergence of u_n to a time-shift of w as $t \to \infty$, and the rate of this convergence is uniform
in n. This will show that any two fronts must be time shifts of each other. The argument
starts with proving local stability of u_n via construction of sub- and super-solutions near u_n ,
which will trap w between small perturbations of two time-shifts of u_n . One then uses the
strong maximum principle to obtain convergence of a fixed time shift of w to u_n . We note
that a similar strategy was employed in the proof of uniqueness of 1-dimensional ignition
fronts [13]. However, our argument applies in a more general setting, and its second part is
different from as well as considerably simpler and shorter than that in [13]. This is in part

thanks to the use of the ideas from Section 2 described above, and in part because unlike [13], we do not establish a rate of the convergence of u_n to a time shift of w.

We recycle this argument in Section 4 to show L_x^{∞} -convergence of u_n to a time shift of any solution u as in Theorem 1.5(iii), again at a uniform rate (in u). Since u_n also converges uniformly to a time shift of w, the same will be true for u.

Of course, this all also proves the special case, Theorem 1.3. Finally, in Appendix A we show how to make our proof independent of previous results by using our arguments to obtain a slight improvement of Lemma 2.6(ii) below, which is from [21].

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2. Existence of Fronts for General Reactions

In this section we will prove Theorem 1.5(i) by finding a front moving to the right. Let us assume that the period of q, A in x_1 is p=1 and thus the unit cell of periodicity is $C = \mathbb{T}^d$ (the general case is treated identically). We will assume without loss that $\zeta > f'_1(0)$ and that there is $\sigma \in (0, \zeta - f'_1(0))$ such that $f(\zeta')$ -majorizes $g(\zeta')$ for each $\zeta' > \zeta - \sigma(\zeta')$. This can be done because if $\zeta' > \zeta$, then $f(\zeta')$ -majorizes $g(\zeta')$, so we only need to change $\zeta(\zeta')$ to $(\max\{\zeta, f'_1(0)\} + \zeta_0)/2$ and pick $\sigma = (\zeta_0 - \max\{\zeta, f'_1(0)\})/4$. The first assumption implies that inf can be replaced by min and ζ' by ζ' is uniformly bounded away from 0 and 1 (see (2.8)). The second guarantees that if a subsequence of reactions ζ' which satisfy (H1) converges locally uniformly to ζ' (each such sequence has a locally uniformly convergent subsequence) and each ζ' majorizes ζ' for each $\zeta' > \zeta - \sigma(\zeta')$, then this ζ' not only satisfies (H1) but also ζ' -majorizes ζ' for each $\zeta' > \zeta - \sigma(\zeta')$. This claim would not be true with ζ' in place of ζ' . We note that all constants in this section will depend on ζ' , ζ' , ζ' , ζ' , ζ' (also on ζ' , ζ' , ζ' , which already only depend on ζ' , ζ' , but not on ζ' .

As in the one-dimensional case [14, 15, 18], we will look for the front as a limit of solutions with initial data specified at increasingly negative times and supported further and further to the left. These solutions will be monotonically increasing in time. We will therefore need

Lemma 2.1. For each $\tilde{\theta} \in (\theta, 1)$ there exists a function v supported in $(-\infty, 0) \times \mathbb{T}^{d-1} \subseteq D$ with $v(x) \in [0, \tilde{\theta} \text{ for all } x \in D \text{ and } v(x) = \tilde{\theta} \text{ for all small enough } x_1 \text{ such that}$

$$-\operatorname{div}(A(x)\nabla v) + q(x)\cdot\nabla v \le f_0(v) \tag{2.1}$$

in the sense of distributions.

Proof. Take a nondecreasing function $\rho \in C(\mathbb{R}) \cup C^2(\mathbb{R}^+)$ with $\rho(v) = 0$ for $v \leq 0$, $\rho(v) = v$ for $v \in [0, \frac{\theta + \tilde{\theta}}{2}]$, $\rho''(v) \leq 0$ for $v \in [\frac{\theta + \tilde{\theta}}{2}, 1]$, and $\rho(v) = \tilde{\theta}$ for $v \geq 1$. Let $\tilde{v} < 0$ be a C^2 solution of $-\operatorname{div}(A\nabla \tilde{v}) + q \cdot \nabla \tilde{v} = q_1 - \operatorname{div}(Ae_1)$ on \mathbb{T}^d , periodically

Let $\tilde{v} < 0$ be a C^2 solution of $-\operatorname{div}(A\nabla \tilde{v}) + q \cdot \nabla \tilde{v} = q_1 - \operatorname{div}(Ae_1)$ on \mathbb{T}^d , periodically continued to D (here $e_1 = (1, 0, \dots, 0)$ and $q_1 = q \cdot e_1$). Such \tilde{v} exists because the integral of the right hand side is zero, and because the left hand side annihilates constants. Then we let $v_{\varepsilon}(x) \equiv \varepsilon(\tilde{v}(x) - x_1)$ and $v(x) \equiv \rho(v_{\varepsilon}(x))$ for $\varepsilon > 0$, so that v is supported in $(-\infty, 0) \times \mathbb{T}^{d-1}$. We have $-\operatorname{div}(A\nabla v_{\varepsilon}) + q \cdot \nabla v_{\varepsilon} = 0$ and so for some distribution $T \geq 0$ supported on the set

 $D_{\varepsilon} \equiv \{x \in D \mid v_{\varepsilon}(x) = 0\},\$

$$-\operatorname{div}(A\nabla v) + q \cdot \nabla v = -\varepsilon^2 \chi_{D \setminus D_{\varepsilon}} \rho''(v_{\varepsilon})(\nabla \tilde{v} - e_1) \cdot A(\nabla \tilde{v} - e_1) - T.$$

If $\rho''(v_{\varepsilon}(x)) < 0$, then $v(x) = \rho(v_{\varepsilon}(x)) \in [\frac{\theta + \tilde{\theta}}{2}, \tilde{\theta}]$. Since f_0 is uniformly positive on this interval and A is a positive matrix, (2.1) follows provided $\varepsilon > 0$ is small enough.

We now fix $\tilde{\theta} < 1$ to be close to 1 so that $\tilde{\theta} > \theta_0$ from (2.8) below (in particular, $\tilde{\theta} > \theta$) and consider the corresponding v along with the functions $v_n(x) \equiv v(x + ne_1)$. For all $n \in \mathbb{N}$ let u_n solve (1.2) for $t > \tau_n$ with initial condition $u_n(\tau_n, x) = v_n(x)$, where $\tau_n \to -\infty$ will be chosen shortly. We then have

Lemma 2.2. The functions u_n satisfy for all $t > \tau_n$ and $x \in D$

$$(u_n)_t(t,x) > 0 (2.2)$$

as well as

$$\lim_{t \to \infty} u_n(t, x) = 1 \tag{2.3}$$

locally uniformly in $x \in D$.

Proof. The time derivative $\dot{u}_n \equiv (u_n)_t$ satisfies $\dot{u}_n(\tau_n, x) \geq 0$ due to (2.1) and periodicity of q, A, and it is not identically 0. Since \dot{u}_n satisfies $(\dot{u}_n)_t + q \cdot \nabla \dot{u}_n = \operatorname{div}(A \nabla \dot{u}_n) + \frac{\partial f}{\partial u}(x, u_n) \dot{u}_n$ and $\frac{\partial f}{\partial u}$ is bounded, the strong maximum principle shows (2.2). This means that for each n the function $\tilde{u}(x) \equiv \lim_{t\to\infty} u_n(t,x)$ is well defined and satisfies

 $\tilde{u}(x) \in (0,1]$ and

$$-\operatorname{div}(A\nabla \tilde{u}) + q \cdot \nabla \tilde{u} = f(x, \tilde{u}).$$

Lemma 2.3(ii) below now shows that \tilde{u} is a constant, which is then 1 due to $\|\tilde{u}\|_{\infty} \geq \theta > \theta$. Parabolic regularity then shows the limit to be locally uniform in D.

We now choose $\tau_n < 0$ to be the unique time such that

$$u_n(0,0) = \theta. (2.4)$$

Note that $\tau_n \to -\infty$ (see the remark after Lemma 2.4).

Despite its apparent simplicity, we were not able to locate the following Liouville-type result in the literature.

Lemma 2.3. (i) Let q, A be as in (H2) but without the assumptions of periodicity and q_1 being mean-zero. If the function u is bounded on D and satisfies

$$-\operatorname{div}(A(x)\nabla u) + q(x)\cdot\nabla u = 0, \tag{2.5}$$

then u is constant.

(ii) Let q, A be as in (H2) but without the assumption of periodicity, and let r be a bounded non-negative measurable function on D. If u is bounded and non-negative on D and satisfies

$$-\operatorname{div}(A(x)\nabla u) + q(x)\cdot\nabla u = r(x)u, \tag{2.6}$$

then u is constant.

Remark. Note that if q is incompressible on D, then its mean $\bar{q} \equiv \int_{\mathbb{T}^{d-1}} q(x_1, x') dx'$ is independent of x_1 . Thus $\int_{[0,p]\times\mathbb{T}^{d-1}} q_1(x) dx = 0$ is satisfied in (ii) for either all p > 0 or none.

Proof. Let u satisfy (2.5) or (2.6). Then $\tilde{u}(x_1) \equiv \min_{x' \in \mathbb{T}^{d-1}} u(x_1, x')$ cannot have a local minimum by the maximum principle unless it is constant. In either case the limits

$$\lim_{x_1 \to -\infty} \tilde{u}(x_1) = l_1 \quad \text{and} \quad \lim_{x_1 \to \infty} \tilde{u}(x_1) = l_2$$

exist. Harnack inequality for the domains $(y, y + 1) \times \mathbb{T}^{d-1} \subset (y - 1, y + 2) \times \mathbb{T}^{d-1}$ with $y \to \pm \infty$ (see, e.g., [9, p. 199]) now shows that uniformly in $x' = (x_2, \dots, x_d)$,

$$\lim_{x_1 \to -\infty} u(x) = l_1 \quad \text{and} \quad \lim_{x_1 \to \infty} u(x) = l_2$$

Parabolic regularity shows that $\lim_{x_1\to\pm\infty} \nabla u(x) = 0$ uniformly in x'.

If $\bar{q}_1 = 0$, integrate (2.6) over D to get

$$0 = \int_{D} \operatorname{div}(A\nabla u) = \int_{D} \operatorname{div}(qu) - \int_{D} ru = (l_{2} - l_{1})\bar{q}_{1} - \int_{D} ru = -\int_{D} ru.$$

Thus $ru \equiv 0$ and (2.6) becomes (2.5). We multiply (2.5) by u and integrate over D to get

$$-\int_{D} \nabla u \cdot A \nabla u = \int_{D} u \operatorname{div}(A \nabla u) = \frac{1}{2} \int_{D} \operatorname{div}(q u^{2}) = \frac{1}{2} (l_{2}^{2} - l_{1}^{2}) \bar{q}_{1} = 0.$$
 (2.7)

Thus u must be constant, proving (ii) and the case $\bar{q}_1 = 0$ in (i).

If $\bar{q}_1 \neq 0$ in (i), integrate (2.5) over D to get

$$0 = \int_D \operatorname{div}(A\nabla u) = \int_D \operatorname{div}(qu) = (l_2 - l_1)\bar{q}_1.$$

Thus $l_2 = l_1$ and (2.7) finishes the proof of this case.

We will now recover the transition front w as a limit of the u_n along a subsequence as $n \to \infty$. Such a limit always exists by parabolic regularity and satisfies $w(0,0) = \theta$ due to (2.4), but the main issue is to show that it is indeed a transition front for (1.2). The following four lemmas will ensure this fact.

