

$$+ \cos \varphi \, d\varphi \wedge \theta_2 + \sin \varphi \, d\theta_2.$$

$$d\theta_2 = -\cos \varphi \, d\varphi \wedge \theta_1 - \sin \varphi \, d\theta_1 \\ - \sin \varphi \, d\varphi \wedge \theta_2 + \cos \varphi \, d\theta_2.$$

$$d\theta_1 = \omega_{12} \wedge \theta_2, \quad d\theta_2 = \omega_{21} \wedge \theta_1$$

$$d\bar{\theta}_1 = \bar{\omega}_{12} \wedge \bar{\theta}_2, \quad d\bar{\theta}_2 = \bar{\omega}_{21} \wedge \bar{\theta}_1$$

$$\Rightarrow d\bar{\theta}_1 = -\sin \varphi \, d\varphi \wedge \theta_1 + \cos \varphi \, d\varphi \wedge \theta_2 \\ + \cos \varphi \, \omega_{12} \wedge \theta_2 + \sin \varphi \, \omega_{21} \wedge \theta_1$$

$$= \sin \varphi (\omega_{21} - d\varphi) \wedge \theta_1 \\ + \cos \varphi (\omega_{12} + d\varphi) \wedge \theta_2$$

$$= -\sin \varphi (\omega_{12} + d\varphi) \wedge \theta_1 \\ + \cos \varphi (\omega_{12} + d\varphi) \wedge \theta_2$$

$$= (\omega_{12} + d\varphi) \wedge (-\sin \varphi \theta_1 + \cos \varphi \theta_2)$$

$$d\bar{\theta}_1 = (\omega_{12} + d\varphi) \wedge \bar{\theta}_2$$

$$\text{Similarly: } d\bar{\theta}_2 = -(\omega_{12} + d\varphi) \wedge \bar{\theta}_1$$

So:

$$\boxed{\bar{\omega}_{12} = \omega_{12} + d\varphi} \quad \left( \begin{array}{l} = -\omega_{12} - d\varphi \\ \text{if opposite} \\ \text{orientations} \end{array} \right)$$

$$\text{Also } \bar{\theta}_1 \wedge \bar{\theta}_2 = \det \begin{pmatrix} \cos \varphi & \sin \varphi \\ -\sin \varphi & \cos \varphi \end{pmatrix} \theta_1 \wedge \theta_2$$

$$\boxed{\bar{\theta}_1 \wedge \bar{\theta}_2 = \theta_1 \wedge \theta_2} \begin{cases} = -\theta_1 \wedge \theta_2 & \text{if} \\ \text{opposite orientation} \end{cases}$$

First consequence: The Gaussian curvature is well-defined.

In other words, ~~for any~~ if

$$d\omega_{12} = -K \theta_1 \wedge \theta_2$$

for one frame field, we also

$$\text{have } d\bar{\omega}_{12} = -K \bar{\theta}_1 \wedge \bar{\theta}_2$$

for any other frame field.

Proof: We already know:

$$\bar{\theta}_1 \wedge \bar{\theta}_2 = \pm \theta_1 \wedge \theta_2$$

$$\text{also } \bar{\omega}_{12} = \pm (\omega_{12} + d\varphi)$$

$$\begin{aligned} \text{so } d\bar{\omega}_{12} &= \pm d(\omega_{12} + d\varphi) \\ &= \pm (d\omega_{12} + d(d\varphi)) \\ &= \pm d\omega_{12} \end{aligned}$$

Second consequence: The covariant derivative is well-defined.

Meaning, it does not depend on the choice of frame field.

Proof: recall:  $\nabla_V(E_1) = \omega_{12}(V)E_2$

$$\nabla_V(E_2) = \omega_{21}(V)E_1$$

for the second frame field, we have

$$\bar{\nabla}_V(\bar{E}_1) = \bar{\omega}_{12}(V)\bar{E}_2$$

$$\bar{\nabla}_V(\bar{E}_2) = \bar{\omega}_{21}(V)\bar{E}_1$$

Because  $\nabla$  &  $\bar{\nabla}$  both satisfy the Leibnitz rules, to show they are equal on all vector fields  $W$ , it is enough to show:

$$\bar{\nabla}_V(E_1) = \nabla_V(E_1) \quad \forall V$$

$$\text{and } \bar{\nabla}_V(E_2) = \nabla_V(E_2)$$

$$\begin{aligned}
\bar{\nabla}_V(E_1) &= \bar{\nabla}_V(\cos \varphi \bar{E}_1 - \sin \varphi \bar{E}_2) \\
&= \bar{\nabla}_V(\cos \varphi) \bar{E}_1 + \cos \varphi \bar{\nabla}_V(\bar{E}_1) \\
&\quad - \bar{\nabla}_V(\sin \varphi) \bar{E}_2 - \sin \varphi \bar{\nabla}_V(\bar{E}_2) \\
&= -\sin \varphi \varphi[V] \bar{E}_1 + \cos \varphi \bar{\omega}_{12}(V) \bar{E}_2 \\
&\quad - \cos \varphi \varphi[V] \bar{E}_2 - \sin \varphi \bar{\omega}_{21}(V) \bar{E}_1
\end{aligned}$$

claim:  $\varphi[V] = d\varphi(V)$

in a coordinate chart:  $\varphi(u, v)$

$$d\varphi = \frac{\partial \varphi}{\partial u} du + \frac{\partial \varphi}{\partial v} dv.$$

