

Introduction

Bo Li, Spring 2019

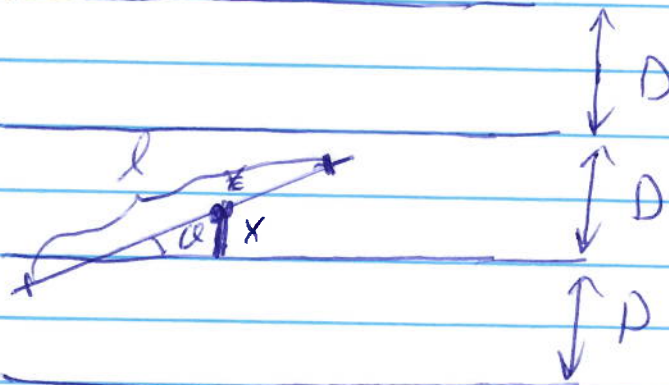
Monte Carlo Methods are a class of computer simulation methods using random numbers. They are used for simulating directly systems that are naturally random or fluctuating. They are also used to simulate artificial random systems to determine or approximate deterministic quantities.

Monte Carlo methods have been ^{extensively} used in physics, chemistry, and biology, and in statistics, operations research, and finance.

(A) Examples

Example 1 (Buffon's needle problem, 1777) This is an example of Monte Carlo (non-computer) experiment, used to determine a deterministic quantity.

Given parallel lines (with separation D) on ground. Throw needles on the ground.



x = dist. of center of needle to line
 l = length of needle
 α = acute angle between needle and line
 $(l \leq D)$

x, θ are both random variables.

x is distributed uniformly. The probability distribution of x is $2/D$ if $0 \leq x \leq D/2$, 0 otherwise. θ is uniformly distributed, with the distribution being $2/\pi$ if $0 \leq \theta \leq \pi/2$, 0 otherwise. The joint distribution of (x, θ) is uniform on $0 \leq x \leq D/2, 0 \leq \theta \leq \pi/2$, i.e., the joint distribution is given by $4/D\pi$ if $0 \leq x \leq D/2, 0 \leq \theta \leq \pi/2$ and 0 otherwise. A needle crosses a line if $x \leq \frac{l}{2} \sin \theta$.

The probability that a needle crosses a line is then

$$p = \int_0^{\pi/2} \int_0^{\frac{l}{2} \sin \theta} \frac{4}{D\pi} dx d\theta = \frac{2l}{\pi D}$$

Hence, $\pi = \frac{2l}{pD}$. On the other hand,

$$p \approx \frac{\# \text{ needles crossing lines}}{\# \text{ needles}}$$

if the $\# \text{ needles} \gg 1$. More precisely, let P_N denote the proportion of "intersects" in N throws. We have

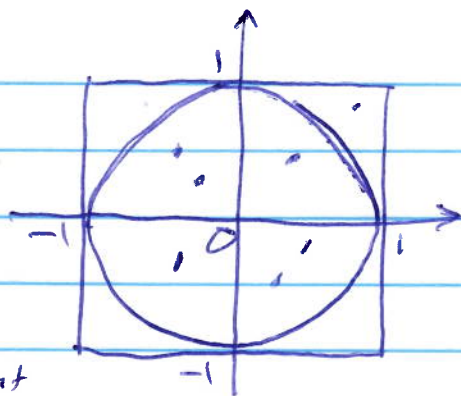
$$\pi = \lim_{N \rightarrow \infty} \frac{2l}{P_N D}$$

Example 2 Computing π again, using a Monte Carlo computer simulation this time.

Area of circle = π

Area of square = 4.

Randomly throw a point into the square. The probability that this point falling into the circle is $\frac{\pi}{4}$.



Suppose we repeat this "experiment" N ($\gg 1$) times. Let I_N be the number of times the point is inside the circle. Then, naturally, $I_N/N \approx \frac{\pi}{4}$.

i.e., $\lim_{N \rightarrow \infty} 4I_N/N = \pi$. Indeed, these N points are uniformly and independently distributed on the square. The random variable I_N has a binomial distribution with the parameters N and $\pi/4$. Therefore

$$E(4I_N/N) = \pi.$$

This means $4I_N/N$ is an unbiased estimator.

This method can be generalised to the calculation of area, volume, etc. using Monte Carlo methods.

Continuing, we can look into the accuracy issue. This is related to the efficiency issue. So, we ask, if we want to have the estimate to have an error within certain range, how large N we need.