Let us take ζ from the statement of Theorem 1.5 and let θ_j (j = 0, 1) be the smallest positive number such that $f_j(\theta_j) = \zeta \theta_j$. Since $\zeta > f'_1(0)$, we have $0 < \theta_1 \le \theta_0 < 1$, $\theta < \theta_0$, and for each $x \in D$,

$$\alpha_f(x) \in [\theta_1, \theta_0]. \tag{2.8}$$

We now let

$$c_{\zeta} \equiv \min_{\lambda > 0} \frac{\zeta + \kappa(\lambda)}{\lambda} \tag{2.9}$$

with $\kappa(\lambda)$ and $\gamma(x;\lambda) > 0$ the principal eigenvalue and eigenfunction for

$$\operatorname{div}(A\nabla\gamma) - (q + 2\lambda Ae_1) \cdot \nabla\gamma + (\lambda^2 e_1^T Ae_1 - \lambda \operatorname{div}(Ae_1) + \lambda q_1)\gamma = \kappa(\lambda)\gamma \tag{2.10}$$

on \mathbb{T}^d , normalized by $\sup_{x \in \mathbb{T}^d} \gamma(x; \lambda) = 1$ (for (1.1) we have $\gamma(x; \lambda) \equiv 1$ and $\kappa(\lambda) = \lambda^2$). We note that the minimum is achieved at some $\lambda_{\zeta} > 0$ because κ is a continuous function of λ and the fraction in (2.9) diverges to ∞ as $\lambda \to 0, \infty$. The latter follows from $\zeta > 0$ and

$$\kappa(\lambda) \geq \underline{A}\lambda^2$$
,

which is obtained after dividing (2.10) by γ and integrating over \mathbb{T}^d :

$$\kappa(\lambda) = \int_{\mathbb{T}^d} (\nabla \log \gamma - \lambda e_1) \cdot A(\nabla \log \gamma - \lambda e_1) \ge \underline{A} \int_{\mathbb{T}^d} |\nabla \log \gamma|^2 + \lambda^2 \ge \underline{A} \lambda^2.$$

We note that $c_{\zeta} < c_0$ because $\zeta < \zeta_0$ and ζ_0 was defined so that the right hand side of (2.9) with ζ_0 in place of ζ equals c_0 .

We continue $\gamma(x; \lambda_{\zeta})$ periodically on D, and define

$$\Psi(s,x) \equiv \left[\inf_{D} \gamma(x;\lambda_{\zeta})\right]^{-1} e^{-\lambda_{\zeta} s} \gamma(x;\lambda_{\zeta}) > 0$$

(for (1.1) this is $\Psi(s,x) = e^{-\lambda_{\zeta}s}$). Notice that $\Psi(0,x) \geq 1$, and $\psi(t,x) \equiv \Psi(x_1 - c_{\zeta}t,x)$ is an (exponentially growing as $x_1 \to -\infty$) pulsating front with speed c_{ζ} for (1.2) with f replaced by ζu . In fact, [4] shows that c_{ζ} is also the minimal speed of a true pulsating front for (1.2) with any x-independent KPP reaction \tilde{f} satisfying $\tilde{f}'(0) = \zeta$. This is why $c_{\zeta} < c_0$ will be a crucial component of our argument (and, in fact, any λ such that the fraction in (2.9) is smaller than c_0 would do in place of λ_{ζ}).

We now let for each $n \in \mathbb{N}$ and $t \geq \tau_n$,

$$X_n(t) \equiv \sup\{x_1 \mid u_n(t,x) \ge \alpha_f(x) \text{ for some } x = (x_1, x')\}$$

and

$$Y_n(t) \equiv \inf\{y \mid u_n(t, x) \le \Psi(x_1 - y, x) \text{ for all } x \in D\}.$$

Both these functions are non-decreasing because u_n is increasing and Ψ decreasing in their respective first variables. Continuity of $u_n(t,x)$, lower semi-continuity of $\alpha_f(x)$, and compactness of \mathbb{T}^{d-1} imply that X_n is continuous from the right.

Lemma 2.4. Let $\xi \equiv \sup_{u \in (0,1)} f_1(u)/u \ge \zeta$ and $c_{\xi} \equiv (\xi + \kappa(\lambda_{\zeta}))/\lambda_{\zeta}$. Then for any n and $t \ge \tau \ge \tau_n$ we have

$$Y_n(t) - Y_n(\tau) \le c_{\xi}(t - \tau). \tag{2.11}$$

Remark. Taking $\tau = \tau_n$ and t = 0, this and (2.4) give $\tau_n \to -\infty$.

Proof. This is immediate from the definition of $Y_n(\tau)$ and the fact that the function $\phi(t,x) \equiv \Psi(x_1 - Y_n(\tau) - c_{\xi}(t - \tau), x)$ is a supersolution of (1.2) for $t > \tau$ (it solves (1.2) with ξu in place of f).

This and $Y_n(\tau_n) < \infty$ show that $Y_n(t) < \infty$ (then also $X_n(t) < \infty$ by (2.8)) and that Y_n is continuous because it is non-decreasing. The following is a crucial step in our proof of existence of fronts.

Lemma 2.5. There is $C_1 < \infty$ such that for all $n \in \mathbb{N}$ and all $t \geq \tau_n$,

$$|Y_n(t) - X_n(t)| \le C_1. (2.12)$$

Proof. The bound $X_n(t) - Y_n(t) \le C_1$ for a large enough C_1 is obvious from the definitions of X_n, Y_n , and (2.8), so let us show $Y_n(t) - X_n(t) \le C_1$. We let C_1 (to be chosen later) be larger than $C_0 \equiv \max\{Y_n(\tau_n) - X_n(\tau_n), 1\}$, the latter independent of n and finite due to v_n being supported in a half-strip, $\tilde{\theta} > \theta_0$, and (2.8). Let us assume that $Y_n(t_1) - X_n(t_1) > C_1$ for some n and $t_1 \ge \tau_n$, and let $t_0 \equiv \max\{t < t_1 \mid Y_n(t) - X_n(t) \le C_0\} \ge \tau_n$. The maximum exists and $Y_n(t_0) - X_n(t_0) = C_0$ because Y_n is continuous and X_n non-decreasing.

Let $\tilde{Y}_n(t) \equiv Y_n(t_0) + c_{\zeta}(t-t_0)$ and if $X_n(t) \geq \tilde{Y}_n(t)$ for some $t \in [t_0, t_1]$, let t_2 be the first such time (recall that X_n is continuous from the right). Then $\psi(t, x) \equiv \Psi(x_1 - \tilde{Y}_n(t), x)$ is a solution and $u_n(t, x)$ a subsolution in $\tilde{D} \equiv \{(t, x) \in (t_0, t_2) \times D \mid x_1 > \tilde{Y}_n(t)\}$ of (1.2) with ζu in place of f. Moreover, $\psi(t_0, x) \geq u_n(t_0, x)$ for $x_1 > \tilde{Y}_n(t_0)$ and $\psi(t, x) \geq u_n(t, x)$ for $t \in [t_0, t_2]$ and $x_1 = \tilde{Y}_n(t)$ (because $\Psi(0, x) \geq 1$). Thus $\psi(t, x) \geq u_n(t, x)$ in \tilde{D} by the comparison principle, meaning that $Y_n(t) \leq \tilde{Y}_n(t)$ for $t \in [t_0, t_2]$. But then $Y_n(t_2) - X_n(t_2) \leq 0 < C_0$, a contradiction with the choice of t_0 . Hence such time $t_2 \in [t_0, t_1]$ does not exist and the above argument gives $Y_n(t) \leq \tilde{Y}_n(t)$ for $t \in [t_0, t_1]$.

That is, Y_n increases with average speed at most c_{ζ} on $[t_0, t_1]$. On the other hand, the next lemma shows that X_n increases with average speed at least $c_0 - \varepsilon$ (for any $\varepsilon > 0$) after an initial time delay t_{ε} (in the sense of (2.13) below with $\tau = t_0$). Since $c_{\zeta} < c_0$, we can pick $\varepsilon \equiv (c_0 - c_{\zeta})/2 > 0$ and obtain

$$Y_n(t)-X_n(t) \leq Y_n(t_0)+c_{\zeta}(t-t_0)-X_n(t_0)-(c_0-\varepsilon)(t-t_0-t_{\varepsilon}) = C_0+(c_0-\varepsilon)t_{\varepsilon}-(c_0-\varepsilon-c_{\zeta})(t-t_0)$$
 for $t \in [t_0,t_1]$. If we now let $C_1 \equiv C_0+c_0t_{\varepsilon}$, then it follows that $Y_n(t_1)-X_n(t_1) < C_1$, a contradiction. Thus $Y_n(t)-X_n(t) \leq C_1$ for a large enough C_1 and all n and $t \geq \tau_n$.

Recall that $c_0^- > 0$ is the speed of the unique left-moving front for (1.2) with reaction f_0 .

Lemma 2.6. (i) For every $\varepsilon > 0$ there is $t_{\varepsilon} < \infty$ such that if $u : [0, \infty) \times D \to [0, 1]$ solves (1.2) with $u_t \geq 0$ and $u(0, \tilde{x}) \geq \alpha_f(\tilde{x})$ for some $\tilde{x} \in D$, then for each $t \geq 0$ we have

$$\inf \left\{ u(t+t_{\varepsilon},x) \mid x_1 - \tilde{x}_1 \in [-(c_0^- - \varepsilon)t, (c_0 - \varepsilon)t] \right\} \ge 1 - \varepsilon.$$

(ii) There is L' > 0 such that for every $\varepsilon > 0$ there is $t'_{\varepsilon} < \infty$ satisfying the following. If $u: [0, \infty) \times D \to [0, 1]$ solves (1.2) and $\inf \{u(0, x) \mid |x_1 - \tilde{x}_1| \leq L'\} \geq (1 + \theta)/2$ for some $\tilde{x} \in D$, then for each $t \geq 0$ we have

$$\inf \left\{ u(t + t'_{\varepsilon}, x) \mid x_1 - \tilde{x}_1 \in [-(c_0^- - \varepsilon)t, (c_0 - \varepsilon)t] \right\} \ge 1 - \varepsilon.$$

Remarks. 1. Part (i) also shows that for $t \geq \tau \geq \tau_n$,

$$X_n(t) - X_n(\tau) \ge (c_0 - \varepsilon)(t - \tau - t_{\varepsilon}).$$
 (2.13)

2. Of course, the constants $t_{\varepsilon}, t'_{\varepsilon}$ are independent of f.

Proof. (ii) This is an immediate consequence of Proposition 3.4 in [21] which proves the same result for $f(x, u) = f_0(u)$. In Appendix A we will provide an alternative proof of this result (in fact, with $c_0 - \varepsilon$, $c_0^- - \varepsilon$ replaced by c_0 , c_0^-), thus making our proof independent of [21].

