MODEL ANSWERS TO THE FIFTH HOMEWORK

2.4.1. Note that Q(x,t) is a solution to the diffusion equation with Q(x,0) = H(x). This jumps from zero to one at x = 0. $\phi(x)$ jumps from 0 to 1 at x = -l and then jumps from 1 to 0 at x = l. Now

$$H(x+l) = \begin{cases} 1 & \text{for } x > -l \\ 0 & \text{for } x < -l. \end{cases}$$

This is almost correct, we just want to adjust this function so that it jumps down at x = l.

$$H(x-l) = \begin{cases} 1 & \text{for } x > l \\ 0 & \text{for } x < l. \end{cases}$$

Thus

$$\phi(x) = H(x+l) - H(x-l).$$

Q(x+l,t) is a solution to the diffusion equation with initial conditions H(x+l and Q(x-l,t)) is a solution to the diffusion equation with initial conditions H(x-l).

It follows that Q(x + l, t) - Q(x - l, t) is a solution to the diffusion equation with initial condition $H(x + l) - H(x - l) = \phi(x)$. Now

$$Q(x,t) = \frac{1}{2} + \frac{1}{2} \operatorname{\mathscr{E}rf}\left(\frac{x}{\sqrt{4kt}}\right).$$

Thus

$$Q(x+l,t) - Q(x-l,t) = \frac{1}{2} \left(\mathscr{E}\mathrm{rf}\left(\frac{x+l}{\sqrt{4kt}}\right) - \mathscr{E}\mathrm{rf}\left(\frac{x-l}{\sqrt{4kt}}\right) \right).$$

2.4.9. u_{xxx} is a solution to the diffusion equation, as any derivative of a solution is a solution. As $u(x,0) = x^2$, we have $u_x(x,0) = 2x$, $u_{xx}(x,0) = 2$ and $u_{xxx}(x,0) = 0$. By uniqueness, it follows that $u_{xxx}(x,t) = 0$. If we integrate thrice we get

$$u(x,t) = A(t)x^{2} + B(t)x + C(t).$$

In this case

$$u_t = A'(t)x^2 + B'(t)x + C'(t)$$
 and $u_{xx} = 2A(t)$.

As u is a solution of the diffusion equation we get

$$A'(t)x^{2} + B'(t)x + C'(t) = 2kA(t).$$
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It follows that A'(t) = B'(t) = 0 and C'(t) = 2A(t). From the first two equations we deduce that A(t) = a and B(t) = b are constants. If we plug in t = 0 we see that a = 1 and b = 0. From the equation C'(t) = 2k we see that $C(t)^2 = 2kt + c$ and from the initial condition we see that c = 0.

Thus

$$u(x,t) = x^2 + 2kt$$

is a solution to the diffusion equation such that $u(x, 0) = x^2$. 2.4.10. (a) The general formula says that

$$u(x,t) = \frac{1}{2\sqrt{\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} y^2 \, \mathrm{d}y.$$

If we let

$$p = \frac{x - y}{\sqrt{4kt}}$$
 then $dp = -\frac{dy}{\sqrt{4kt}}$

and

$$y^2 = (x - \sqrt{4kt}p)^2.$$

So the integral becomes

$$u(x,t) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x - \sqrt{4kt}p)^2 \, \mathrm{d}p.$$

(b) If we expand the square in the integral, we get three terms,

$$x^2 - 2\sqrt{4kt}px + 4ktp^2.$$

Now consider what happens when we integrate. For the first term we can pull out x^2 and the resulting integral

$$\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} \,\mathrm{d}p = 1.$$

It follows that the coefficient of x^2 is one, as expected. As p is odd and e^{-p^2} is even, pe^{-p^2} is odd and so the integral of the second term is zero. Hence the coefficient of x is zero. Thus

$$u(x,t) = x^2 + \frac{4kt}{\sqrt{\pi}} \int_{-\infty}^{\infty} p^2 e^{-p^2} dp.$$

Comparing we must have

$$\int_{-\infty}^{\infty} p^2 e^{-p^2} \,\mathrm{d}p = \frac{\sqrt{\pi}}{2}.$$

2.4.11. (a) Consider v(x,t) = u(x,t) + u(-x,t). By linearity v is a solution of the diffusion equation. We have

$$v(x,0) = u(x,0) + u(-x,0)$$

= $\phi(x) + \phi(-x)$
= 0.

Thus v is a solution to the diffusion equation such that v(x,0) is identically zero. Another such function is the function which is identically zero. By uniqueness v is identically zero. But then

$$u(x,t) + u(-x,t) = 0,$$

so that u is odd.

(b) Consider v(x,t) = u(x,t) - u(-x,t). By linearity v is a solution of the diffusion equation. We have

$$v(x,0) = u(x,0) - u(-x,0)$$

= $\phi(x) - \phi(-x)$
= 0.

Thus v is a solution to the diffusion equation such that v(x,0) is identically zero. Another such function is the function which is identically zero. By uniqueness v is identically zero. But then

$$u(x,t) - u(-x,t) = 0,$$

so that u is even.

(c) For the wave equation we need that both $\phi(x)$ and $\psi(x)$ are odd (respectively even).

