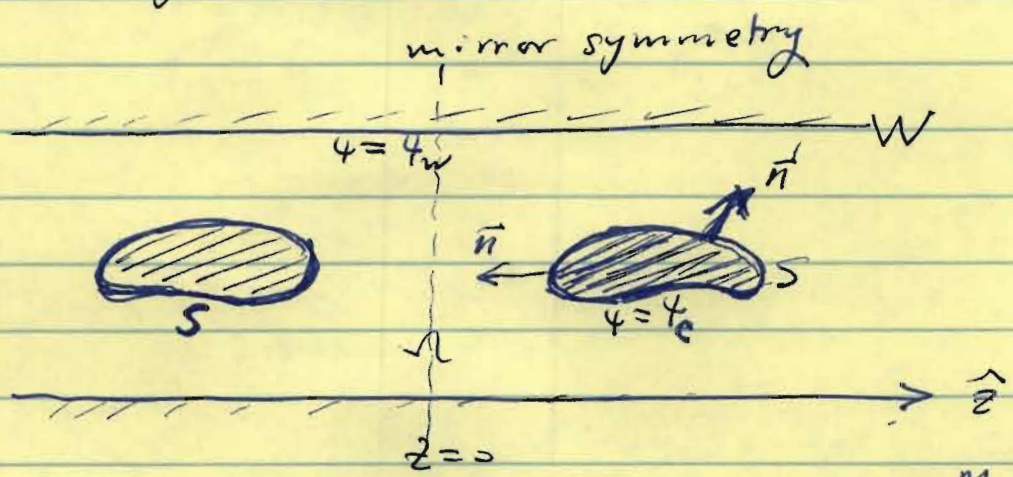


The PB theory does not predict the wall-mediated like-charge attraction



- Assume:
- ⊙ symmetry
 - ⊙ constant body data ψ_w, ψ_c

$$\begin{cases} \Delta\psi - \beta'(\psi) = 0 \\ \psi = \psi_w \text{ on walls } W \\ \psi = \psi_c \text{ on bdry of colloidal bodies } S \\ \psi(\infty) = 0 \end{cases}$$

$$-\beta'(\psi) = \sum_i^M q_i z_i e^{-z_i \psi} \quad (\text{in reduced units})$$

Force acting on a colloidal surface

$$\vec{F} = \frac{1}{2} \int_S \left(\frac{\partial \psi}{\partial n} \right)^2 \vec{n} da$$

$$\begin{aligned} \vec{F} \cdot \hat{z} < 0 & \text{ attraction} \\ \vec{F} \cdot \hat{z} > 0 & \text{ repulsion.} \end{aligned}$$

$$\beta(\psi) = \sum_i q_i e^{-z_i \psi}$$

Let $T = \nabla\psi \otimes \nabla\psi - f I$
 where $f = \frac{1}{2} |\nabla\psi|^2 + \beta(\psi)$.

Ⓘ

$$\vec{F} = \frac{1}{2} \int_S \left(\frac{\partial \psi}{\partial n} \right)^2 \vec{n} da = \int_S T \vec{n} da$$

Verify: $Tn = (\nabla\psi \otimes \nabla\psi)n - (fI)n$

$$= (\nabla\psi \cdot n) \nabla\psi - fn$$

$$= (\nabla\psi \cdot n) \nabla\psi - \frac{1}{2} |\nabla\psi|^2 n - \beta(\psi)n$$

$$\int_S Tnda = \int_S \left[\frac{\partial\psi}{\partial n} \nabla\psi - \frac{1}{2} |\nabla\psi|^2 n - \beta(\psi)n \right] da$$

on S : $\psi = \psi_0 = \text{const.}$ $\beta(\psi) = \beta(\psi_0)$.

$$\int_S n da = \vec{0}. \quad \int_S n_j da = - \int_V \frac{\partial 1}{\partial x_j} dV = 0.$$

$$\nabla\psi = \frac{\partial\psi}{\partial n} n + \frac{\partial\psi}{\partial \tau_1} \tau_1 + \frac{\partial\psi}{\partial \tau_2} \tau_2 \quad (n, \tau_1, \tau_2) \text{ orthonormal}$$

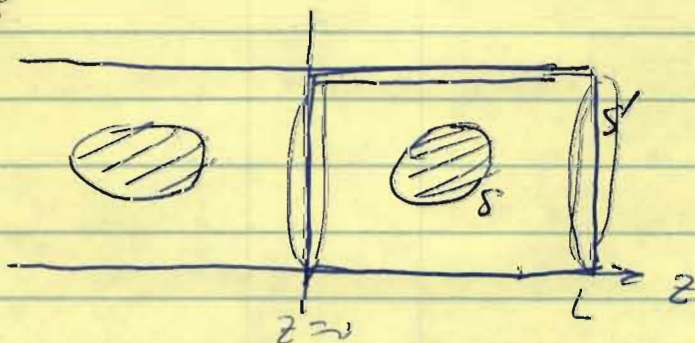
τ_1, τ_2 : tangential direction.