(i) Consider σ from the beginning of this section and let

$$\alpha'_f(x) \equiv \inf\{u \in (0,1] \mid f(x,u) \ge (\zeta - \sigma/2)u\}.$$

Then $\alpha'_f(x) \leq \alpha_f(\tilde{x}) - \sigma\theta_1/4K$ for $x \in B_{\sigma\theta_1/4K}(\tilde{x})$ since by the definition of $\alpha_f(\tilde{x})$ and (2.8),

$$f\left(x,\alpha_f(\tilde{x}) - \frac{\sigma\theta_1}{4K}\right) \ge \zeta\alpha_f(\tilde{x}) - \frac{\sigma\theta_1}{2} \ge \left(\zeta - \frac{\sigma}{2}\right)\alpha_f(\tilde{x}) \ge \left(\zeta - \frac{\sigma}{2}\right)\left(\alpha_f(\tilde{x}) - \frac{\sigma\theta_1}{4K}\right)$$

for those x. (Also, $\zeta - \sigma/2 > f_1'(0)$ shows that $\alpha_f'(x)$ is the minimum of the set in its definition, as well as uniformly bounded away from 0.) So $u(0, \tilde{x}) \geq \alpha_f(\tilde{x})$, $u_t \geq 0$, and parabolic regularity show that there is $\delta > 0$ (independent of f, \tilde{x}) such that

$$u(1,x) \ge \alpha_f'(x)$$
 for all $x \in B_\delta(\tilde{x})$. (2.14)

Assume that (i) is false. Taking L' from (ii) and using (ii), this means that for each $n \in \mathbb{N}$, there is a solution w_n of (1.2) with some reaction f_n satisfying all the hypotheses (in particular, f_n ζ' -majorizes g for each $\zeta' > \zeta - \sigma$), such that $w_n(0, \tilde{x}^n) \ge \alpha_{f_n}(\tilde{x}^n)$, $(w_n)_t \ge 0$, and $w_n(n, x^n) < (1 + \theta)/2$ for some x^n with $|x_1^n - \tilde{x}_1^n| \le L'$. After possible translation in x_1 it is sufficient to consider $\tilde{x}^n \in \mathbb{T}^d$, so we can assume $\tilde{x}^n \to \tilde{x}$ (otherwise we choose a subsequence). Then (2.14) gives $w_n(1, \tilde{x}) \ge \alpha'_{f_n}(\tilde{x})$ for large n.

By parabolic regularity, the functions w_n are uniformly bounded in $C^{1,\eta;2,\eta}([1,\infty)\times D)$ for some $\eta>0$. Thus there is a subsequence (which we again denote w_n) converging in $C^{1,2}_{loc}([1,\infty)\times D)$ to a solution $\tilde{w}\geq 0$ of (1,2) on $(1,\infty)\times D$ with some reaction f (local uniform limit of f_n), satisfying $\tilde{w}_t\geq 0$. But then $w(x)\equiv \lim_{t\to\infty}\tilde{w}(t,x)$ exists and satisfies

$$q \cdot \nabla w = \operatorname{div}(A\nabla w) + f(x, w).$$

Lemma 2.3 and boundedness of f(x,u)/u show that w is a constant and $f(x,w(x)) \equiv 0$. This constant is then 1 because $w(\tilde{x}) \geq \liminf_n \alpha'_{f_n}(\tilde{x}) \geq \alpha'_f(\tilde{x})$ and $f(\zeta - \sigma/2)$ -majorizes $g(\zeta)$ (being local uniform limit of f_n). But $w_n(n,x^n) < (1+\theta)/2$ and $(w_n)_t \geq 0$ show for all $t \geq 1$,

$$||1 - \tilde{w}(t, \cdot)||_{L^{\infty}([\tilde{x}_1 - L', \tilde{x}_1 + L'])} \ge \frac{1 - \theta}{2} > 0.$$

Parabolic regularity again shows that this contradicts $w \equiv 1$, thus finishing the proof.

For $\varepsilon > 0$ we define

$$Z_{n,\varepsilon}^{-}(t) \equiv \sup \left\{ y \mid u_n(t,x) \ge 1 - \varepsilon \text{ when } x_1 \le y \right\},$$

$$Z_{n,\varepsilon}^{+}(t) \equiv \inf \left\{ y \mid u_n(t,x) \le \varepsilon \text{ when } x_1 \ge y \right\}.$$

Clearly both are finite. The following will ensure a bounded width of the constructed transition front.

Lemma 2.7. For any $\varepsilon \in (0, \min\{c_0, c_0^-\})$ let $\tilde{t}_{\varepsilon} \equiv t_{\varepsilon} + C_1(\min\{c_0, c_0^-\} - \varepsilon)^{-1}$ with t_{ε} from Lemma 2.6(i) and C_1 from Lemma 2.5. Then there is $L_{\varepsilon} < \infty$ such that for all n and $t \geq \tau_n + \tilde{t}_{\varepsilon}$,

$$Z_{n,\varepsilon}^+(t) - Z_{n,\varepsilon}^-(t) \le L_{\varepsilon}.$$

Proof. Notice that Lemma 2.5 and continuity of Y_n show that if X_n has jumps, they cannot be larger than $2C_1$. This and $(u_n)_t > 0$ mean that for each $t \geq \tau_n + \tilde{t}_{\varepsilon}$ and any closed subinterval I of $(-\infty, X_n(t - \tilde{t}_{\varepsilon})]$ of length $2C_1$ there is $x_1 \in I$ and $x' \in \mathbb{T}^{d-1}$ with $u_n(t - \tilde{t}_{\varepsilon}, x_1, x') \geq \alpha_f(x_1, x')$. Then Lemma 2.6(i) shows that $u_n(t, x) \geq 1 - \varepsilon$ whenever $x_1 \leq X_n(t - \tilde{t}_{\varepsilon}) + C_1$. On the other hand, Lemmas 2.4 and 2.5 show that $u_n(t, x) \leq \varepsilon$ whenever

$$x_1 \ge X_n(t - \tilde{t}_{\varepsilon}) + C_1 + c_{\xi}\tilde{t}_{\varepsilon} + l_{\varepsilon},$$

where l_{ε} is such that $\Psi(l_{\varepsilon}, x) \leq \varepsilon$ for all $x \in D$. Thus $L_{\varepsilon} \equiv c_{\varepsilon} \tilde{t}_{\varepsilon} + l_{\varepsilon}$ works.

Having Lemma 2.7, the proof of Theorem 1.5 is now standard. Parabolic regularity shows that the functions u_n are uniformly bounded in $C^{1,\eta;2,\eta}([\tau_n+1,\infty)\times D)$, so we can find a subsequence converging in $C^{1;2}_{loc}(\mathbb{R}\times D)$ to a function w on $\mathbb{R}\times D$ which then is also a solution of (1.2). Moreover, (2.4) gives $w(0,0)=\theta$, which together with Lemma 2.7, (2.11), and (2.12) ensures (1.3) as well as a bounded width of w. Thus w is a transition front in the sense of Definition 1.1. The claim $w_t>0$ is immediate from (2.2) and the strong maximum principle for w_t . The exponential decay in Remark 1 follows from Lemma 2.5.

3. Uniqueness of Fronts for Ignition Reactions

We will now prove Theorem 1.5(ii), again assuming that the period of q in x_1 is p = 1. Since now $f'_1(0) = 0$, we have automatically $\zeta > f'_1(0)$, and we again assume that $f(\zeta')$ -majorizes g for each $\zeta' > \zeta - \sigma$. All constants in this section will depend on $f(\zeta')$ and $f(\zeta')$ but not on $f(\zeta')$. Without loss we only need to consider fronts moving to the right, which we will denote by $f(\zeta')$.

We can assume $\theta'' \ge \theta_0$ (otherwise we change θ'' to θ_0) and let

$$\varepsilon_0 \equiv \frac{1}{2} \min\{\theta', 1 - \theta''\}$$
 and $\tilde{\theta} \equiv 1 - \frac{\varepsilon_0}{2}$,

thus fixing v from Lemma 2.1.

We let w be an arbitrary front for (1.2) (in particular, we do not assume $w_t > 0$) and u the solution of (1.2) with the fixed initial condition v. Our strategy is as follows. First we will show that w has to decrease exponentially as $x \to \infty$ in the same way as u_n in the last section. Then we will use this to show that u has to converge to some time shift of w in L^{∞} as $t \to \infty$ (thus any two fronts for f must approach each other up to a time shift as $t \to \infty$). Finally, we will show that the rate of this convergence is uniform for all f and all fronts w with uniformly bounded width. This gives a uniform convergence of solutions u_n with initial data $v_n(x) \equiv v(x+n)$ to time shifts of any front w, and the uniqueness of the front follows.

Let w be an arbitrary transition front for (1.2). Consider Ψ from the last section and let

$$X_w(t) \equiv \sup\{x_1 \mid w(t, x) \ge \alpha_f(x) \text{ for some } x = (x_1, x')\},$$

$$Y_w(t) \equiv \inf\{y \mid w(t, x) \le \Psi(x_1 - y, x) \text{ for all } x \in D\},$$

$$Z_{w, \varepsilon}^-(t) \equiv \sup\{y \mid w(t, x) \ge 1 - \varepsilon \text{ when } x_1 \le y\},$$

$$Z_{w, \varepsilon}^+(t) \equiv \inf\{y \mid w(t, x) \le \varepsilon \text{ when } x_1 \ge y\},$$

as well as

$$L_{w,\varepsilon} \equiv \sup_{t \in \mathbb{D}} \{ Z_{w,\varepsilon}^+(t) - Z_{w,\varepsilon}^-(t) \}, \qquad Z_w(t) \equiv Z_{w,\varepsilon_0}^-(t), \quad \text{and} \quad L_w \equiv L_{w,\varepsilon_0}.$$

All of these, except possibly $Y_w(t)$, are finite because w is a transition front, and we have $X_w(t) \in [Z_{w,\varepsilon}^-(t), Z_{w,\varepsilon}^+(t)]$ for $\varepsilon \leq \varepsilon_0 (\leq \min\{\theta_1, 1 - \theta_0\})$. The next lemma shows $Y_w(t) < \infty$.

Lemma 3.1. There is $\tilde{C}_2 < \infty$ (depending on L_w if $w_t \geq 0$) such that for all t we have

$$|Y_w(t) - Z_w(t)| \le \tilde{C}_2. \tag{3.1}$$

Proof. Again $Z_w(t) - Y_w(t) \leq \tilde{C}_2$ (with a uniform bound) is immediate so we are left with proving $Y_w(t) - Z_w(t) \leq \tilde{C}_2$. We fix any $\varepsilon \in (0, \varepsilon_0)$ and for $t \in \mathbb{R}$ define

$$\alpha_{f,\varepsilon}(x) \equiv \inf\{u \in (\varepsilon, 1] \mid f(x, u) \ge \zeta(u - \varepsilon)\} \uparrow \alpha_f(x) \quad \text{as } \varepsilon \to 0,$$

$$X_{w,\varepsilon}(t) \equiv \sup\{x_1 \mid w(t,x) \ge \alpha_{f,\varepsilon}(x) \text{ for some } x = (x_1,x')\} \downarrow X_w(t) \text{ as } \varepsilon \to 0,$$

$$Y_{w,\varepsilon}(t) \equiv \inf\{y \mid w(t,x) \leq \Psi(x_1 - y, x) + \varepsilon \text{ for all } x \in D\} \uparrow Y_w(t) \text{ as } \varepsilon \to 0.$$

The convergences hold because $\theta' > 0$, w and f are continuous, and \mathbb{T}^{d-1} is compact. Note that $\alpha_{f,\varepsilon}(x) \in [\theta',\theta'']$ because $f(x,\cdot)$ decreases on $[\theta'',1]$. Thus for any t and $\varepsilon \leq \varepsilon_0$,

$$0 \le X_{w,\varepsilon}(t) - Z_w(t) \le L_w \tag{3.2}$$

by $\varepsilon_0 < \theta', 1 - \theta''$. Hence it is sufficient to show

$$Y_{w,\varepsilon}(t) - X_{w,\varepsilon}(t) \le C_2' \tag{3.3}$$

with C_2' independent of ε (then use (3.2) and take $\varepsilon \to 0$ to obtain $Y_w(t) - Z_w(t) \le \tilde{C}_2 \equiv C_2' + L_w$). We will prove (3.3) using the argument from Lemma 2.5. We do not have $w_t > 0$ here but Lemma 2.6(i) will not be needed. Instead, Lemma 2.6(ii) will suffice thanks to (3.2).