$$V = a \varphi_u + b \varphi_v = "a \frac{\partial}{\partial u} + b \frac{\partial}{\partial v}"$$

$$\varphi[V] = a \frac{\partial \varphi}{\partial u} + b \frac{\partial \varphi}{\partial v}$$

$$= d\varphi(V)$$

---


$$\begin{aligned}
\text{So } \bar{\nabla}_V(E_1) &= -\sin \varphi d\varphi(V) \bar{E}_1 \\
&+ \cos \varphi (\omega_{12} + d\varphi(V)) \bar{E}_2 - \cos \varphi (d\varphi(V)) \bar{E}_2 \\
&+ \sin \varphi (\omega_{12} + d\varphi(V)) \bar{E}_1
\end{aligned}$$

$$\begin{aligned}
 \bar{\nabla}_V(E_1) &= \cos \psi \, \omega_{12}(V) \, \bar{E}_2 \\
 &\quad + \sin \psi \, \omega_{12}(V) \, \bar{E}_1 \\
 &= \cancel{\cos \psi} \, \omega_{12}(V) (\cos \psi \, \bar{E}_2 + \sin \psi \, \bar{E}_1) \\
 &= \omega_{12}(V) \, E_2 = \nabla_V(E_1)
 \end{aligned}$$

Similarly  $\bar{\nabla}_V(E_2) = \nabla_V(E_2)$ .

---

~~The~~  $\alpha$  a curve on  $M$ .

First property of a vector field  $W$  parallel on  $\alpha$ :  $W$  has constant length along  $\alpha$ :

$$\|W\|^2 = \langle W, W \rangle$$

$$\text{and } (\|W\|^2)' = \langle W, W \rangle'$$

$$= \nabla_V \langle W, W \rangle$$

$$= \langle \nabla_V W, W \rangle + \langle W, \nabla_V W \rangle$$

$$= 2 \langle \nabla_V W, W \rangle = 0 \text{ if } W \text{ is parallel.}$$

Rotation operator on  $M$ :

Choose an orientation on  $M$ .

This could be the orientation  
of a frame field,  
or it could be an area form:

$$\theta_1 \wedge \theta_2 \leftrightarrow E_1 \times E_2$$

an area form is a 2-form  $\Omega$  s.t.

$$\iint_{\text{subset of } M} \Omega = \text{area (subset of } M)$$

in coordinates:  ~~$\Omega$~~   $u, v$ .

$$\Omega = f(u, v) du \wedge dv.$$

$$\iint_{\gamma} \Omega = \iint_{\gamma} f(u, v) du dv$$

last quarter:  $\Omega = \sqrt{EG - F^2} du dv$

$\forall \{E_1, E_2\}$  frame field with given orientation

$$J(E_1) = E_2$$

$$J(E_2) = -E_1$$

rotation by  $\pi/2$  in  $T_p M$ .

linear operator :  $T_p M \rightarrow T_p M$ .

$$J^2 = -\text{Id}$$

"almost complex structure"

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Lemma: Suppose  $Y$  is a vector field of constant length on a curve  $\alpha$  in  $M$ . Then:

$$Y' = (\varphi' + \omega_{12}(\alpha')) J(Y)$$

where  $\varphi$  is the angle of  $Y$  with  $E_1$ ,  $E_1$  ~~is~~ a vector from a frame field  $\{E_1, E_2\}$ .

---

Proposition: Given  $\alpha$  and

$t_0 \in \mathbb{R}$  where  $\alpha$  is defined with  $\alpha(t_0) = p \in M$ , given a vector  $V_0 \in T_p M$ ,  $\exists!$

vector field  $V$  parallel on  $\alpha$  with  $V(t_0) = V_0$ .

### Proof of the lemma:

put  $c = \text{length of } \gamma$   
then  $\gamma = c(\cos \varphi E_1 + \sin \varphi E_2)$

$$\gamma' = c \left( -\sin \varphi \varphi' E_1 + \cos \varphi \omega_{12}(\alpha') E_2 \right. \\ \left. + \cos \varphi \varphi' E_2 + \sin \varphi \omega_{21}(\alpha') E_1 \right)$$

$$= c \left[ (\varphi' + \omega_{12}(\alpha')) (-\sin \varphi E_1) \right. \\ \left. + (\varphi' + \omega_{12}(\alpha')) (\cos \varphi E_2) \right]$$

$$= c (\varphi' + \omega_{12}(\alpha')) (\cos \varphi E_2 - \sin \varphi E_1)$$

$$= (\varphi' + \omega_{12}(\alpha')) J(\gamma) \quad \square$$

### Proof of the proposition:

By the lemma, a vector field

$V = c(\cos \varphi E_1 + \sin \varphi E_2)$   
of constant length is parallel

$$\Leftrightarrow \varphi' + \omega_{12}(\alpha') = 0$$



note:  $c = \|V_0\|$  is given.

So to obtain  $V$ , we need  $\varphi$ .

the equation gives  $\varphi' = -\omega_{12}(\alpha')$

$$\text{So } \varphi = \int_{t_0}^t -\omega_{12}(\alpha')$$

$\varphi(t_0) =$  angle of  $V_0$  with  $E_1$

$$\cos \varphi(t_0) = \frac{\langle E_1, V_0 \rangle}{\|V_0\|}$$