$$\frac{\partial\psi}{\partial \tau_1} = 0 \quad \frac{\partial\psi}{\partial \tau_2} = 0 \text{ along } S. \text{ since } \psi = \psi_0 = \text{const.}$$

$$\int_S Tnda = \int_S \left[\left(\frac{\partial\psi}{\partial n} \right)^2 n - \frac{1}{2} \left(\frac{\partial\psi}{\partial n} \right)^2 n \right] da$$

$$= \int_S \frac{1}{2} \left(\frac{\partial\psi}{\partial n} \right)^2 n da.$$

Let \hat{z} be the unit vector in the positive direction of the z -axis



(II)

$$\hat{z} \cdot \vec{F} = \int_S \hat{z} \cdot Tnda = \int_{S'} \hat{z} \cdot Tnda$$

Since T is divergence free.

~~Pr. 1.4~~

check $T = \nabla\psi \otimes \nabla\psi - \frac{1}{2} |\nabla\psi|^2 I - B(\psi) I$

$$\begin{aligned} \nabla \cdot T &= 2 \Delta\psi \nabla\psi - \Delta\psi \nabla\psi - B'(\psi) \nabla\psi \\ &= [\Delta\psi - B'(\psi)] \nabla\psi = \vec{0} \end{aligned}$$

$$\int_{\text{boundary}} T n da = \int_{\text{volume}} \nabla \cdot T dV = 0$$

boundary = $S + S'$ different directions of n .

(14) Calculate $\hat{z} \cdot \vec{F} = \int_S \hat{z} \cdot T n da$

on the ~~base~~ of the cylinder section, $\psi = \text{const.} = \psi_w$

$$\begin{aligned} \hat{z} \cdot T n &= 0 \\ \hat{z} \cdot T n &= \hat{z} \cdot \left[\nabla \left(\frac{\partial\psi}{\partial n} \nabla\psi \right) - \frac{1}{2} |\nabla\psi|^2 n - B(\psi) n \right] \\ &= \frac{\partial\psi}{\partial n} \cdot \frac{\partial\psi}{\partial z} - \frac{1}{2} |\nabla\psi|^2 (\underbrace{n \cdot \hat{z}}_{=0}) - B(\psi) (\underbrace{n \cdot \hat{z}}_{=0}) \\ &= 0 \end{aligned}$$

on $z=L$: $\hat{z} \cdot T n = \left(\frac{\partial\psi}{\partial z} \right)^2 (y,L) - f(y,L)$
 $\vec{n} = \hat{z}$

on $z=0$: $n = -\hat{z}$
 $\hat{z} \cdot T \vec{n} = \cancel{f(y,0)} \left(\frac{\partial\psi}{\partial z} (y,0) \right)^2 + f(y,0)$
 $= -f(y,0)$

$\left. \frac{\partial\psi}{\partial z} \right|_{z=0} = 0$ Due to mirror symmetry.

$$\begin{aligned} \hat{z} \cdot \vec{F} &= \int_S \hat{z} \cdot T n da \\ &= \int \left[\left(\frac{\partial\psi}{\partial z} \right)^2 (y,L) + f(y,0) - f(y,L) \right] dy \\ &\geq \int_0^D [f(y,0) - f(y,L)] dy \end{aligned}$$

cross section of cylinder $\rightarrow D$

$$\hat{z} \cdot \vec{F} \geq \lim_{L \rightarrow +\infty} \int_0^L [f(y,0) - f(y,L)] dy.$$

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Assume $\psi(y,L) \rightarrow \Phi(y)$ in D .

$$\text{2 dim. } \begin{cases} \Delta \Phi - V'(\Phi) = 0 & \text{in } D \\ \Phi = u_w & \text{on } \partial D \end{cases}$$

$$\hat{z} \cdot \vec{F} \geq \int_0^L \left[\frac{1}{2} |\nabla \psi|^2 + B(\psi) \right] dy$$

$$= \int_0^L \left[\frac{1}{2} |\nabla \psi(y, z \rightarrow \infty)|^2 + B(\psi(y, z \rightarrow \infty)) \right] dy$$

$$- \int_0^L \left[\frac{1}{2} |\nabla \psi(y, z \rightarrow 0)|^2 + B(\psi(y, z \rightarrow 0)) \right] dy$$

Let $G[h] = \int_0^L \left[\frac{1}{2} |h'|^2 + B(h) \right] dy$.

$$\hat{z} \cdot \vec{F} \geq G[\psi^{(0)}] - G[\Phi] \geq 0.$$

Since Φ is the global minimum of $G[\psi^{(0)}]$. \square