Let X_k be the random variable: $X_k = 1$ if the point in the k th "throwing experiment" lies inside the circle, and 0 otherwise. Then $X_1, X_2, \dots, X_N, \dots$ are independent and identically distributed (i.i.d.) with the expectation $E(X_k) = \pi$ ($\forall k \geq 1$). Define

$$\bar{X}_N = \frac{1}{N} (X_1 + \dots + X_N).$$

call it a sample mean. Then the Law of Large Numbers (LLN) implies that $\bar{X}_N \rightarrow E(X_1) = \pi$ with probability 1. Moreover, the Central Limit Theorem (CLT) implies that the sequence

$$\frac{\bar{X}_N - \pi}{\sigma/\sqrt{N}}$$

converges to a random variable that has the standard normal distribution. (Here $\sigma^2 = \text{Var}(X_k)$ ($\forall k \geq 1$) is the variance (equivalently, σ is the standard deviation). Given $\varepsilon > 0$, an error tolerance, we have then for $N \gg 1$ that the probability

$$P\left(-\varepsilon \leq \frac{\bar{X}_N - \pi}{\sigma/\sqrt{N}} \leq \varepsilon\right) \approx \int_{-\varepsilon}^{\varepsilon} g(x) dx$$

where $g(x) = \frac{1}{\sqrt{2\pi}} e^{-x^2/2}$ ($x \in \mathbb{R}$) is the probability density function (PDF) of the standard Gaussian (i.e., normal) random variable. We can rewrite this as

$$P\left(\left|\bar{X}_N - \pi\right| \leq \frac{\varepsilon\sigma}{\sqrt{N}}\right) \approx \int_{-\varepsilon}^{\varepsilon} g(x) dx.$$

Thus, for $N \gg 1$, the probability that π (which is unknown in our "experiment" and analysis) falls into the interval

$$\left[\bar{X}_N - \frac{\varepsilon\sigma}{\sqrt{N}}, \bar{X}_N + \frac{\varepsilon\sigma}{\sqrt{N}}\right]$$

is approximately

$$\gamma := \int_{-\varepsilon}^{\varepsilon} g(x) dx \in (0, 1).$$

This γ -value (percentage) can be numerically estimated for a given $\varepsilon > 0$. We call the interval γ -confidence interval for the estimator \bar{X}_N of π .

Note that, in this case, one has $\sigma^2 = \pi(1-\pi)$. But, π is supposed to be unknown in this analysis. So, we need estimate it. We can estimate $\sqrt{x(1-x)}$ for $x \in (0, 1)$, e.g., $\sigma \leq \sqrt{4 \cdot 4} = 4$.

Example 3 Let D be a bounded region in \mathbb{R}^d and $g: D \rightarrow \mathbb{R}$ a ~~non~~ continuous and bounded function. We consider the integral

$$I = \int_D g(x) dx.$$

If $d \gg 1$ then the usual numerical integration

methods will be prohibited. We should try to use Monte Carlo methods. The starting point of Monte Carlo integration is to reformulate the integral as the expectation of some random variable. Let $f: D \rightarrow \mathbb{R}$ be such that $f \geq 0$ on D . ~~(that $f \geq 0$)~~ and $\int_D f(x) dx = 1$. So f is the probability density function of some random variable $Z \in D$. A simple choice of f is $f(x) = 1/|D|$ ($x \in D$), where $|D|$ is the volume of D (if that is known or easy to calculate). Now, we have

$$I = \int_D \frac{g(x)}{f(x)} f(x) dx = \int_D h(x) f(x) dx,$$

$h(x) = g(x)/f(x)$. Hence

$$I = \mathbb{E}_f(h(Z)),$$

i.e., the expectation of random variable $h(Z)$ with respect to the density f .

The next step is to generate random variables $X_1, X_2, \dots, X_n, \dots$:

(1) All $X_k \in D$.

(2) X_1, \dots, X_n, \dots are independent.

(3) Each X_k is distributed according to f .

Now, by the Law of Large Numbers (LLN)

$$I = E(h(Z)) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{k=1}^N h(X_k)$$

So, for $N \gg 1$, $I \approx \frac{1}{N} \sum_{k=1}^N h(X_k)$.

The error in this estimation is

$$\begin{aligned} \text{error} &= \left| \frac{1}{N} \sum_{k=1}^N h(X_k) - I \right| \\ &= \frac{\sigma_h}{\sqrt{N}} \left| \frac{\frac{1}{N} \sum_{k=1}^N h(X_k) - I}{\sigma_h / \sqrt{N}} \right|, \end{aligned}$$

where $\sigma_h^2 = \text{Var}(h(X_k))$ ($\forall k \geq 1$). We have

$$\begin{aligned} \sigma_h^2 &= \text{Var}(h(Z)) \\ &= \int_D |h(x) - I|^2 f(x) dx \end{aligned}$$

is the variance of $h(Z)$, and hence $h(X_k)$ for all $k \geq 1$. The Central Limit Theorem (CLT) implies that, for $N \gg 1$,

$$\frac{\frac{1}{N} \sum_{k=1}^N h(X_k) - I}{\sigma_h / \sqrt{N}}$$

is distributed approximately according to the standard Gaussian $N(0,1)$. As before, we can determine a γ -confidence interval for estimating I by $\frac{1}{N} \sum_{k=1}^N h(X_k)$. The question is again: how to estimate the variance σ_h^2 as that is unknown.

