

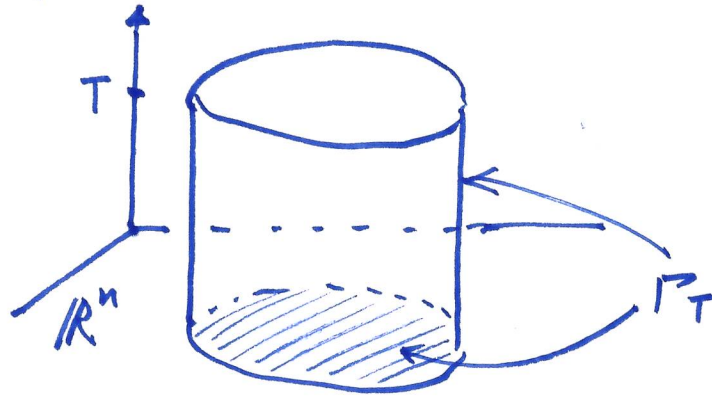
## Section 2.2 Mean-Value Theorem, Maximum Principle, Energy Method

Notation:  $\Omega \subseteq \mathbb{R}^n$ : open, bounded

$T > 0$ .

$$\Omega_T = \{(x, t) : x \in \Omega, t \in [0, T]\} = \Omega \times [0, T]$$

$$\Gamma_T = \Omega \times \{0\} \cup \partial\Omega \times [0, T] = \overline{\Omega_T} \setminus \Omega_T$$



Mean values for harmonic functions are averages over spheres: level sets of fundamental sol'n for Laplace's equation  $(\frac{1}{4\pi|x-y|})$ .

Recall: the fundamental solution for  $u_t - \Delta u = 0$

is

$$K(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{|x|^2}{4t}}$$

So, try to average of solution over level sets of  $K(x, t)$ .

For  $x \in \mathbb{R}^n$ ,  $t \in \mathbb{R}$ ,  $r > 0$ , denote

$$E(x, t; r) = \left\{ (y, s) : y \in \mathbb{R}^n, s \in \mathbb{R}, s \leq t, \right. \\ \left. K(x-y, t-s) \geq \frac{1}{r^n} \right\}$$

Call it a "heat ball".

Notation

$C_1^2(\Omega_T) = \{ u = u(x,t) : \Omega_T \rightarrow \mathbb{R}, u, u_t, \nabla u, \nabla^2 u \in C(\Omega_T) \}$   
 $u \in C_1^2(\Omega_T)$  means  $u = u(x,t)$  is continuous in  $\Omega_T$  and all first-order space and time partial derivatives and all second-order space partial derivatives.

Theorem (Mean Value) Let  $u \in C_1^2(\Omega_T)$  solve the heat equation  $u_t - \Delta u = 0$  in  $\Omega_T$ . Then, for any  $E(x,t;r) \subset \Omega_T$ , we have

$$u(x,t) = \frac{1}{4r^n} \iint_{E(x,t;r)} u(y,s) \frac{|x-y|^2}{(t-s)^2} dy ds.$$

Sketch of Proof

Without loss of generality, set

$(x,t) = (0,0)$ . Denote  $E(r) = E(0,0,r)$ , and

$$\begin{aligned} \phi(r) &= \frac{1}{r^n} \iint_{E(r)} u(y,s) \frac{|y|^2}{s^2} dy ds \\ &\stackrel{y \rightarrow ry, s \rightarrow r^2 s}{=} \iint_{E(1)} u(ry, r^2 s) \frac{|y|^2}{s^2} dy ds \end{aligned}$$

Idea: ① show  $\phi'(r) \equiv 0$ .

②  $\phi(r) = \lim_{t \rightarrow 0} \phi(t) = 4u(0,0)$

$$\begin{aligned} \phi'(r) &= \iint_{E(1)} \left( \sum_{i=1}^n u_{y_i} y_i \frac{|y|^2}{s^2} + 2ry_s \frac{|y|^2}{s} \right) dy ds \\ &= \frac{1}{r^{n+1}} \iint_{E(r)} \left( \sum_{i=1}^n u_{y_i} y_i \frac{|y|^2}{s^2} + 2u_s \frac{|y|^2}{s} \right) dy ds \\ &\equiv A + B \end{aligned}$$