Pick any $t_0 \in \mathbb{R}$ and notice that $\Psi(0,x) \geq 1$ implies $Y_{w,\varepsilon}(t_0) \leq Z_{w,\varepsilon}^+(t_0)$. We also have $X_{w,\varepsilon}(t_0) \geq Z_w(t_0) \geq Z_{w,\varepsilon}^-(t_0)$ by (3.2), so

$$Y_{w,\varepsilon}(t_0) - X_{w,\varepsilon}(t_0) \le L_{w,\varepsilon}.$$

As long as $X_{w,\varepsilon}(t) \leq Y_{w,\varepsilon}(t)$ for $t \geq t_0$, the argument in Lemma 2.5 shows that $Y_{w,\varepsilon}(t)$ increases with average speed at most c_{ζ} (i.e., $Y_{w,\varepsilon}(t) \leq Y_{w,\varepsilon}(t_0) + c_{\zeta}(t-t_0)$) because $\psi(t,x) \equiv \Psi(x_1 - Y_{w,\varepsilon}(t_0) - c_{\zeta}(t-t_0), x) + \varepsilon$ solves (1.2) with $\zeta(u-\varepsilon)$ in place of f. On the other hand, Lemma 2.6(ii) means that $Z_w(t)$ (and thus also $X_{w,\varepsilon}(t)$ due to (3.2)) increases with average speed at least $(c_0 + c_{\zeta})/2 > c_{\zeta}$ after an initial time delay $t'_{(c_0-c_{\zeta})/2}$ (independent of t_0). Thus the faster moving $X_{w,\varepsilon}(t)$ will catch up with $Y_{w,\varepsilon}(t)$ and we have $X_{w,\varepsilon}(t_1) \geq Y_{w,\varepsilon}(t_1)$ for some $t_1 \in [t_0, t_0 + t''_{\varepsilon}]$. Here t''_{ε} is independent of t_0 because the speed difference $z_0 = (c_0 - c_{\zeta})/2$ and initial distance $z_0 = (c_0 - c_{\zeta})/2$ and initial distance $z_0 = (c_0 - c_{\zeta})/2$ and $z_0 = (c_0 - c_{\zeta})/2$ independent of $z_0 = (c_0 - c_{\zeta})/2$ independent of $z_0 = (c_0 - c_{\zeta})/2$ and increases with average speed at most $z_0 = (c_0 - c_{\zeta})/2$ and increases with average speed at most $z_0 = (c_0 - c_{\zeta})/2$ and Lemma 2.6(ii). This proves the existence

of C_2' (depending on L_w but independent of ε) such that (3.3) holds for all $t \geq t_0 + t_\varepsilon''$. Since t_0 has been arbitrary, (3.3) holds for all t and all $\varepsilon \in (0, \varepsilon_0)$, and taking $\varepsilon \to 0$ gives (3.1). In particular, $Y_w(t)$ is finite, and as in Lemma 2.4 we obtain for $t \geq \tau$,

$$Y_w(t) - Y_w(\tau) \le c_{\xi}(t - \tau). \tag{3.4}$$

Finally, note that if $w_t \geq 0$, then Lemma 2.6(i) applies to w. As a result, the proof of (3.3) for small enough ε is identical to that of (2.12) (using that $f(\zeta - \sigma/2)$ -majorizes g). Hence we do not need (3.2) to show that $X_{w,\varepsilon}(t)$ increases with average speed larger than $(c_0 + c_\zeta)/2$, and C'_2 becomes L_w -independent. Then $\varepsilon \to 0$ gives $Y_w(t) - X_w(t) \leq C'_2$ and thus

$$|Y_w(t) - X_w(t)| \le C_2' \tag{3.5}$$

This and L_w -independent upper bounds on $X_w(t - t_{\varepsilon_0}) - Z_w(t)$ (from Lemma 2.6(i)) and on $X_w(t) - X_w(t - t_{\varepsilon_0})$ (from (3.5) and (3.4)) show that \tilde{C}_2 in (3.1) is also L_w -independent. \square

Remark. Notice that this result, together with the definition of Y_w and (3.8) below, shows that $L_{w,\varepsilon}$ depends only on L_w and ε .

Let u be the solution of (1.2) with initial condition v(x) from Lemma 2.1 (with fixed $\tilde{\theta} \equiv 1 - \varepsilon_0$). Let us define

$$X_{u}(t) \equiv \sup\{x_{1} \mid u(t, x) \geq \alpha_{f}(x) \text{ for some } x = (x_{1}, x')\},$$

$$Y_{u}(t) \equiv \inf\{y \mid u(t, x) \leq \Psi(x_{1} - y, x) \text{ for all } x \in D\},$$

$$Z_{u,\varepsilon}^{-}(t) \equiv \sup\{y \mid u(t, x) \geq 1 - \varepsilon \text{ when } x_{1} \leq y\},$$

$$Z_{u,\varepsilon}^{+}(t) \equiv \inf\{y \mid u(t, x) \leq \varepsilon \text{ when } x_{1} \geq y\},$$

as well as

$$L_{u,\varepsilon} \equiv \sup_{t > t'} \{ Z_{u,\varepsilon}^+(t) - Z_{u,\varepsilon}^-(t) \}, \qquad Z_u(t) \equiv Z_{u,\varepsilon_0}^-(t) \quad \text{and} \quad L_u \equiv L_{u,\varepsilon_0},$$

with t'_{ε} from Lemma 2.6(ii). All these are finite as in Section 2 and again

$$Y_u(t) - Y_u(\tau) \le c_{\xi}(t - \tau). \tag{3.6}$$

Lemmas 2.5–2.7 and $u_t > 0$ again show

$$|Y_u(t) - Z_u(t)| \le C_2 \tag{3.7}$$

for some f-independent C_2 . We also have that for each $\varepsilon > 0$ there are $C_{\varepsilon}, C_{\varepsilon} < \infty$ (the latter L_w -dependent if $w_t \not\geq 0$) such that for any $t \geq \tau$,

$$Z_{w,\varepsilon}^{-}(t) \ge Z_w(\tau) + \frac{c_0 + c_{\zeta}}{2}(t - \tau) - \tilde{C}_{\varepsilon}, \tag{3.8}$$

$$Z_{u,\varepsilon}^{-}(t) \ge Z_u(\tau) + \frac{c_0 + c_{\zeta}}{2}(t - \tau) - C_{\varepsilon} \quad (\tau \ge t_{\varepsilon}' \text{ if } \varepsilon \le 1 - \tilde{\theta} = \frac{\varepsilon_0}{2}). \tag{3.9}$$

Here (3.9) holds because Lemma 2.6(ii) shows that

$$Z_{u,\varepsilon}^{-}(t) \ge Z_u(\tau - t_{\varepsilon}') + \frac{c_0 + c_{\zeta}}{2}(t - \tau)$$

and $Z_u(\tau) - Z_u(\tau - t'_{\varepsilon})$ is uniformly bounded in τ due to (3.6) and (3.7). The same argument works for w but it uses (3.1) and so \tilde{C}_{ε} depends on L_w via \tilde{C}_2 (unless $w_t \geq 0$).

Before we can show that u converges to some time shift of w, we need to prove that once u is close to a time shift of w, it will not depart far from it.

Lemma 3.2. For each $\varepsilon > 0$ there is a $\delta > 0$ (depending also on L_w if $w_t \geq 0$) such that:

- (i) If $w(t_1, x) \leq u(t_0, x) + \delta$ for some $t_0 \geq 1$, $t_1 \in \mathbb{R}$ and all $x \in D$, then $w(t + t_1 t_0, x) \leq u(t, x) + \varepsilon$ for all $t \geq t_0$ and $x \in D$.
- (ii) If $w(t_1, x) \ge u(t_0, x) \delta$ for some $t_0 \ge 1$, $t_1 \in \mathbb{R}$ and all $x \in D$, then $w(t + t_1 t_0, x) \ge u(t, x) \varepsilon$ for all $t \ge t_0$ and $x \in D$.

Proof. This is proved via a construction of suitable supersolution and subsolution. To do that, let $\kappa(\lambda_{\zeta}/2) \geq 0$ and $\gamma(x; \lambda_{\zeta}/2) > 0$ be from (2.10) with $\lambda = \lambda_{\zeta}/2$. If we continue $\gamma(x; \lambda_{\zeta}/2)$ periodically on D and let

$$\Phi(s,x) \equiv [\inf_{D} \gamma(x;\lambda_{\zeta}) \inf_{D} \gamma(x;\lambda_{\zeta}/2)]^{-1} e^{-\lambda_{\zeta}s/2} \gamma(x;\lambda_{\zeta}/2) > 0,$$

then (recall that $\sup_{D} \gamma(x; \lambda_{\zeta}) = 1$)

$$\Phi(s,x) \ge e^{\lambda_{\zeta} s/2} \Psi(s,x) \tag{3.10}$$

for $(s,x) \in \mathbb{R} \times D$. We also have from the convexity of $\kappa(\lambda)$ [1, Proposition 5.7(iii)] and $\kappa(0) = 0$ that

$$\kappa\left(\frac{\lambda_{\zeta}}{2}\right) \le \frac{\kappa(\lambda_{\zeta})}{2} < \frac{c_{\zeta}\lambda_{\zeta}}{2}.$$

Thus for each $y \in \mathbb{R}$ the function $\phi(t, x) \equiv \Phi(x_1 - y - c_{\zeta}t, x)$ satisfies

$$\phi_t + q \cdot \nabla \phi - \operatorname{div}(A\nabla \phi) = \left[\frac{c_\zeta \lambda_\zeta}{2} - \kappa \left(\frac{\lambda_\zeta}{2}\right)\right] \phi \ge 0.$$
 (3.11)

Next pick $\omega \leq 1 \leq \Omega$ so that for each f as in the statement of Theorem 1.5 (with fixed $q, A, f_0, f_1, \zeta, g, K, \theta''$),

$$0 < \omega \le \inf\{u_t(t, x) \mid t \ge 1 \text{ and } x_1 \in [Z_u(t), Z_{u, \varepsilon_0}^+(t)]\},$$
(3.12)

$$\infty > \Omega \ge \sup\{u_t(t, x) \mid t \ge 1 \text{ and } x \in D\}.$$
(3.13)

The existence of such Ω follows from parabolic regularity and boundedness of u. The existence of ω is guaranteed by $u_t>0$ and is proved as follows. Assume the contrary, that is, there are sequences f_n and $(t_n,x^n)\in[1,\infty)\times[Z_u(t_n),Z_{u,\varepsilon_0}^+(t_n)]\times\mathbb{T}^{d-1}$ such that $u_t(t_n,x^n)\to 0$. As at the end of Section 2, the functions $u_n(t,x)\equiv u(t+t_n,x+\lfloor x_1^n\rfloor e_1)$ contain a subsequence which converges in $C_{\text{loc}}^{1;2}$ to a solution \tilde{u} of (1.2) on $(-1,\infty)\times D$ (with the same q,A and some Lipschitz reaction $\tilde{f}(x,u)\in[f_0(u),f_1(u)]$ which is a locally uniform limit of a subsequence of $f_n(x+\lfloor x_1^n\rfloor e_1,u)$). But then $\tilde{u}_t(0,\tilde{x})=0$ for some $\tilde{x}\in\mathbb{T}^d$, and so $\tilde{u}_t\geq 0$ and the strong maximum principle for \tilde{u}_t show $\tilde{u}_t\equiv 0$. Since $Z_{u,\varepsilon_0}^+(t)-Z_u(t)$ is uniformly bounded (in t and t) due to (3.7), we have t lim t sup t such that t is uniformly bounded (in t and t in t such that t is uniformly bounded (in t and t in t in t such that t is uniformly bounded (in t and t in t i

$$Z_{\tilde{u}}(t) \equiv \sup\{y \mid \tilde{u}(t,x) \ge 1 - \varepsilon_0 \text{ when } x_1 \le y\}$$

again grows with a positive average speed after an initial time delay, by Lemma 2.6(ii).