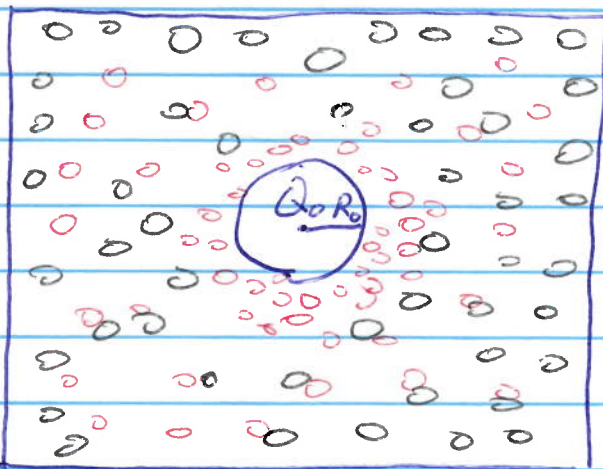
Monte Carlo integration is important as integrals represent the means of random variables. From this example, we see immediately there are two questions:

- (1) How to generate random variables?
- (2) How to reduce the variance? (see, e.g., the term σ_n in the error expression on page [7]).

Example 4 Monte Carlo simulations of ionic systems to investigate the ionic size effect.

○ counter ions

○ co ions



Consider the region $(-L, L)^3$, our simulation box. At the center $(0, 0)$, we place (and fix it) a relatively large colloidal ball of radius R_0 (e.g., $R_0 = 20 \text{ \AA}$), and also place a charge $Q_0 (> 0)$ at the center $(0, 0)$. This is an effective point charge of the

central ball. (The point charge can be equivalently redistributed as the surface charges on the surface of the ball.)

We put many ^{mobile} ions in the box, say, M types (or species) of them.

Each ion of type i carries a point charge $Q_j^{(i)}$ and has the radius $R_j^{(i)}$, $j=1, \dots, N_i$, $i=1, \dots, M$, where N_i is the total number of ions of type i .

The solvent (e.g., water) is treated implicitly through the dielectric coefficient $\epsilon > 0$.

We require the charge neutrality

$$Q_0 + \sum_{i=1}^M \sum_{j=1}^{N_i} Q_j^{(i)} = 0$$

Often $M=2$ or 3 or 4 , or 5 . $Q_j^{(i)}$ can be $+1, +2$, and $-1, -2, -3$. Moreover, N_i can be a few thousands, or more.

Let us relabel all the mobile ions as $1, \dots, N$. So, their positions, radii, and charges are labeled by

$$\vec{r}_k = (x_k, y_k, z_k), R_k, \text{ and } Q_k \quad (1 \leq k \leq N)$$

We set $\vec{r}_0 = (0, 0, 0)$.

~~We use Monte Carlo methods to simulate this system to get equilibrium distributions.~~

All the ions (including the central ball) interact through a potential $U = U(\vec{r}_0, \vec{r}_1, \dots, \vec{r}_N)$, given by

$$A U(\vec{r}_0, \vec{r}_1, \dots, \vec{r}_N) = \sum_{i,j=0, i \neq j}^N U_{ij}(|\vec{r}_i - \vec{r}_j|),$$

where A is a constant that depends on the temperature T , dielectric coefficient ϵ , etc.,

$$U_{ij}(|\vec{r}_i - \vec{r}_j|) = \begin{cases} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|} & \text{if the balls } B(\vec{r}_i, R_i), B(\vec{r}_j, R_j) \\ & \text{do not intersect,} \end{cases}$$

We use the periodical boundary condition. ∞ otherwise.

We use the Metropolis algorithm to simulate this system to get the equilibrium distributions of different types of mobile ions. (Intuitively, counterions will be dense around the central ball; and coions will be mostly away from the central charge.) This algorithm is perhaps the first modern Monte Carlo algorithm, and is also perhaps the most popular Monte Carlo method.

The method generates a sequence of configurations $X_k = (\vec{r}_0, \vec{r}_1^{(k)}, \dots, \vec{r}_N^{(k)})$ (note that $\vec{r}_0 = (0, 0, 0)$ is fixed). The sample means $\bar{X}_N = \frac{1}{N} \sum_{k=1}^N X_k$

provide good estimates of the equilibrium

configuration as $N \gg 1$.

