

# Solutions to Midterm 1

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**Problem 0.4.** Let  $f$  be a scalar field. The gradient  $\nabla f$  of  $f$  is given by  $(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z})$ .

1. Calculate  $\nabla \times \nabla f$ .
2. Calculate  $\nabla \times \mathbf{y}\mathbf{i}$ .
3. Does there exist a scalar field  $f$  such that  $\nabla f = \mathbf{y}\mathbf{i}$ ?

*Solution.*

1. Consider the standard Euclidean metric on the  $C^\omega$ -manifold  $M = \mathbb{R}^3$  with the standard differentiable structure given by the single analytic chart  $(U = \mathbb{R}^3, \varphi = Id)$ . We construct the de Rham complex on  $\mathbb{R}^3$  and show that the gradient and the curl are given by the Hodge star operator. Let  $\Omega^k(M)$  be the vector space of sections of the bundle of  $k$ -fold alternating tensors on  $M$ , spanned by the wedge products of the standard linear operators  $dx^1, dx^2, dx^3$  where  $dx^i(e_j) = \delta_j^i = 0$  if  $i \neq j$ , and 1 if  $i = j$ , and  $e_1 = \mathbf{i}$ ,  $e_2 = \mathbf{j}$ , and  $e_3 = \mathbf{k}$ . The wedge products of the  $dx^i$ ,  $dx^{i_1} \wedge \cdots \wedge dx^{i_k}$  where  $1 \leq i_1 < \cdots < i_k \leq 3$  form a basis for  $\Lambda^k(\mathbb{R}^3)$ , although you did not need to state this in order to receive full credit for this problem. We define the operator  $d : \Omega^k(M) \rightarrow \Omega^{k+1}(M)$  as follows:

$$(0.1) \quad d\omega(\xi_1, \dots, \xi_{k+1}) := \sum_{\ell=1}^{k+1} \sum_{j=1}^3 \sum_{1 \leq i_1 < \cdots < i_k \leq 3} (-1)^{\ell-1} \xi_\ell^j \frac{\partial \omega_{i_1 \dots i_k}}{\partial x^j} dx^{i_1} \wedge \cdots \wedge dx^{i_k}(\xi_1, \dots, \hat{\xi}_\ell, \dots, \xi_{k+1})$$

where the  $\omega_{i_1 \dots i_k}$  are the unique coefficients such that

$$\omega = \sum_{1 \leq i_1 < \cdots < i_k \leq 3} \omega_{i_1 \dots i_k} dx^{i_1} \wedge \cdots \wedge dx^{i_k},$$

and  $\hat{\xi}_\ell$  means to omit the  $\ell$ th vector in the expression above.  $d$  is called the **exterior derivative on the Grassmann algebra**  $\Omega^*(M)$ .

**Theorem 0.4.1.**  $d^2 = 0$ , that is,  $d \circ d : \Omega^k(M) \rightarrow \Omega^{k+2}(M)$  is the 0 operator.

*Proof.* It is enough to check the operation of  $d(d\omega)$  on a basis set  $(e_{j_1}, \dots, e_{j_{k+2}})$ :

$$\begin{aligned}
(0.2) \quad d(d\omega)(e_{j_1}, \dots, e_{j_{k+2}}) &= \\
& \sum_{\ell=1}^{k+2} \sum_{j=1}^3 \sum_{1 \leq i_1 < \dots < i_{k+1} \leq 3} (-1)^{\ell-1} e_{j_\ell}^j \frac{\partial(d\omega)_{i_1 \dots i_{k+1}}}{\partial x^j} dx^{i_1} \wedge \dots \wedge dx^{i_{k+1}}(e_{j_1}, \dots, \hat{e}_{j_\ell}, \dots, e_{j_{k+2}}) \\
&= \sum_{\ell=1}^{k+2} (-1)^{\ell-1} \frac{\partial(d\omega)_{j_1 \dots \hat{j}_\ell \dots j_{k+2}}}{\partial x^{j_\ell}} \\
&= \sum_{1 \leq p < \ell \leq k+2} (-1)^{\ell+p} \frac{\partial^2 \omega_{j_1 \dots \hat{j}_p \dots \hat{j}_\ell \dots j_{k+2}}}{\partial x^{j_\ell} \partial x^{j_p}} + \sum_{1 \leq \ell \leq p \leq k+1} (-1)^{\ell+p} \frac{\partial^2 \omega_{j_1 \dots \hat{j}_\ell \dots \hat{j}_{p+1} \dots j_{k+2}}}{\partial x^{j_\ell} \partial x^{j_{p+1}}} \\
&= \sum_{1 \leq p < \ell \leq k+2} (-1)^{\ell+p} \frac{\partial^2 \omega_{j_1 \dots \hat{j}_p \dots \hat{j}_\ell \dots j_{k+2}}}{\partial x^{j_\ell} \partial x^{j_p}} + \sum_{1 \leq \ell < q \leq k+2} (-1)^{\ell+q-1} \frac{\partial^2 \omega_{j_1 \dots \hat{j}_\ell \dots \hat{j}_q \dots j_{k+2}}}{\partial x^{j_\ell} \partial x^{j_q}} \\
&= \sum_{1 \leq p < q \leq k+2} \left( \frac{\partial^2 \omega_{j_1 \dots \hat{j}_p \dots \hat{j}_q \dots j_{k+2}}}{\partial x^{j_p} \partial x^{j_q}} - \frac{\partial^2 \omega_{j_1 \dots \hat{j}_p \dots \hat{j}_q \dots j_{k+2}}}{\partial x^{j_q} \partial x^{j_p}} \right) = 0
\end{aligned}$$