Finally, assume without loss that $\varepsilon \leq \varepsilon_0$ and $t_1 = t_0$ (otherwise we shift w in t by $t_1 - t_0$), and increase K so that

$$K \ge \frac{\lambda_{\zeta}(c_0 - c_{\zeta})}{4}.\tag{3.14}$$

(i) Notice that (3.7) and the hypothesis show that (after possibly increasing C_2 by a constant only depending on Ψ and thus not on f, u, w)

$$Z_u(t_0) \ge Z_w(t_0) - C_2. \tag{3.15}$$

Let

$$C'_{\varepsilon} \equiv C_2 + \tilde{C}_2 + C_{\varepsilon} + \tilde{C}_{\varepsilon} + 1,$$

$$b_{\varepsilon} \equiv \frac{\varepsilon \lambda_{\zeta} (c_0 - c_{\zeta}) e^{-\lambda_{\zeta} C'_{\varepsilon}/2} \omega}{4\Omega K \sup_{D} \Phi(0, x)},$$

$$\beta(t) \equiv \frac{\varepsilon}{\Omega} \left(1 - e^{-\lambda_{\zeta} (c_0 - c_{\zeta})(t - t_0)/4} \right),$$

$$\phi_{+}(t, x) \equiv b_{\varepsilon} \Phi(x_1 - Y_w(t_0) - c_{\zeta}(t - t_0), x).$$
(3.16)

Then $b_{\varepsilon} \leq \varepsilon$ by (3.14), $\omega \leq \Omega$, and $\Phi(0,x) \geq 1$. Also, (3.11) holds for ϕ_+ and we define

$$z_{+}(t,x) \equiv \tilde{u}_{+}(t,x) + \phi_{+}(t,x) \equiv u(t+\beta(t),x) + \phi_{+}(t,x)$$

for $t \geq t_0$. Our aim is to show

$$z_{+}(t,x) \ge w(t,x) \tag{3.18}$$

for all $x \in D$ and $t \ge t_0$. This estimate might not appear very useful because ϕ_+ is unbounded but it will suffice. The reason is that at $t = t_0$, the function ϕ_+ is large only where both u, ware close to 1 and therefore also to each other. This setup will persist for all $t \ge t_0$ because ϕ_+ travels with speed c_{ζ} which is strictly smaller than the speeds of propagation of u and w. Thus, in fact, ϕ_+ decays near the reaction zones of u, w as t grows.

Let $0 < \delta \le b_{\varepsilon} \inf_{D} \Phi(2|\ln b_{\varepsilon}|/\lambda_{\zeta}, x)$. Then

$$z_{+}(t_0, x) \ge w(t_0, x) \tag{3.19}$$

for $x_1 \leq Y_w(t_0) + 2|\ln b_{\varepsilon}|/\lambda_{\zeta}$ by the hypothesis $w(t_0, x) \leq u(t_0, x) + \delta$ and our choice of δ , and for $x_1 \geq Y_w(t_0) + 2|\ln b_{\varepsilon}|/\lambda_{\zeta}$ by the definition of $Y_w(t_0)$ and by $b_{\varepsilon}\Phi(2|\ln b_{\varepsilon}|/\lambda_{\zeta} + y, x) \geq \Psi(2|\ln b_{\varepsilon}|/\lambda_{\zeta} + y, x)$ when $y \geq 0$ (see (3.10)).

Moreover, we will prove that z_+ is a supersolution of (1.2) for $t \ge t_0$, with f(x, u) = 0 when $u \ge 1$. Since (3.11) gives

$$(z_+)_t + q \cdot \nabla z_+ - \operatorname{div}(A\nabla z_+) \ge f(x, z_+) + [f(x, \tilde{u}_+) - f(x, z_+) + \beta'(t)u_t(t + \beta(t), x)],$$

this will be established if we show that the square bracket is non-negative.

This is clearly true for $x_1 \leq Z_u(t+\beta(t))$ since then $\tilde{u}_+(t,x) \geq \theta''$ and so $f(x,\tilde{u}_+) \geq f(x,z_+)$. Next, (3.9), (3.15), (3.1), and $\beta(t) \geq 0$ for $t \geq t_0$ give

$$Z_{u,\varepsilon}^{-}(t+\beta(t)) - Y_w(t_0) - c_{\zeta}(t-t_0) \ge \frac{c_0 - c_{\zeta}}{2}(t-t_0) - C_{\varepsilon}', \tag{3.20}$$

so for $x_1 \geq Z_{u,\varepsilon}^-(t+\beta(t))$ we find using (3.14),

$$\phi_{+}(t,x) \le b_{\varepsilon} e^{-\lambda_{\zeta}[(c_{0}-c_{\zeta})(t-t_{0})/2 - C_{\varepsilon}']/2} \sup_{D} \Phi(0,x) = \frac{\beta'(t)\omega}{K} \le \varepsilon \frac{\lambda_{\zeta}(c_{0}-c_{\zeta})}{4K} \le \varepsilon.$$
 (3.21)

This gives for $x_1 \geq Z_{u,\varepsilon}^-(t+\beta(t))$.

$$|f(x, \tilde{u}_+) - f(x, z_+)| \le K\phi_+ \le \beta'(t)\omega.$$

Thus the square bracket is again non-negative for $Z_u(t + \beta(t)) \leq x_1 \leq Z_{u,\varepsilon_0}^+(t + \beta(t))$ due to (3.12) and $Z_u(t + \beta(t)) \geq Z_{u,\varepsilon}^-(t + \beta(t))$. The same is true for $x_1 \geq Z_{u,\varepsilon_0}^+(t + \beta(t))$ because then (3.21) implies $z_+(t,x) \leq \varepsilon_0 + \varepsilon \leq 2\varepsilon_0 \leq \theta'$, yielding $f(x,z_+) = 0$.

Hence z_+ is a supersolution of (1.2) with (3.19), meaning that (3.18) holds. Thus for $x_1 \geq Z_{u,\varepsilon}^-(t+\beta(t))$,

$$u(t,x) - w(t,x) \ge z_+(t,x) - w(t,x) - \phi_+(t,x) - \beta(t)\Omega \ge 0 - \varepsilon - \varepsilon = -2\varepsilon$$

using (3.21) and (3.16), and for $x_1 \leq Z_{u,\varepsilon}^-(t + \beta(t))$,

$$u(t,x) - w(t,x) \ge \tilde{u}_+(t,x) - \beta(t)\Omega - 1 \ge \tilde{u}_+(t,x) - \varepsilon - 1 \ge -2\varepsilon.$$

This proves (i) with 2ε in place of ε . Note that δ also depends on C'_{ε} and thus on L_w when $w_t \not \geq 0$.

(ii) Recall that we assume $\varepsilon \leq \varepsilon_0$, $t_1 = t_0$, and (3.14), and let us also assume $\varepsilon \leq c_0^{-1}$. This time (3.1) and the hypothesis give (after increasing \tilde{C}_2 by an f, u, w-independent constant)

$$Z_w(t_0) \ge Z_u(t_0) - \tilde{C}_2.$$
 (3.22)

Then the proof goes along the same lines as in (i) but using

$$\phi_{-}(t,x) \equiv b_{\varepsilon} \Phi(x_{1} - Y_{u}(t_{0}) - c_{\zeta}(t - t_{0}), x),$$

$$z_{-}(t,x) \equiv \tilde{u}_{-}(t,x) - \phi_{-}(t,x) \equiv u(t - \beta(t), x) - \phi_{-}(t,x).$$
(3.23)

This time

$$(z_{-})_{t} + q \cdot \nabla z_{-} - \operatorname{div}(A\nabla z_{-}) \leq f(x, z_{-}) - [f(x, z_{-}) - f(x, \tilde{u}_{-}) + \beta'(t)u_{t}(t - \beta(t), x)].$$

with $f(x, u) = 0$ for $u \leq 0$, and

$$z_{-}(t_0, x) \le w(t_0, x) \tag{3.24}$$

if δ is as in (i). We again need to show that the square bracket is non-negative.

For $x_1 \geq Z_{u,\varepsilon_0}^+(t-\beta(t))$ we have $\tilde{u}_-(t,x) \leq \varepsilon_0$, so $f(x,\tilde{u}_-)=0$ and the square bracket is non-negative. For $Z_u(t-\beta(t)) \leq x_1 \leq Z_{u,\varepsilon_0}^+(t-\beta(t))$ the same is true because

$$Z_{u,\varepsilon}^{-}(t-\beta(t)) - Y_u(t_0) - c_{\zeta}(t-t_0) \ge \frac{c_0 - c_{\zeta}}{2}(t-t_0) - C_{\varepsilon}'$$
(3.25)

(from (3.9), (3.7), and $\beta(t) \le \varepsilon \le c_0^{-1} \le 2/(c_0 + c_\zeta)$) again gives for $x_1 \ge Z_{u,\varepsilon}(t - \beta(t))$,

$$\phi_{-}(t,x) \leq b_{\varepsilon} e^{-\lambda_{\zeta}[(c_{0}-c_{\zeta})(t-t_{0})/2 - C_{\varepsilon}']/2} \sup_{D} \Phi(0,x) = \frac{\beta'(t)\omega}{K} \leq \varepsilon \frac{\lambda_{\zeta}(c_{0}-c_{\zeta})}{4K} \leq \varepsilon, \qquad (3.26)$$
$$|f(x,\tilde{u}_{-}) - f(x,z_{-})| \leq K\phi_{-} \leq \beta'(t)\omega.$$

For $x_1 \leq Z_u(t - \beta(t))$ we have $\tilde{u}_-(t, x) \geq 1 - \varepsilon_0$, so the bracket is non-negative as long as $\phi_-(t, x) \leq \varepsilon_0$ (because then $1 - \theta'' \leq z_-(t, x) \leq \tilde{u}_-(t, x)$). This means that z_- is a subsolution of (1.2) where $\phi_-(t, x) \leq \varepsilon_0$.

Since (3.8), (3.22), and (3.7) imply

$$Z_{w,\varepsilon}^{-}(t) - Y_u(t_0) - c_{\zeta}(t - t_0) \ge \frac{c_0 - c_{\zeta}}{2}(t - t_0) - C_{\varepsilon}'$$

(3.26) also holds for $x_1 \geq Z_{w,\varepsilon}^-(t)$. Thus z_- is a subsolution of (1.2) on the set where $\phi_-(t,x) \leq \varepsilon_0$ while on the complement of that set we have $x_1 \leq Z_{w,\varepsilon}^-(t)$ and so

$$w(t,x) \ge 1 - \varepsilon \ge 1 - \varepsilon_0 \ge 1 - \phi_-(t,x) \ge z_-(t,x).$$

This together with (3.24) gives $z_{-}(t,x) \leq w(t,x)$ for $t \geq t_0$ and $x \in D$. The rest of the proof is analogous to (i), with $Z_{w,\varepsilon}^{-}(t)$ in place of $Z_{u,\varepsilon}^{-}(t+\beta(t))$.