Introduce  $\psi = -\frac{n}{2} \log(-4\pi s) + \frac{|y|^2}{4s} + n \log r$ .  
 $\psi = 0$  on  $\partial E(r)$ . So,

$$B = \frac{1}{r^{n+1}} \iint_{E(r)} 4u_s \sum_{i=1}^n y_i y_i dy ds$$

Integration by parts  
(IBP)

$$= -\frac{1}{r^{n+1}} \iint_{E(r)} \left( 4n u_s r + 4 \sum_{i=1}^n u_{y_i} y_i r \right) dy ds$$

$$\stackrel{\text{IBP m.r.t. } s}{=} \frac{1}{r^{n+1}} \iint_{E(r)} \left( -4n u_s r + 4 \sum_{i=1}^n u_{y_i} y_i r \right) dy ds$$

$$= \frac{1}{r^{n+1}} \iint_{E(r)} \left( -4n u_s r + 4 \sum_{i=1}^n u_{y_i} y_i \left( -\frac{n}{2s} - \frac{|y|^2}{4s^2} \right) \right) dy ds$$

$$= \frac{1}{r^{n+1}} \iint_{E(r)} \left( -4n u_s r - \frac{2n}{s} \sum_{i=1}^n u_{y_i} y_i \right) dy ds - A.$$

$$\phi'(r) = A + B \stackrel{u_s - \Delta u = 0}{=} \frac{1}{r^{n+1}} \iint_{E(r)} \left( -4n u_s r - \frac{2n}{s} \sum_{i=1}^n u_{y_i} y_i \right) dy ds$$

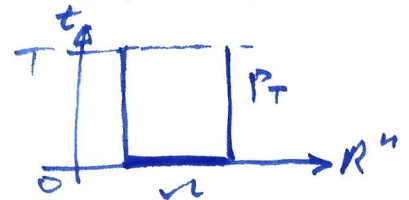
$$\stackrel{\text{IBP}}{\longrightarrow} = \sum_{i=1}^n \iint_{E(r)} \left( 4n u_{y_i} y_i - \frac{2n}{s} u_{y_i} y_i \right) dy ds$$

$$= 0 \quad \text{since } u_{y_i} = \frac{y_i}{2s}$$

$$\text{Now, } \phi(r) = \lim_{t \rightarrow 0} \phi(t) \stackrel{= u(0,0)}{=} \left( \lim_{t \rightarrow 0} \frac{1}{t^n} \iint_{E(t)} \frac{|y|^2}{s} dy ds \right) = u(0,0) 4. \quad \underline{\text{Q.E.D.}}$$

Theorem (Strong Maximum Principle) Assume  $u \in C(\bar{\Omega}_T) \cap C^2(\Omega_T)$  solves the heat equation  $u_t - \Delta u = 0$  in  $\Omega_T$ . Then

$$\max_{\bar{\Omega}_T} u = \max_{\Gamma_T} u$$



Moreover, if  $\Omega$  is connected and  $\exists (x_0, t_0) \in \Omega_T$  s.t.  $u(x_0, t_0) = \max_{\bar{\Omega}_T} u$  then  $u = \text{const}$  in  $\Omega_T$ .

Sketch of Proof Similar to that for Laplace's equation, use the mean-value Theorem.

If  $u(x_0, t_0) = M \equiv \frac{\max \{u\}}{\sqrt{T}}$ , then ~~for~~  $r > 0$  small

$$M = u(x_0, t_0) = \frac{1}{4r^n} \iint_{E(x_0, t_0, r)} u(y, s) \frac{|x_0 - y|^2}{(t_0 - s)^2} dy ds \leq M$$

Since  $1 = \frac{1}{4r^n} \iint_{E(x_0, t_0, r)} \frac{|x_0 - y|^2}{(t_0 - s)^2} dy ds.$

So,  $u \equiv M$  on  $E(x_0, t_0, r)$ .

If  $(y_0, s_0) \in \sqrt{T}$ ,  $s_0 < t_0$ , and  $\exists$  line segment  $L$  connecting  $(y_0, s_0)$  and  $(x_0, t_0)$ , then  $u \equiv M$  on  $L$ . This is because, if we set,

$$r_0 := \inf \{s \geq s_0 : u(x, t) = M \text{ for all } (x, t) \in L, s \leq t \leq t_0\},$$

then  $r_0 = s_0$ .

The last part: use the connectedness of  $\sqrt{T}$ .

Connect  $(x_0, t_0)$  to any  $(y_0, s_0)$ , with  $s_0 < t_0$ ,

by finitely many line segments. Q.E.D.

Thm (Weak Maximum Principle) Let  $u \in C(\sqrt{T}) \cap C_1^2(\Omega_T)$  be such that  $\Delta u \geq u_t$  in  $\sqrt{T}$  ( $\Delta u \leq u_t$  in  $\sqrt{T}$ ). Then  $u$  achieves its maximum on  $\Gamma_T$ . ( $u$  achieves its minimum on  $\Gamma_T$ ).