where we recursively substitute for the second  $d\omega$  and juggle indices, taking into account the fact that all the indices get pushed up by one if  $\ell$  is elided earlier than  $p$ , setting the index  $q = p+1$ , and recognizing symmetry. The last line follows, of course, from the equality of mixed partials. Therefore  $d(d\omega)$  is the zero  $(k+2)$ -form.  $\square$

To finish off the theorem, let  $\mathfrak{X}(M)$  denote the set of all smooth vector fields on  $\mathbb{R}^3$ . We would like to show that the following diagram commutes:

$$\begin{array}{ccccccc}
C^\infty(M) & \xrightarrow{\nabla} & \mathfrak{X}(M) & \xrightarrow{\nabla \times} & \mathfrak{X}(M) & \xrightarrow{\nabla \cdot} & C^\infty(M) \\
\downarrow \text{Id} & & \downarrow g & & \downarrow \star \circ g & & \downarrow \text{dvol} \\
\Omega^0(M) & \xrightarrow{d} & \Omega^1(M) & \xrightarrow{d} & \Omega^2(M) & \xrightarrow{d} & \Omega^3(M)
\end{array}$$

where  $g$  denotes the **metric dual identification**  $(A, B, C) \mapsto Adx + Bdy + Cdz$  and  $\star$  denotes the **Hodge Star Operator** defined by

$$\star(Adx + Bdy + Cdz) = Ady \wedge dz + Bdz \wedge dx + Cdx \wedge dy.$$

and  $\text{vol}$  sends a scalar function to the function times the **volume form** on  $M$  given by  $dx \wedge dy \wedge dz$  (we will see that in other coordinate systems, transforming the volume form gives us the Jacobian). Then by examining components we see that it is plain as day that  $g, \star$  are invertible linear operators and  $\nabla f = g^{-1}(df)$ , and

$$\begin{aligned}
(0.3) \quad d(F_1 dx + F_2 dy + F_3 dz) &= \\
&= \left( \frac{\partial F_3}{\partial y} - \frac{\partial F_2}{\partial z} \right) dy \wedge dz + \left( \frac{\partial F_1}{\partial z} - \frac{\partial F_3}{\partial x} \right) dz \wedge dx + \left( \frac{\partial F_2}{\partial x} - \frac{\partial F_1}{\partial y} \right) dx \wedge dy
\end{aligned}$$

which shows  $\nabla \times \mathbf{F} = g^{-1} \star^{-1} (d(g(\mathbf{F})))$ . This shows the commutativity of the diagram.

Now, since  $d^2 = 0$ ,

$$\nabla \times \nabla f = g^{-1} \star^{-1} d(g(g^{-1}(df))) = g^{-1} \star^{-1} (d^2 f) = g^{-1} \star^{-1} (0) = 0$$

which shows us that  $\nabla \times \nabla f = 0$  for any scalar field  $f$ .

2. To compute this we just note that under metric duality,  $y\mathbf{i}$  is identified with  $ydx$  and then taking  $d$  of it, we have  $d(ydx) = dy \wedge dx$ . Now the inverse of the Hodge star of this is  $-dz$  (by the antisymmetry of the wedge product) which in turn is re-identified via metric dual  $-\mathbf{k}$ . Therefore  $\nabla \times y\mathbf{i} = -\mathbf{k}$ .
3. Most certainly not! We just computed by part (a) that the curl of any vector field that's the gradient of something must be zero. And by part (b)  $y\mathbf{i}$  has the manifestly nonzero curl  $-\mathbf{k}$ .

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