Consider $v(x,t) = u(x,t) \pm u(-x,t)$. By linearity v is a solution of the diffusion equation. We have

$$v(x,0) = u(x,0) \pm u(-x,0)$$

= $\phi(x) \pm \phi(-x)$
= 0,

and

$$v_t(x,0) = u_t(x,0) \pm u_t(-x,0)$$
$$= \psi(x) \pm \psi(-x)$$
$$= 0.$$

Thus v is a solution to the diffusion equation such that both v(x,0)and $v_t(x,0)$ are identically zero. Another such function is the function which is identically zero. By uniqueness v is identically zero.

But then

$$u(x,t) \pm u(-x,t) = 0,$$

so that u is odd (respectively even). 2.4.12. (a)

$$Q(x,t) = \frac{1}{2} + \frac{1}{2} \mathscr{E}\mathrm{rf}\left(\frac{x}{\sqrt{4kt}}\right).$$

(b) We have

$$e^{z} = 1 + z + \frac{1}{2}z^{2} + \frac{1}{6}z^{3} + \dots + \frac{1}{n!}z^{n} + \dots$$

Thus

$$e^{-y^2} = 1 - y^2 + \frac{1}{2}y^4 - \frac{1}{6}y^6 + \dots + (-1)^n \frac{1}{n!}y^{2n} + \dots$$

It follows that

$$\int e^{-y^2} dy = y - \frac{1}{3}y^3 + \frac{1}{10}y^5 - \frac{1}{42}y^7 + \dots + (-1)^n \frac{1}{(2n+1)n!}y^{2n+1} + \dots$$

It follows that

$$Q(x,t) = \frac{1}{2} + \frac{1}{2} \frac{x}{\sqrt{4kt}} - \frac{1}{6} \left(\frac{x}{\sqrt{4kt}}\right)^3 + \frac{1}{20} \left(\frac{x}{\sqrt{4kt}}\right)^5 + \dots + (-1)^n \frac{1}{2(2n+1)n!} \left(\frac{x}{\sqrt{4kt}}\right)^{2n+1} + \dots$$

(c) We have

$$Q(x,t) \approx \frac{1}{2} + \frac{1}{2} \frac{x}{\sqrt{4kt}}.$$

(d) If x is fixed and t is large then

$$y = \frac{x}{\sqrt{4kt}}$$

is small and then the Taylor series is a good approximation. 2.4.18. Consider the change of variable y = x - Vt, s = t. We have

$$\delta_x = \delta_y$$
 and $\delta_s = -V\delta_y + \delta_t$

It follows that

$$u_t - ku_{xx} + Vu_x = u_s - ku_{yy}$$

The second equation is the diffusion equation and it has solution

$$u(y,s) = \frac{1}{2\sqrt{\pi ks}} \int_{-\infty}^{\infty} e^{-(y-z)^2/4ks} \phi(z) \, \mathrm{d}z.$$

It follows that

$$u(x,t) = \frac{1}{2\sqrt{\pi kt}} \int_{-\infty}^{\infty} e^{-(x-Vt-y)^2/4kt} \phi(y) \, \mathrm{d}y.$$

2.5.1. Suppose that we solve the wave equation on the whole real line with initial data

$$\phi(x) = 0$$
 and $\psi(x) = e^{-x^2}$.

Then there is no boundary and initially u(x,0) = 0 but u(x,t) > 0 for t > 0 small. So the maximum is an interior point of any suitable rectangle.

2.5.2. (a) If
$$u = f(x - at)$$
 then
 $u_{tt} = a^2 f''(x - at)$ and $u_{xx} = f''(x - at)$

This gives

$$a^{2}f''(x-at) = c^{2}f''(x-at).$$

If f'' is not identically zero it follows that $a^2 = c^2$ so that $a = \pm c$. If f'' is identically zero then f must be linear. (b) If u = f(x - at) then

(b) If u = f(x - at) then

$$u_t = af'(x - at)$$
 and $u_{xx} = f''(x - at)$

This gives

$$af'(x-at) = kf''(x-at).$$

Substituting for y = ax - at this gives

$$af'(y) = kf''(y).$$

Integrating we get

$$kf'(y) = af(y) + b.$$

This is an inhomogeneous linear ODE for f. f(y) = by/k is a particular solution. The associated homogeneous is

$$kf'(y) = af(y).$$

This has general solution

$$f(y) = e^{ay/k}.$$

Thus the original equation has general solution

$$f(y) = e^{ay/k} + \frac{by}{k}.$$

This shows that a is arbitrary.

Challenge Problems: (Just for fun)

2.4.16. Let

$$u(x,t) = e^{-bt}v(x,t).$$

Then

$$u_t = -be^{-bt}v(x,t) + e^{-bt}v_t(x,t)$$
 and $v_{xx} = e^{-bt}u_{xx}$.

It follows that

$$v_t - ku_{xx} = 0$$

and so

$$u(x,t) = \frac{1}{2\sqrt{\pi kt}} e^{-bt} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \phi(y) \, \mathrm{d}y.$$

2.4.17. Let

$$u(x,t) = e^{-bt^3/3}v(x,t).$$

Then

$$u_t = -bt^2 e^{-bt^3/3} v(x,t) + e^{-bt} v_t(x,t)$$
 and $v_{xx} = e^{-bt^3/3} u_{xx}$.

It follows that

$$v_t - ku_{xx} = 0$$

and so

$$u(x,t) = \frac{1}{2\sqrt{\pi kt}} e^{-bt^3/3} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \phi(y) \, \mathrm{d}y.$$