

Lecture 9: 9.1 Invariance of the wave operator under Lorentz transformations and 9.3 The wave equation with a Source.

6.1 Invariance of the Laplacian under rotations. A rotation or orthogonal transformation $\mathbf{R}^n \ni \mathbf{x} \rightarrow Q\mathbf{x} \in \mathbf{R}^n$, is multiplication by an orthogonal matrix $Q = (q_{ij})$.

An orthogonal matrix satisfies $Q^T Q = I$ or equivalently $\langle Q\mathbf{x}, Q\mathbf{y} \rangle = \langle \mathbf{x}, \mathbf{y} \rangle$ for any $\mathbf{x}, \mathbf{y} \in \mathbf{R}^n$.

For an arbitrary $n \times n$ matrix $A = (a_{ij})$ the k th component of the vector $A\mathbf{x}$ is $(A\mathbf{x})_k = \sum_{\ell} a_{k\ell} x_{\ell}$. By the change rule we have

$$\partial_i u(A\mathbf{x}) = \sum_{k=1}^n (\partial_k u)(A\mathbf{x}) \partial_i (A\mathbf{x})_k = \sum_{k=1}^n (\partial_k u)(A\mathbf{x}) a_{ki}$$

and hence

$$(9.1) \quad \partial_i \partial_j u(A\mathbf{x}) = \sum_{k=1}^n \sum_{\ell=1}^n (\partial_k \partial_{\ell} u)(A\mathbf{x}) a_{ki} a_{\ell j}$$

If we now apply this to an orthogonal matrix A and sum over $i = j$ we get

$$\Delta u(Q\mathbf{x}) = \sum_{k=1}^n \sum_{\ell=1}^n (\partial_k \partial_{\ell} u)(Q\mathbf{x}) \sum_{i=1}^n q_{ki} q_{\ell i}$$

If $Q_{ij} = q_{ij}$ then $Q_{ij}^T = q_{ji}$. Hence $(QQ^T)_{k\ell} = \sum_{i=1}^n Q_{ki} Q_{i\ell} = \sum_{i=1}^n q_{ki} q_{\ell i}$. But $(QQ^T)_{k\ell} = I_{k\ell}$ so

$$\Delta(u(Q\mathbf{x})) = \sum_{k=1}^n \sum_{\ell=1}^n (\partial_k \partial_{\ell} u)(Q\mathbf{x}) I_{k\ell} = (\Delta u)(Q\mathbf{x}).$$

In two dimensions the rotation matrices are $Q = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$.

9.1 Invariance of the wave operator under Lorentz transformations. A Lorentz transformation $\mathbf{x} \rightarrow L\mathbf{x}$ is multiplication with a matrix $L = (\ell_{ij})$ such that $L^T \Gamma L = \Gamma$, where $\Gamma = \text{diag}\{1, \dots, 1, -1\}$, i.e. Γ is the diagonal matrix with one's in the diagonal apart from a negative one in the lower right corner. Equivalently $\langle L\mathbf{x}, L\mathbf{x} \rangle = \langle \Gamma\mathbf{x}, \mathbf{x} \rangle = x_1^2 + \dots + x_n^2 - x_{n+1}^2$.

If we apply (9.1) to $A = L$, multiply (9.1) by Γ_{ij} and sum over i, j we get

$$\square u(L\mathbf{x}) = \sum_{i=1}^{n+1} \sum_{j=1}^{n+1} \Gamma_{ij} \partial_i \partial_j u(L\mathbf{x}) = \sum_{k=1}^{n+1} \sum_{m=1}^{n+1} (\partial_k \partial_m u)(L\mathbf{x}) \sum_{i=1}^{n+1} \sum_{j=1}^{n+1} \ell_{ki} \ell_{mj} \Gamma_{ij}$$

But $\sum_{i=1}^{n+1} \sum_{j=1}^{n+1} \ell_{ki} \ell_{mj} \Gamma_{ij} = \sum_{i=1}^{n+1} \sum_{j=1}^{n+1} L_{ki} \Gamma_{ij} L_{jm}^T = (L\Gamma L^T)_{km} = \Gamma_{km}$ so we get

$$\square(u(L\mathbf{x})) = \sum_{k=1}^{n+1} \sum_{m=1}^{n+1} (\partial_k \partial_m u)(Q\mathbf{x}) \Gamma_{mk} = (\square u)(L\mathbf{x}).$$

In \mathbf{R}^{1+1} the Lorentz transformations are given by

$$(x, t) \rightarrow L(x, t) = \left(\frac{x - \beta t}{\sqrt{1 - \beta^2}}, \frac{t - \beta x}{\sqrt{1 - \beta^2}} \right), \quad L = \frac{1}{\sqrt{1 - \beta^2}} \begin{bmatrix} 1 & -\beta \\ -\beta & 1 \end{bmatrix}.$$