Lemma 3.3. If

$$\tau_w \equiv \inf\{\tau \mid \liminf_{t \to \infty} \inf_{x \in D} [w(t+\tau, x) - u(t, x)] \ge 0\},\tag{3.27}$$

then $-\infty < \tau_w < \infty$. Moreover, the infimum is also a minimum and so

$$\liminf_{t \to \infty} \inf_{x \in D} \left[w(t + \tau_w, x) - u(t, x) \right] \ge 0.$$
(3.28)

Proof. The set in (3.27) is an interval (a, ∞) for some $a \leq \infty$ due to $u_t > 0$. Inequality (3.8) shows $\lim_{t\to\infty} Z_w(t) = \infty$, so $f_0 > 0$ on $[1-\varepsilon_0,1)$ yields $w(\tau,x) \geq \tilde{\theta}\chi_{(-\infty,0]}(x_1)$ for some $\tau < \infty$ and all $x \in D$. Thus $w(\tau,x) \geq u(0,x)$ for all $x \in D$, and the comparison principle shows $w(t+\tau,x) \geq u(t,x)$ for all $t \geq 0$ and $x \in D$. Hence $\tau_w < \infty$.

In the opposite direction, notice that (3.9) gives $\lim_{t\to\infty} Z_{u,\delta}^-(t) = \infty$. Hence for each $\delta > 0$ the hypothesis of Lemma 3.2(i) is satisfied with $t_1 = 0$ and a large t_0 . Then Lemma 3.2(i) and (3.12) prove for all $t \geq t_0$ that $\inf_{x \in D} [w(t-t_0,x)-u(t+2\varepsilon/\omega,x)] \leq -\varepsilon$, provided we choose $\varepsilon > 0$ small enough and then δ, t_0 according to Lemma 3.2(i). Thus $\tau_w > -t_0 - 2\varepsilon/\omega > -\infty$. Hence τ_w is finite and then the infimum must be a minimum by (3.13).

Lemma 3.4. We have

$$\lim_{t \to \infty} ||w(t + \tau_w, x) - u(t, x)||_{L_x^{\infty}} = 0.$$
(3.29)

Proof. We can assume without loss of generality that $\tau_w = 0$ (otherwise we shift w in t). Then (3.28) reads

$$\liminf_{t \to \infty} \inf_{x \in D} [w(t, x) - u(t, x)] \ge 0$$
(3.30)

and we are left with proving

$$\limsup_{t \to \infty} \sup_{x \in D} [w(t, x) - u(t, x)] \le 0.$$
(3.31)

Assume this is not true. Then by Lemma 3.2(i), there is $\delta_0 > 0$ such that for all $t \ge 1$,

$$\sup_{x \in D} [w(t, x) - u(t, x)] \ge \delta_0. \tag{3.32}$$

Moreover, the definition of $\tau_w = 0$ and Lemma 3.2(ii) show that for each $\tau > 0$ there is $\delta_{\tau} > 0$ such that for all $t \geq 1$,

$$\inf_{x \in D} [w(t - \tau, x) - u(t, x)] \le -\delta_{\tau}. \tag{3.33}$$

Finally, we claim that $Z_w(t) - Z_u(t)$ stays bounded as $t \to \infty$. The lower bound follows from (3.30) and $L_u, L_w < \infty$. The upper bound follows from (3.33) for $\tau = 1, L_{u,\delta_1}, L_{w,\delta_1} < \infty$, and a uniform upper bound on $Z_w(t) - Z_w(t-1)$ (due to (3.1) and (3.4)).

As before, there is a sequence $t_n \to \infty$ such that the functions $w(t+t_n, x+\lfloor Z_w(t_n)\rfloor e_1)$ and $u(t+t_n,x+\lfloor Z_w(t_n)\rfloor e_1)$ converge in $C^{1,2}_{loc}(\mathbb{R}\times D)$ to two solutions \tilde{w},\tilde{u} of (1.2) with some reaction \tilde{f} which has all the properties of f. Moreover, \tilde{w}, \tilde{u} are both transition fronts because of the boundedness of $Z_w(t) - Z_u(t)$ and the properties of w, u (namely, (3.1), (3.4), (3.8), (3.6), (3.7), and (3.9). We also have $\tilde{u}_t \geq 0$ as well as

$$\tilde{w}(t,x) \ge \tilde{u}(t,x) \quad \text{for all } (t,x) \in \mathbb{R} \times D,$$
 (3.34)

$$\sup_{x \in D} [\tilde{w}(t, x) - \tilde{u}(t, x)] \ge \delta_0 \quad \text{for all } t \in \mathbb{R}, \tag{3.35}$$

$$\sup_{x \in D} [\tilde{w}(t, x) - \tilde{u}(t, x)] \ge \delta_0 \quad \text{for all } t \in \mathbb{R},$$
$$\inf_{x \in D} [\tilde{w}(t - \tau, x) - \tilde{u}(t, x)] \le -\delta_{\tau} \quad \text{for all } t \in \mathbb{R}, \, \tau > 0.$$

This is thanks to (3.30), (3.32), (3.33), $t_n \to \infty$, and the uniform boundedness in t of $\max\{Z_{w,\varepsilon}^+(t),Z_{u,\varepsilon}^+(t)\} - \min\{Z_{w,\varepsilon}^-(t),Z_{u,\varepsilon}^-(t)\} \text{ (for any } \varepsilon>0).$

We define $Z_{\tilde{w},\varepsilon}^{\pm}(t)$, $Z_{\tilde{u},\varepsilon}^{\pm}(t)$, $L_{\tilde{w},\varepsilon}$, $L_{\tilde{u},\varepsilon}$ analogously to $Z_{w,\varepsilon}^{\pm}(t)$, $L_{w,\varepsilon}$. Then $Z_{\tilde{w},\varepsilon}^{\pm}(t) \geq Z_{\tilde{u},\varepsilon}^{\pm}(t)$ for any $\varepsilon > 0$ by (3.34), and $Z_{\tilde{w},\varepsilon}^{+}(t) - Z_{\tilde{u},\varepsilon}^{-}(t)$ is uniformly bounded in t because $Z_w(t) - Z_u(t)$ stays bounded as $t \to \infty$. We let $Z^+(t) \equiv \max\{Z^+_{\tilde{u},\varepsilon_0}(t), Z^+_{\tilde{u},\varepsilon_0}(t+1)\}$ and $Z^-(t) \equiv Z^-_{\tilde{u},\varepsilon_0}(t)$ so that $Z^+(t) - Z^-(t)$ is also uniformly bounded in t. Inequality (3.35) shows that for each t there is $x^t \in [Z^-_{\tilde{u},\delta_0}(t),Z^+_{\tilde{w},\delta_0}(t)] \times \mathbb{T}^{d-1}$ such that

$$\tilde{w}(t, x^t) - \tilde{u}(t, x^t) \ge \delta_0.$$

Then (3.34) and Harnack inequality give the existence of $\delta' > 0$ (which is t-independent by $\sup_{\mathbb{R}} (Z^+(t) - Z^-(t)) < \infty$) such that

$$\tilde{w}(t,x) - \tilde{u}(t,x) \ge \delta'$$
 whenever $x_1 \in [Z^+(t), Z^-(t)],$

and so (3.13) yields the existence of $\tau \in (0,1)$ such that

$$\tilde{w}(t,x) \ge \tilde{u}(t+\tau,x)$$
 whenever $x_1 \in [Z^+(t), Z^-(t)]$.

We finish the proof with an argument similar to [13]. We define $z(t,x) \equiv \tilde{w}(t,x) - \tilde{u}(t+1)$ $(\tau,x) \in C^{1,\eta;2,\eta}(\mathbb{R} \times D)$ and notice that z then satisfies

$$z_t + q \cdot \nabla z - \operatorname{div}(A\nabla z) = r(t, x)z$$

with $|r(t,x)| \leq K$. We also have

$$z(t,x) \ge 0$$
 when $x_1 \in [Z^+(t), Z^-(t)],$ (3.36)

 $\inf_{x \in D} z(t, x) \le -\delta_{\tau} \quad \text{for each } t \in \mathbb{R},$

$$r(t,x) \le 0$$
 when $x_1 \notin [Z^+(t), Z^-(t)],$ (3.37)

the last inequality holding because \tilde{f} is non-increasing outside $[\varepsilon_0, 1 - \varepsilon_0]$ and $\tilde{u}_t \geq 0$. Moreover, $z(t, x) \to 0$ as $\operatorname{dist}(x_1, [Z^+(t), Z^-(t)]) \to \infty$ uniformly in t because \tilde{w}, \tilde{u} are transition fronts and hence have bounded width. Let (t_n, x^{t_n}) be such that

$$\lim_{n \to \infty} z(t_n, x^{t_n}) = \delta'' \equiv \inf_{(t, x) \in \mathbb{R} \times D} z(t, x) < 0.$$

Notice that we then have a uniform bound on $\operatorname{dist}(x_1^{t_n}, [Z^+(t_n), Z^-(t_n)])$. Again a subsequence of the sequence of functions $z(t+t_n, x+\lfloor x_1^{t_n} \rfloor e_1)$ converges in $C^{1;2}_{\operatorname{loc}}(\mathbb{R} \times D)$ to a function \tilde{z} with

$$\tilde{z}(0, \tilde{x}) = \delta'' = \inf_{(t, x) \in \mathbb{R} \times D} \tilde{z}(t, x) < 0$$

for some $\tilde{x} \in \mathbb{T}^{d-1}$, and satisfying (due to (3.36) and (3.37))

$$\tilde{z}_t + q \cdot \nabla \tilde{z} - \operatorname{div}(A \nabla \tilde{z}) \ge 0$$
 where $\tilde{z}(t, x) \le 0$.

The strong maximum principle then forces $\tilde{z}(t,x) = \delta'' < 0$ for t < 0, a contradiction with uniform boundedness of $\operatorname{dist}(x_1^{t_n}, [Z^+(t_n), Z^-(t_n)])$ and (3.36). This proves (3.31) and we are done.

Our final ingredient is the claim that the convergence in (3.29) is uniform in f and w.

Lemma 3.5. For any C > 0 and fixed $q, A, f_0, f_1, \zeta, g, K, \theta''$, the convergence in Lemma 3.4 is uniform in all f as above and all fronts w with $L_w \leq C$.

Remark. We will see at the end of this section that the hypothesis $L_w \leq C$ is satisfied for some $C < \infty$ and all f, w. Thus the convergence is uniform in all f, w as in Theorem 1.5(ii).

Proof. Assume the contrary. Thus for some $C, \varepsilon > 0$ and each $n \in \mathbb{N}$, there are w_n, u_n as in Lemma 3.4 — solving (1.2) with reactions f_n (which satisfy the hypotheses of Theorem 1.5(ii) with uniform $q, A, f_0, f_1, \zeta, g, K, \theta''$) and with $\tau_{w_n} = 0$ after a translation of w_n in t — such that $L_{w_n} \leq C$ and for some $t_n \to \infty$

$$||w_n(t_n, x) - u_n(t_n, x)||_{L_x^{\infty}} > \varepsilon.$$
(3.38)

We will obtain a contradiction by finding a subsequence of $\{(f_n, w_n, u_n)\}_n$ which converges locally uniformly to (f, w, u) such that $L_w \leq C$ and (3.29) is violated.

By parabolic regularity, for some $\eta > 0$, the w_n are uniformly bounded in $C^{1,\eta;2,\eta}(\mathbb{R} \times D)$ and the u_n in $C^{1,\eta;2,\eta}([a,\infty) \times D)$ (for any a > 0). We can thus choose a subsequence (which we again index by n) such that $f_n \to f$ in $C_{loc}(D)$, $w_n \to w$ in $C_{loc}^{1;2}(\mathbb{R} \times D)$ and $u_n \to u$ in $C_{loc}^{1;2}((0,\infty) \times D)$. Therefore w,u solve (1.2) on $\mathbb{R} \times D$ and $(0,\infty) \times D$, respectively. Also, $u(0,x) = u_n(0,x) = v(x)$ holds because the f_n are uniformly bounded and v is continuous, so $||u_n(t,x)-v(x)||_{L^\infty_x} \to 0$ as $t \downarrow 0$, uniformly in n. We note that the limiting reaction f again satisfies all the hypotheses, including ζ' -majorization of g for $\zeta' > \zeta - \sigma$.