## Basic idea of Proof

Assume for a moment  $\Delta u > u_t$  in  $\mathcal{V}_T$ .  
 If  $u$  reaches its maximum at some interior point of  $\mathcal{V}_T$ :  $(x, t)$ ,  $x \in \mathcal{V}$ ,  $0 < t < T$ . then  $\partial_t u(x, t) = 0$ ,  $\nabla u(x, t) = 0$  and  $\Delta u(x, t) \leq 0$  (since each  $u_{x_i x_i}(x, t) \leq 0$ ), Contradicting  $\Delta u > u_t$ .  
 Hence  $\max_{\overline{\mathcal{V}_T}} u = \max_{\Gamma_T} u$  for any  $0 < \tau < T$ .  
 Let  $\tau \rightarrow T$ . Then, we can replace  $\tau$  by  $T$ .

Now, just assume  $\Delta u \geq u_t$  in  $\mathcal{V}_T$ . Introduce  $v = u - kt$  for  $k > 0$ . Thus,  $v \leq u$  in  $\overline{\mathcal{V}_T}$ , and  $\Delta v - v_t = \Delta u - u_t + k > 0$ . Hence, by what is proved above,

$$\begin{aligned} \max_{\overline{\mathcal{V}_T}} u &= \max_{\overline{\mathcal{V}_T}} (v + kt) \leq \max_{\overline{\mathcal{V}_T}} v + kT \\ &= \max_{\Gamma_T} v + kT = \max_{\Gamma_T} u + kT. \end{aligned}$$

Let  $k \rightarrow 0$ .

Q.E.D.

Theorem (Uniqueness) There exists at most one solution to 
$$\begin{cases} u_t - \Delta u = f(x, t) & \mathcal{V}(0, T) \\ u(x, t) = g(x, t), & x \in \partial \mathcal{V}, t > 0. \\ u(x, 0) = \phi(x), & x \in \mathcal{V}. \end{cases}$$

Pf. If  $v, w$  both solve the problem, then  $u = v - w$  solves 
$$\begin{cases} u_t - \Delta u = 0 & \mathcal{V}(0, T) \\ u(x, t) = 0 & x \in \partial \mathcal{V}, t > 0 \\ u(x, 0) = 0 & x \in \mathcal{V} \end{cases}$$

So,  $u = 0$  on  $\Gamma_T$ .

Max. Principle  $\Rightarrow u \leq 0$  in  $\mathcal{V}_T$ .

Apply MP to  $-u$ .  $(-u)_t - \Delta(-u) = 0$ ,  $-u = 0$  on  $\Gamma_T$ /  
 $\Rightarrow -u \leq 0$ . So,  $u \equiv 0$  Q.E.D.

Theorem (Maximum principle for the Cauchy problem) Suppose  $u \in C_1^2(\mathbb{R}^n \times (0, T]) \cap C(\mathbb{R}^n \times [0, T])$  solves

$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbb{R}^n \times (0, T) \\ u = g & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases}$$

and satisfies the growth estimate

$$|u(x, t)| \leq A e^{a|x|^2} \quad (x \in \mathbb{R}^n, 0 \leq t \leq T)$$

for some constants  $A, a > 0$ . Then

$$\sup_{\mathbb{R}^n \times [0, T]} u = \sup_{\mathbb{R}^n} g.$$

Note: Without the growth assumption, this can fail.

Theorem (Uniqueness for the Cauchy problem)

Let  $f \in C(\mathbb{R}^n \times [0, T])$  and  $g \in C(\mathbb{R}^n)$ . Then there exists at most one solution  $u \in C_1^2(\mathbb{R}^n \times (0, T]) \cap C(\mathbb{R}^n \times [0, T])$  of the initial-value problem

$$\begin{cases} u_t - \Delta u = f & \text{in } \mathbb{R}^n \times (0, T) \\ u = g & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases}$$

satisfying the growth estimate

$$|u(x, t)| \leq A e^{a|x|^2} \quad (x \in \mathbb{R}^n, 0 \leq t \leq T)$$

for some constants  $A > 0$  and  $a > 0$ .

Note: Without the growth assumption, the solution is not unique!

For the initial-boundary-value problem of the (nonhomogeneous) heat equation, we can prove the solution by the energy method.

Basically, if 
$$\begin{cases} u_t - \Delta u = 0 & \text{in } \Omega \times (0, T] \\ u(x, t) = 0 & \text{on } \partial\Omega, t \in (0, T] \\ u(x, 0) = 0 & \text{in } \Omega. \end{cases}$$

then we want to show  $u \equiv 0$  in  $\Omega$ . Here we assume  $u \in C^1(\bar{\Omega}_T) \cap C_1^2(\Omega_T)$ .