There are different ways of writing this. It is simply a linear transformation such that

$$\text{if } (y, s) = L(x, t), \quad \text{then } y^2 - s^2 = x^2 - t^2.$$

9.3 The wave equation with a Source. Recall the formula for the solution of the initial value problem for the wave equation in 3 space dimensions:

$$(9.2) \quad u_{tt} - \Delta u = 0, \quad u(\mathbf{x}, 0) = f(\mathbf{x}), \quad u_t(\mathbf{x}, 0) = g(\mathbf{x}).$$

We have

$$u(\mathbf{x}, t) = \mathcal{S}g(\mathbf{x}, t) + \frac{\partial}{\partial t} \mathcal{S}f(\mathbf{x}, t),$$

where

$$(9.3) \quad \mathcal{S}g(\mathbf{x}, t) \equiv \frac{1}{4\pi t} \iint_{|\mathbf{z}-\mathbf{x}|=t} g(\mathbf{z}) dS(\mathbf{z}) = \frac{1}{4\pi t} \iint_{|\mathbf{y}|=t} g(\mathbf{x} + \mathbf{y}) dS(\mathbf{y}),$$

where $dS(\mathbf{y})$ is the surface measure on the sphere of radius t centered at \mathbf{x} . Note that the sphere $|\mathbf{z} - \mathbf{x}| = t$ is the intersection of the backward light cone from (\mathbf{x}, t) with the initial surface $t = 0$.

By *Duhamel's principle* the solution to

$$(9.4) \quad u_{tt}(\mathbf{x}, t) - \Delta u(\mathbf{x}, t) = F(\mathbf{x}, t), \quad u(\mathbf{x}, 0) = 0, \quad u_t(\mathbf{x}, 0) = 0.$$

is given by

$$(9.5) \quad u(\mathbf{x}, t) = \int_0^t \mathcal{S}F(\cdot, s)(\mathbf{x}, t-s) ds = \int_0^t \frac{1}{4\pi(t-s)} \iint_{|\mathbf{y}|=t-s} F(\mathbf{x} + \mathbf{y}, s) dS(\mathbf{y}) ds,$$

where $\mathcal{S}F(\cdot, s)$ acts on the function of the first variables only and s is a fixed parameter. (section 3.4) The proof of this used that by definition $v(\mathbf{x}, t; s) = \mathcal{S}F(\cdot, s)(\mathbf{x}, t-s)$ satisfy the initial value problem

$$v_{tt} - \Delta v = 0, \quad v(\mathbf{x}, s; s) = 0, \quad v_t(\mathbf{x}, s; s) = F(\mathbf{x}, s).$$

Hence

$$\frac{\partial}{\partial t} u(\mathbf{x}, t) = \frac{\partial}{\partial t} \int_0^t v(\mathbf{x}, t; s) ds = v(\mathbf{x}, t; t) + \int_0^t v_t(\mathbf{x}, t; s) ds = \int_0^t v_t(\mathbf{x}, t; s) ds$$

whereas

$$\frac{\partial^2}{\partial t^2} u(\mathbf{x}, t) = \frac{\partial}{\partial t} \int_0^t v_t(\mathbf{x}, t; s) ds = v_t(\mathbf{x}, t; t) + \int_0^t v_{tt}(\mathbf{x}, t; s) ds = F(\mathbf{x}, t) + \int_0^t v_{tt}(\mathbf{x}, t; s) ds$$

and it follows that

$$(\partial_t^2 - \Delta)u = F(\mathbf{x}, t) + \int_0^t (v_{tt} - \Delta v)(\mathbf{x}, t; s) ds = 0.$$

We can write (9.5)

$$(9.6) \quad u(\mathbf{x}, t) = \frac{1}{4\pi} \int_0^t \iint_{|\mathbf{y}|=t-s} \frac{F(\mathbf{x} + \mathbf{y}, t - |\mathbf{y}|)}{|\mathbf{y}|} dS(\mathbf{y}) ds.$$

This is an iterated integral in spherical coordinates and $d\mathbf{y} = dS(\mathbf{y})dr = dS(\mathbf{y})ds$, where $r = t - s = |\mathbf{y}|$. The domain of integration is the backward light cone from (\mathbf{x}, t) . Hence we can also write:

$$(9.7) \quad u(\mathbf{x}, t) = \frac{1}{4\pi} \iiint_{|\mathbf{y}| \leq t} \frac{F(\mathbf{x} + \mathbf{y}, t - |\mathbf{y}|)}{|\mathbf{y}|} d\mathbf{y} = \frac{1}{4\pi} \iiint_{|\mathbf{x}-\mathbf{z}| \leq t} \frac{F(\mathbf{z}, t - |\mathbf{x} - \mathbf{z}|)}{|\mathbf{x} - \mathbf{z}|} d\mathbf{z}.$$