Next we show that w is a front. The $Z_{w_n}(0)$ must be uniformly bounded above because otherwise $w_n(-1,x) \geq v(x)$ for large n and all $x \in D$, meaning that $\tau_{w_n} \leq -1$, a contradiction. Similarly, the $Y_{w_n}(0)$ are uniformly bounded below because of (3.4) for w_n and the fact that $Y_{w_n}(t'_{\delta})$ are uniformly bounded below for each $\delta > 0$ (by the argument in the second

part of the proof of Lemma 3.3 and $\tau_{w_n} = 0$). Then Lemma 3.1 and $L_{w_n} \leq C$ show that $Z_{w_n}(0)$ and $Y_{w_n}(0)$ are uniformly bounded below and above, as are the average growth rates of $Z_{w_n}(t)$ and $Y_{w_n}(t)$ (due to (3.4), (3.8), and (3.1)). It follows from (3.8) and $L_{w_n} \leq C$ that the locally uniform limit w is indeed a transition front with $L_w \leq C$.

Thus Lemma 3.4 applies to this w and we have (3.29) for some τ_w . So for each $\delta > 0$ there is $s_{\delta} \geq t'_{\delta}$ such that $\|w(s_{\delta} + \tau_w, x) - u(s_{\delta}, x)\|_{L_x^{\infty}} < \delta$. Hence for each $M, \delta > 0$ and all large enough n we have $\|w_n(s_{\delta} + \tau_w, x) - u_n(s_{\delta}, x)\|_{L_x^{\infty}(-M,M)} < 2\delta$. Since $Z_{w_n,\delta}^-(s_{\delta} + \tau_w)$, $Y_{w_n}(s_{\delta} + \tau_w)$, $Y_{u_n}(s_{\delta})$ are uniformly bounded in n by the argument above and $s_{\delta} \geq t'_{\delta}$, it follows that, in fact, $\|w_n(s_{\delta} + \tau_w, x) - u_n(s_{\delta}, x)\|_{L_x^{\infty}} < 2\delta$ for all large enough n. Then $\delta > 0$ being arbitrary and Lemma 3.2 show that for each $\varepsilon' > 0$ there are $N_{\varepsilon'}$, $r_{\varepsilon'}$ such that for all $n > N_{\varepsilon'}$ and $t > r_{\varepsilon'}$,

$$||w_n(t+\tau_w,x)-u_n(t,x)||_{L_x^{\infty}}<\varepsilon'.$$
 (3.39)

Now (3.12) and $\varepsilon_0 \leq \frac{1}{4}$ show for these n, t,

$$\|w_n(t+\tau_w,x)-u_n(t+\tau_w,x)\|_{L_x^\infty} \ge \min\left\{\tau_w\omega,\frac{1}{2}\right\}-\varepsilon'.$$

If $\tau_w \neq 0$, then this contradicts $\tau_{w_n} = 0$ and (3.29) when we take ε' small enough. Therefore $\tau_w = 0$. But then after taking $\varepsilon' = \varepsilon$ and $n > N_{\varepsilon}$ such that $t_n > r_{\varepsilon}$, we obtain a contradiction between (3.38) and (3.39) with $t = t_n$.

We can now proceed to prove Theorem 1.5(ii). Let w be a transition front for (1.2) and translate it in t so that $w(0,0) = \theta$. This is possible by (3.8) although such translation may not be unique. Define $f_n(x,u) \equiv f(x-ne_1,u)$ and let u_n solve (1.2) with reaction f_n and $u_n(0,x) \equiv v(x)$. Pick t_n so that $u_n(t_n,ne_1) = \theta$ and consider the front $w_n(t,x) \equiv w(t-t_n,x-ne_1)$ for (1.2) with f_n .

We have $t_n \to \infty$ by (3.4) as well as $L_{w_n} = L_w$ for each n. Then Lemma 3.5 shows that for any $T \in \mathbb{R}$, uniformly in $t \geq T$,

$$||w(t+\tau_{w_n},x)-u_n(t+t_n,x+ne_1)||_{L_x^{\infty}} = ||w_n(t+t_n+\tau_{w_n},x+ne_1)-u_n(t+t_n,x+ne_1)||_{L_x^{\infty}} \to 0$$
(3.40)

as $n \to \infty$. This and $u_n(t_n, ne_1) = \theta$ show $w(\tau_{w_n}, 0) \to \theta$ as $n \to \infty$. Then τ_{w_n} must be bounded in n by $w(0,0) = \theta$, (3.8), and $L_w < \infty$. So there is a subsequence converging to some $\tau \in \mathbb{R}$. It follows from (3.13) and (3.40) that for each $t \in \mathbb{R}$,

$$||w(t+\tau,x)-u_n(t+t_n,x+ne_1)||_{L_x^{\infty}}\to 0$$

along this subsequence.

If now w_1, w_2 are two fronts for the same f, then we can choose the same subsequence for both, which gives the existence of τ_1, τ_2 such that $w_1(t + \tau_1, x) = w_2(t + \tau_2, x)$ for each t, x. Thus the two fronts are time shifts of each other, that is, each front for this f is a time shift of the front w constructed at the end of Section 2.

Since this front satisfies $w_t > 0$, the constants in this section do not depend on L_w . In particular, \tilde{C}_2 in (3.1) does not, which in turn gives a uniform in f bound on L_w . This proves the remark after Lemma 3.5.

4. Stability of Fronts for Ignition Reactions

We will now prove Theorem 1.5(iii). We make the same assumptions as at the beginning of the last section and all constants will again depend on $q, A, f_0, f_1, \zeta, g, K, \theta''$ as well as on Y, μ, ν , but not on f. Let us denote by w_{\pm} the unique right- and left-moving fronts from the last section.

Without loss we will assume $\mu \leq \lambda_{\zeta}/2$ and (3.14). We can also assume a=0, after possibly shifting the domain. It will be notationally convenient to let u be the solution of (1.2) with initial condition v (with $\tilde{\theta} \equiv 1 - \varepsilon_0/2$ as in the last section), and w the solution of (1.2) in question, with initial condition $w(0,x) \equiv w_0(x)$.

Let us prove claim (a) in Theorem 1.5(iii) (i.e., Definition 1.4(a) with f, w_0 -uniform constants). We have

$$w_0(x) \le e^{-\mu(x_1 - Y)},\tag{4.1}$$

$$w_0(x) \ge \tilde{\theta}\chi_{(-\infty,0)}(x_1),\tag{4.2}$$

where we have also assumed that $\nu = \tilde{\theta} - \theta$. This we can do without loss due to the following. Lemma 2.6(ii) in fact holds with $\theta + \nu$ in place of $(1+\theta)/2$ for any $\nu > 0$ (see [21] or Lemma 5.1) but with ν -dependent t'_{ε} and L'. So if (4.2) holds with $\theta + \nu$ in place of $\tilde{\theta}$, then $w(\tau', x) \geq \tilde{\theta}\chi_{(-\infty,0)}(x_1)$ for $\tau' \equiv t'_{1-\tilde{\theta}} + L'/(c_0 - \varepsilon_0)$. We also have that if

$$Y_w(t) \equiv \inf\{y \mid w(t, x) \le \Phi(x_1 - y, x) \text{ for all } x \in D\}$$

$$\tag{4.3}$$

with

$$\Phi(s,x) \equiv 2[\inf_{D} \gamma(x;\lambda_{\zeta}) \inf_{D} \gamma(x;\mu)]^{-1} e^{-\mu s} \gamma(x;\mu) > 0,$$

corresponding to $\lambda = \mu$ in (2.10), then (3.4) holds with $c_{\xi} = (\xi + \kappa(\mu))/\mu$ (and ξ from Lemma 2.4). In particular, $Y_w(t)$ is again finite because (4.1) and

$$\Phi(s,x) \ge 2e^{-\mu s} \tag{4.4}$$

imply $Y_w(0) \leq Y$. Thus $w(\tau', x) \leq e^{-\mu(x_1 - Y - Y')}$ holds with the f, w-independent constant $Y' \equiv c_\xi \tau' + \mu^{-1} \sup_D \log \Phi(0, x)$. So (4.1) and (4.2) are satisfied for $w(\tau', x)$ in place of $w_0(x)$ and Y + Y' in place of Y, with τ', Y' independent of f, w.

We define $X_w, Z_{w,\varepsilon}^{\pm}, L_{w,\varepsilon}, Z_w, L_w$ as before and $Y_w(t)$ by (4.3). The proof of claim (a) in Theorem 1.5(iii) will be essentially identical to the argument in Lemmas 3.2–3.5, after we have established the basic properties (3.1), (3.4), (3.8) for w. We will then show uniform convergence of u to a time shift of w. Since u also uniformly converges to a time shift of w_+ , claim (a) in Theorem 1.5(iii) will thus be proved. The proof of claim (b) at the end of this section will be a slight variation on the same theme.

Lemma 4.1. The estimates (3.1), (3.4), and (3.8) hold with f, u, w-independent constants.

Proof. We have already proved (3.4) and obviously we also have

$$Y_w(t) \ge Z_w(t) - \tilde{C}_2 \tag{4.5}$$

for some f, w-independent \tilde{C}_2 . We next notice that (4.2) gives $w_0(x) \geq v(x)$, thus $w(t,x) \geq u(t,x)$ and so

$$Z_{w,\varepsilon}^{\pm}(t) \ge Z_{u,\varepsilon}^{\pm}(t)$$
 and $Z_w(t) \ge Z_u(t)$ (4.6)

for all $t \geq 0$. It is therefore sufficient to show that there is a f, u, w-independent t_1 such that

$$Y_w(t) \le Y_u(t+t_1) \tag{4.7}$$

for all $t \ge 0$, because then (3.1) and (3.8) (with new f, u, w-independent constants) follow from (4.5), (4.6) and (3.6), (3.7), (3.9).

We prove this by bounding w above by a uniformly bounded time shift of u plus a small perturbation. The argument is very similar to that in the proof of Lemma 3.2(i), with μ in place of $\lambda_{\zeta}/2$. We let

$$b_0 \equiv \frac{\varepsilon_0 \mu(c_0 - c_\zeta)\omega}{2\Omega K \sup_D \Phi(0, x)} \qquad (\leq \varepsilon_0 \text{ due to (3.14) and } \mu \leq \lambda_\zeta/2),$$

$$Y_0 \equiv Y + \frac{|\log b_0|}{\mu},$$

and choose t_0 so that

$$Z_u(t_0) \ge Y_0 + C_{\varepsilon_0} + C_2. \tag{4.8}$$

This can be done uniformly in f, u, w thanks to (3.9) and $Z_u(0)$ depending only on v. Let

$$\beta(t) \equiv \frac{\varepsilon_0}{\Omega} \left(1 - e^{-\mu(c_0 - c_\zeta)(t - t_0)/2} \right),$$

$$\phi(t, x) \equiv b_0 \Phi(x_1 - Y_0 - c_\zeta(t - t_0), x),$$

$$z(t, x) \equiv \tilde{u}(t, x) + \phi(t, x) \equiv u(t + \beta(t), x) + \phi(t, x),$$

so that (4.4) and (4.1) give

$$z(t_0, x) \ge \phi(t_0, x) \ge w_0(x) \tag{4.9}$$

for all $x \in D$. Convexity of $\kappa(\lambda)$, $\kappa(0) = 0$, and $\mu < \lambda_{\zeta}$ yield

$$\kappa(\mu) \le \frac{\mu}{\lambda_{\zeta}} \kappa(\lambda_{\zeta}) < c_{\zeta} \mu,$$

so that again $\phi_t + q \cdot \nabla \phi - \operatorname{div}(A\nabla \phi) \ge 0$.