Set 
$$e(t) = \int_{\Omega} u(x, t)^2 dx \quad 0 \leq t \leq T.$$

Then 
$$\begin{aligned} e'(t) &= \int_{\Omega} 2u u_t dx \\ &= 2 \int_{\Omega} u \Delta u dx \\ &= -2 \int_{\Omega} |\nabla u|^2 dx + 2 \int_{\partial\Omega} u \frac{\partial u}{\partial n} dS. \end{aligned}$$

since  $u=0$  on  $\partial\Omega$   
$$= -2 \int_{\Omega} |\nabla u|^2 dx \leq 0$$

So,  $e(t) \leq e(0) = 0$ . So,  $e(t) = 0 \quad \forall t \in [0, T]$ .  
Hence,  $u(x, t) \equiv 0 \quad \forall x \in \Omega, \forall t \in [0, T]$ .

Theorem (Backward Uniqueness) Let  $v, w$  be in  $C_1^2(\Omega_T) \cap C(\bar{\Omega}_T)$ . Suppose 
$$\begin{cases} v_t - \Delta v = f(x, t), & \Omega_T \\ v = g & \text{on } \partial\Omega \times [0, T] \end{cases} \quad \begin{cases} w_t - \Delta w = f(x, t), & \Omega_T \\ w = g & \text{on } \partial\Omega \times [0, T] \end{cases}$$
 Suppose  $v(x, T) = w(x, T) \quad (x \in \Omega)$ . Then  $v \equiv w$  on  $\bar{\Omega}_T$ .

Proof Set  $u = v - w$ . Then  $\begin{cases} u_t - \Delta u = 0 & \text{in } \Omega_T \\ u = 0 & \text{on } \partial \Omega \times [0, T] \\ u(x, T) = 0 & \forall x \in \Omega. \end{cases}$

Let  $e(t) = \int_{\Omega} u(x, t)^2 dx \quad (0 \leq t \leq T)$

Then  $\dot{e}(t) = -2 \int_{\Omega} |\nabla u|^2 dx$   
 ↑ as before

$$\begin{aligned} \ddot{e}(t) &= -4 \int_{\Omega} \nabla u_t \cdot \nabla u_t dx = 4 \int_{\Omega} \Delta u u_t dx \\ &= 4 \int_{\Omega} (\Delta u)^2 dx \end{aligned}$$

Since  $u = 0$  on  $\partial \Omega$ ,

$$\int_{\Omega} |\nabla u|^2 dx = - \int_{\Omega} u \Delta u dx \leq \left( \int_{\Omega} u^2 dx \right)^{1/2} \left( \int_{\Omega} (\Delta u)^2 dx \right)^{1/2}$$

So,  $(\dot{e}(t))^2 = 4 \left( \int_{\Omega} |\nabla u|^2 dx \right)^2$   
 $\leq \int_{\Omega} u^2 dx \cdot 4 \int_{\Omega} (\Delta u)^2 dx$   
 $= e(t) \ddot{e}(t)$   
 i.e.,  $\ddot{e}(t) e(t) \geq (\dot{e}(t))^2$ .

If  $e(t) = 0$  for all  $t \in [0, T]$ , then we are done.

Otherwise,  $\exists [t_1, t_2] \subseteq [0, T]$  with

$$e(t) > 0, \quad t_1 \leq t < t_2, \quad e(t_2) = 0$$

With these, we define  $f(t) \equiv \log e(t)$ ,  $t \in (t_1, t_2)$ .

and get  $f''(t) = \left( \frac{\dot{e}(t)}{e(t)} \right)' = \frac{\ddot{e}e - \dot{e}^2}{e(t)^2} \geq 0$

So,  $f$  is convex on  $(t_1, t_2)$ . Therefore, if  $0 \leq t_1 < t_2 < T$ , then



$$f((1-\tau)t_1 + \tau t) \leq (1-\tau)f(t_1) + \tau f(t).$$

$$\log e((1-\tau)t_1 + \tau t) \leq (1-\tau)\log e(t_1) + \tau\log e(t)$$

$$e((1-\tau)t_1 + \tau t) \leq e(t_1)^{1-\tau} e(t)^\tau$$

$$t \rightarrow t_2: \quad 0 \leq e((1-\tau)t_1 + \tau t_2) \leq e(t_1)^{1-\tau} \underbrace{e(t_2)^\tau}_{=0} = 0$$

$\Rightarrow e \equiv 0$  in  $[t_1, t_2]$ , contradiction.

Q.E.D.

Theorem (Smoothness) Suppose  $u \in C_1^2(\Omega_T)$  solves the heat equation  $u_t - \Delta u = 0$  in  $\Omega_T$ . Then  $u \in C^\infty(\Omega_T)$ .

Q.E.D.