The algorithm

Given $X^{(k)} = (\vec{r}_0, \vec{r}_1, \dots, \vec{r}_N)$.

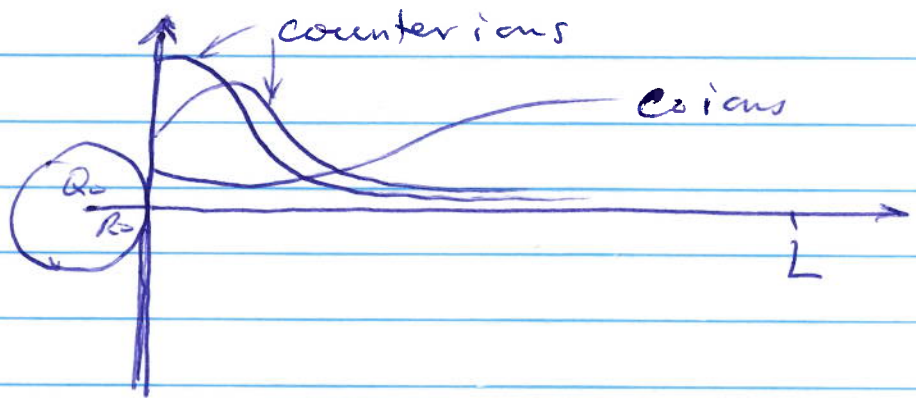
- (1) Choose $J \in \{1, \dots, N\}$ uniformly at random.
- (2) Define $Y = (\vec{r}_0, \vec{r}_1, \dots, \vec{r}_{J-1}, \vec{r}_J^*, \vec{r}_{J+1}, \dots, \vec{r}_N)$ by $\vec{r}_J^* = \vec{r}_J + \vec{\Delta r}$ where $\vec{\Delta r} = (\Delta x, \Delta y, \Delta z)$ with $\Delta x, \Delta y, \Delta z$ independent, and uniform in $(-l_x, l_x), (-l_y, l_y), (-l_z, l_z)$, respectively. Here, l_x, l_y , and l_z are prefixed maximal perturbation sizes, (or more sizes).
- (3) Define $\alpha = \frac{\min(1, -[U(Y) - U(X^{(k)})])}{e} \in (0, 1)$

And accept Y with probability α . i.e., generate a random number U that is uniformly distributed on $[0, 1]$.

If $U \leq \alpha$ then accept Y : $X^{(k+1)} \leftarrow Y$.
 If $U > \alpha$ then reject Y , and $X^{(k+1)} \leftarrow X^{(k)}$.

The sequence $\{X^{(k)}\}_{k=0}^{\infty}$ is a Markov chain. It is reversible, i.e., satisfying the detailed balance. Moreover, it is ergodic so that each configuration will be visited infinitely many times (in principle).

Once the Markov chain is obtained, we need to process the simulation data. For instance we can plot histogram to show the particle density for each type of mobile ions. A typical result (after smoothing out the histogram) is a plot of densities.



(B) Basic Questions / Topics

Questions

- (1) Design estimators for some quantities.
Algorithms.
- (2) Accuracy and efficiency.
- (3) Data processing.

Topics

- ⊙ Generating (or sampling) random variables according to a give distribution, continuous

or discrete. This is a basic step in a Monte Carlo method. Here, one generates a sequence of random variables, e.g., i.i.d. sequence or a Markov chain, according to distribution. There are many methods for generating a random variable. The starting point is to use existing ^{uniform} random number generators. [In these notes, we do not discuss how uniform random numbers can be generated — on computers.]

- ① Variance reduction. While there are many methods to sample random variables. Some of them may give large variances and others small variances. Techniques that can reduce the variance will then be useful.
- ② Markov chain Monte Carlo. A class of methods for generating Markov chains with given targeted distributions. Some of the well-known methods include: The Metropolis method, or a generalized version, the Metropolis-Hastings Method, the Gibbs sampler, the simulated annealing method, etc.

- Statistical analysis of simulation data, including methods/techniques of presenting data, postprocessing data to get better results, etc.

Additional topics include: Monte Carlo integration and summation — which are largely in the part variance reduction; Monte Carlo optimization; rare event simulation; and application in statistical mechanics.

I have collected some basics of probability and Markov chains in the Appendix. Useful analysis tools include: LLN, CLT, and other random variable convergence techniques. Martingales will be also useful sometimes.

It is always good to look into carefully some useful examples/systems to understand the method. These include: A particle moves in a potential; a many-particle system; random walk on a graph; some ^{combinatoric} optimization problems (such as the ^{travelling} salesman problem, the knapsack problem, etc.

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