It then follows, as in the proof of Lemma (3.2)(i), that z is a supersolution of (1.2). The argument is identical, with $\varepsilon = \varepsilon_0$, μ in place of $\lambda_{\zeta}/2$, (3.20) replaced by

$$Z_u(t+\beta(t)) - Y_0 - c_{\zeta}(t-t_0) \ge \frac{c_0 - c_{\zeta}}{2}(t-t_0) + C_2 \tag{4.10}$$

(which is immediate from (4.8) and (3.9)), and (3.21) by

$$\phi(t,x) \le b_0 e^{-\mu[(c_0 - c_\zeta)(t - t_0)/2 + C_2]} \sup_D \Phi(0,x) = \frac{\beta'(t)\omega}{K} e^{-\mu C_2} \le \varepsilon_0 \frac{\mu(c_0 - c_\zeta)}{2K} \le \varepsilon_0 \tag{4.11}$$

for $x_1 \ge Z_u(t + \beta(t))$.

Thus (4.9) yields $z(t+t_0, x) \ge w(t, x)$ for all $t \ge 0$ and $x \in D$. But then $\Psi(s, x) \le \frac{1}{2}\Phi(s, x)$ for $s \ge 0$ and $b_0 \le \frac{1}{2}$ give

$$w(t,x) \le \tilde{u}(t+t_0,x) + \phi(t+t_0,x) \le \frac{1}{2}\Phi(x_1 - Y_u(t+t_0+\beta(t)),x) + \frac{1}{2}\Phi(x_1 - Y_0 - c_{\zeta}(t-t_0),x)$$

for $x_1 \ge Y_u(t+t_0+\beta(t))$. Since $Y_u(t+t_0+\beta(t)) \ge Y_0+c_\zeta(t-t_0)$ (by (4.10) and (3.7)) and $\Phi(s,x) \ge 1$ for $s \le 0$, we obtain $w(t,x) \le \Phi(x_1-Y_u(t+t_0+\beta(t)),x)$ for $t \ge 0$ (and all x_1). This and $\beta(t) \le \varepsilon_0/\Omega$ now yield (4.7) with $t_1 \equiv t_0 + \varepsilon_0/\Omega$.

Lemma 4.2. Lemmas 3.2, 3.3, 3.4 hold for u, w as above, and the convergence in Lemma 3.4 is uniform in all f, u, w as above (with fixed $q, A, f_0, f_1, \zeta, g, K, \theta'', Y, \mu, \nu$).

Proof. This is virtually identical to the proofs of Lemmas 3.2–3.5, with the following adjustments. In Lemma 3.2 one considers only $t_1 \geq 1$ and λ_{ζ} is replaced by μ in the proof. In the proof of Lemma 3.4 we cannot assume $\tau_w = 0$ but this causes no change to the argument. This is also the case in the proof of Lemma 3.5 but uniformity in f of the argument in Lemma 3.3 shows that the τ_w are uniformly bounded (for fixed Y, μ, ν). Thus in the proof of Lemma 3.5 we can simply consider a sequence w_n with convergent τ_{w_n} and the proof is unchanged. The condition $L_w \leq C$ can be omitted because it is automatic from Lemma 4.1. Finally, the limit w obtained in the proof of Lemma 3.5 along a subsequence of $\{w_n\}_n$ is not a front but satisfies (4.1) and (4.2), and thus the just proved Lemma 3.4 holds for w, too. \square

This and f-uniform convergence of u to a time shift of w_+ prove f, w-uniform convergence of w to a time shift of w_+ in L_x^{∞} , claim (a) in Theorem 1.5(iii).

The proof of claim (b) is virtually identical, with a separate treatment of the two reaction zones of w (one on either side of $x_1 = a$) moving right and left. This requires the adjustment of the definition of $Z_{w,\varepsilon}^-(t)$ (for the right-moving reaction zone) to

$$Z_{w,\varepsilon}^-(t) \equiv \sup\{y \ge a \mid w(t,x) \ge 1 - \varepsilon \text{ when } x_1 \in [a,y]\},$$

and a restriction of all the estimates to $x_1 \geq a$. The rest of the proof is unchanged because our subsolution z_- in Lemma 3.2(ii) is in fact negative for $t \geq t_0$ and $x_1 < a$, as long as t_0 is large enough (depending on ε) so that $\phi_-(t,x) \geq 1$ for these t,x_1 (see (3.23)). Thus we still obtain $z_- \leq w$ and ultimately prove f, w-uniform convergence in $L^\infty_x(D^+_a)$ of w to a time shift of w (and hence of w_+), with $D^+_a \equiv [a,\infty) \times \mathbb{T}^{d-1}$. A similar treatment of the left-moving reaction zone of w gives a f, w-uniform convergence in $L^\infty_x(D^-_a)$ of w to a time shift of w_- , with $D^-_a \equiv (-\infty, a] \times \mathbb{T}^{d-1}$. Since w_\pm converge f-uniformly to 1 in $L^\infty_x(D^+_a)$ by Lemma 2.6, the claim follows.

Appendix A. The Spreading Lemma and Transition Fronts for Homogeneous Ignition Reactions

We will now show how one can use our arguments to obtain a proof of Lemma 2.6(ii), which is from [21], without the use of [21]. In fact, we will prove a slightly stronger result. In the course of its proof we will also prove Theorem 1.5 for $f(x, u) = f_0(u)$, showing that

 c_0 in that theorem is well defined. Recall that $c_0, c_0^- > 0$ are the speeds of the unique right-and left-moving fronts for (5.1) below.

Lemma 5.1. Let q, A, f_0 be as in (H1), (H2). Then for each $\nu > 0$ there is $L_{\nu} > 0$ such that for every $\varepsilon > 0$ there is $t'_{\varepsilon} < \infty$ satisfying the following. If $u : [0, \infty) \times D \to [0, 1]$ solves

$$u_t + q(x) \cdot \nabla u = \operatorname{div}(A(x)\nabla u) + f_0(u), \tag{5.1}$$

and $\tilde{x} \in D$ is such that $\inf \{u(0,x) \mid |x_1 - \tilde{x}_1| \leq L_{\nu}\} \geq \theta + \nu$, then for each $t \geq 0$ we have

$$\inf \left\{ u(t + t_{\varepsilon}', x) \mid x_1 - \tilde{x}_1 \in [-c_0^- t, c_0 t] \right\} \ge 1 - \varepsilon. \tag{5.2}$$

Remark. The comparison principle then gives the same for solutions of (1.2) with $f \geq f_0$. This also gives Lemma 2.6(i) with c_0, c_0^- in place of $c_0 - \varepsilon, c_0^- - \varepsilon$. Thus $(c_0 + c_\zeta)/2$ can be replaced by c_0 in (3.8) and (3.9), proving the first claim in Remark 1 after Theorem 1.5.

Proof. Let $f(x,u) \equiv f_0(u)$ and let us reprove Theorem 1.5 without relying on Lemma 2.6 as originally stated. We first let $\tilde{\theta} = \theta + \nu$ and construct a compactly supported initial datum $v(x) \leq \tilde{\theta}$ that satisfies (2.1). This is done as in Lemma 2.1 but cutting off v on both sides. We let \tilde{v}_+ be as \tilde{v} in that lemma and \tilde{v}_- be a C^2 solution of $-\operatorname{div}(A\nabla \tilde{v}_-) + q \cdot \nabla \tilde{v}_- = -q_1 + \operatorname{div}(Ae_1)$ on \mathbb{T}^d , periodically continued to D. Let $v_{\pm,\varepsilon(x)} \equiv \varepsilon(\tilde{v}_{\pm}(x) \mp x_1)$ and $v(x) \equiv \rho(\min\{v_{+,\varepsilon}(x), v_{-,\varepsilon}(x) + 4\})$ for $\varepsilon > 0$. So v is compactly supported and if ε is small, then there is $a \in \mathbb{R}$ and l > 0 such that

$$v(x) = \begin{cases} \tilde{\theta} & x_1 \in [a-l, a+l], \\ \rho(v_{+,\varepsilon}(x)) & x_1 \ge a, \\ \rho(v_{-,\varepsilon}(x)+4) & x_1 \le a \end{cases}$$

(a is such that $v_{\pm,\varepsilon}(x) \approx \pm 2$ when $x_1 \approx a$). We also have $-\operatorname{div}(A\nabla v_{\pm,\varepsilon}) + q \cdot \nabla v_{\pm,\varepsilon} = 0$, so for some distribution $T \geq 0$ supported on the set $D_{\varepsilon} \equiv \{x \in D \mid v_{+,\varepsilon}(x) = 0 \text{ or } v_{-,\varepsilon}(x) = 0\}$ and with \tilde{v} standing for \tilde{v}_{\pm} when $\pm (x_1 - a) \geq 0$,

$$-\operatorname{div}(A\nabla v) + q \cdot \nabla v = -\varepsilon^2 \chi_{D \setminus D_{\varepsilon}} \rho''(\min\{v_{+,\varepsilon}(x), v_{-,\varepsilon}(x) + 4\}) (\nabla \tilde{v} - e_1) \cdot A(\nabla \tilde{v} - e_1) - T,$$
which is again less than or equal to $f_0(v)$ if ε is small enough.

Next let $L_{\nu} < \infty$ be such that v is supported in $[-L_{\nu}, L_{\nu}] \times \mathbb{T}^{d-1}$. Let u solve (5.1) with u(0,x) = v(x), so that Lemma 2.2 holds for u and t > 0. Thus there is $\tau_{\nu} < \infty$ such that

$$u(\tau_{\nu}, x) \ge \tilde{\theta}\chi_{[-L_{\nu}-p, L_{\nu}+p]}(x_1) \ge \max\{v(x - pe_1), v(x), v(x + pe_1)\}.$$

This, (2.3) for u, and the comparison principle then prove Lemma 5.1 with $c_0' \equiv p/\tau_{\nu} > 0$ in place of c_0, c_0^- . This in turn proves Lemma 2.6 with c_0' in place of $c_0 - \varepsilon, c_0^- - \varepsilon$.

We now notice that f_0 being independent of x and positive on $(\theta, 1)$ shows that f_0 ζ majorizes some (ζ -dependent) g as in Theorem 1.6 for each $\zeta > 0$. We therefore choose $\zeta > 0$ small enough (and a corresponding g) so that $c_{\zeta} < c'_0$ in (2.9) and f_0 ζ' -majorizes g for all $\zeta' > \zeta - \sigma$ (with $\sigma \equiv \zeta/2$). This can be done because $\kappa(0) = \kappa'(0) = 0$ [1, Proposition 5.7(iii)].

Now we can perform the rest of the proof of Theorem 1.5 for (5.1) using c'_0 in place of c_0 , c_0^- (and, in particular, $c_{\zeta} < c'_0$ in place of $c_{\zeta} < c_0$). This yields the existence of a unique rightmoving transition front for (5.1) with some speed $c_0 > 0$, and similarly a left-moving one

with speed $c_0^- > 0$, as well as convergence of general solutions to them as in Theorem 1.5(iii). So the solution u above converges in L_x^{∞} to some time shifts of these fronts in the sense of (1.6). This and the comparison principle now proves (5.2).

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