## Alternate approach to local $L^{2}$-estimate of the pressure

Stokes Problem: Let $\Omega \subset \mathbb{R}^{3}$ be a open bounded domain with smooth boundary. Additionally we assume that $\pi_{1}(\Omega)=\pi_{1}\left(\Omega_{\epsilon}\right)=0$ for $\epsilon>0$ sufficient small.
Let $I(u)=\int_{\Omega} \frac{1}{2}|\nabla u|^{2}-\langle f, u\rangle d x$ be the functional defined for $u \in H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right)$ and $f \in$ $L^{2}\left(\Omega, \mathbb{R}^{3}\right)$.
The technique via a minimizing sequence gives a minimizer $u \in H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right)$ within the admissible set $\mathcal{A}$ which is defined by

$$
\mathcal{A}=\left\{v \in H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right) \mid \operatorname{div}(v)=0\right\}
$$

since $\mathcal{A}$ is a linear subspace (hence a convex subset) of $H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right)$. In particular, the minimizer $u$ satisfies $\operatorname{div} u=0$. It is also easy to see, via the convexity of $L(p, z)=\frac{1}{2}|p|^{2}-\langle f, z\rangle$ on $p$, the minimizer is unique.

Theorem 0.1. There exists a (pressure) function $p(x) \in L_{\mathrm{loc}}^{2}(\Omega, \mathbb{R})$ such that

$$
\begin{equation*}
\int_{M}\langle\nabla u, \nabla \eta\rangle=\int_{M} p \operatorname{div}(\eta)+\langle f, \eta\rangle \tag{1}
\end{equation*}
$$

for all $\eta \in H^{1}\left(\Omega, \mathbb{R}^{3}\right)$ and $\eta$ has compact support in $\Omega$. Namely $u$ and $p$ satisfy weakly that

$$
\begin{equation*}
-\Delta u=-\nabla p+f, \quad \text { and pointwisely } \operatorname{div} u=0 \tag{2}
\end{equation*}
$$

Proof. The proof is built upon that $u$ satisfies the Euler-Lagrange equation

$$
\begin{equation*}
\int_{\Omega}\langle\nabla u, \nabla \eta\rangle-\langle f, \eta\rangle=0, \quad \forall \eta \in H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right), \operatorname{div} \eta=0 \tag{3}
\end{equation*}
$$

The first half of the proof is the same as Ch 8.4 of Evans' book in terms of using the mollifiers. Namely we let $\eta_{\epsilon}, u_{\epsilon}, f_{\epsilon}$ be the mollified approximation of $\eta, u$ and $f$. They are defined in $\Omega_{\delta}$ if $\epsilon<\delta$. It is easy to check that $\operatorname{curl}\left(\eta_{\epsilon}\right)=(\operatorname{curl}(\eta))_{\epsilon}$. And they are controlled by the corresponding norms of $\eta, u, f$ (see Theorem 1 of Ch5.3 of Evans).
Moreover since

$$
\left.\int_{\Omega}\left\langle F, G_{\epsilon}\right\rangle=\int_{\Omega} F_{\epsilon}, G\right\rangle, \quad \frac{\partial}{\partial x_{j}} F_{\epsilon}=\left(\frac{\partial}{\partial x_{j}} F\right)_{\epsilon}
$$

we have that for any $v \in H_{0}^{1}\left(\Omega, \mathbb{R}^{3}\right)$

$$
\begin{aligned}
0 & =\int_{\Omega}\left\langle\nabla u, \nabla \operatorname{curl}\left(v_{\epsilon}\right)\right\rangle-\left\langle f, \operatorname{curl}\left(v_{\epsilon}\right)\right\rangle \\
& =\int_{\Omega}\left\langle-\Delta u_{\epsilon}-f_{\epsilon}, \operatorname{curl}(v)\right\rangle \\
& =\int_{\Omega}\left\langle\operatorname{curl}\left(-\Delta u_{\epsilon}-f_{\epsilon}\right), v\right\rangle
\end{aligned}
$$

Here in the last line we have used the fact that, after identifying a vector with a 1 -form, $\operatorname{curl} \alpha=* d \alpha$ and hence

$$
(\alpha, * d \beta)=\left(\alpha, d^{*} *^{-1} \beta\right)=\left(d \alpha, *^{-1} \beta\right)=(* d \alpha, \beta)
$$

It can also be checked via a direct calculation. From the above we infer that

$$
\begin{equation*}
\operatorname{curl}\left(-\Delta u_{\epsilon}-f_{\epsilon}\right)=0, \text { over } \Omega_{\delta} \tag{4}
\end{equation*}
$$

The difference of the alternate proof is on the construction of $p_{\epsilon}$ and the way of proving a uniform upper bound for the $L^{2}$-norm of $p_{\epsilon}$ over a relative compact domain $\Omega_{\delta}$ for $\delta>0$ sufficiently small.
Let $\Delta u_{\epsilon}+f_{\epsilon}=\left(\Delta u_{\epsilon}^{1}+f_{\epsilon}^{1}, \Delta u_{\epsilon}^{2}+f_{\epsilon}^{2}, \Delta u_{\epsilon}^{3}+f_{\epsilon}^{3}\right)$. Here we construct $p_{\epsilon}(X)$ by

$$
p_{\epsilon}(X)=\int_{X_{0}}^{X}\left(\Delta u_{\epsilon}^{1}+f_{\epsilon}^{1}\right) d x+\left(\Delta u_{\epsilon}^{2}+f_{\epsilon}^{2}\right) d y+\left(\Delta u_{\epsilon}^{3}+f_{\epsilon}^{3}\right) d z
$$

along any smooth path from a fixed point $X_{0} \in \Omega_{\delta}$ to $X$. Namely we use $\Delta u_{\epsilon}+f_{\epsilon}$ as a 1-form. The function is well defined, namely independent of the choice of the paths, since $\pi_{1}\left(\Omega_{\delta}\right)=\{0\}$ and $d \omega_{\epsilon}=0$, where $\omega_{\epsilon}=\left(\Delta u_{\epsilon}^{1}+f_{\epsilon}^{1}\right) d x+\left(\Delta u_{\epsilon}^{2}+f_{\epsilon}^{2}\right) d y+\left(\Delta u_{\epsilon}^{3}+f_{\epsilon}^{3}\right) d z$. Here $(x, y, z)$ are the coordinate functions of $\mathbb{R}^{3}$. The claim $d \omega_{\epsilon}=0$ is equivalent to (4).
In order to establish a uniform $L^{2}$-bound for $p_{\epsilon}$ it suffices to estimate the $L^{2}$-norm of $\Delta u_{\epsilon}$ since we may choose the connecting path so that the length of the path is uniformly bounded from above.

To estimate the $L^{2}$-norm of $\Delta u_{\epsilon}$ we apply (3) by letting

$$
\eta=\left[\operatorname{curl}\left(\xi^{2} \operatorname{curl}\left(u_{\epsilon}\right)\right)\right]_{\epsilon} .
$$

Here $\xi$ is a cut-off function with support in $\Omega_{\delta^{\prime}}$ with $\delta^{\prime}<\delta$ and being 1 in $\Omega_{\delta}$. Clearly $\operatorname{div} \eta=0$. Now $\operatorname{div} u_{\epsilon}=0$ implies the equation

$$
\operatorname{curl}\left(\xi^{2} \operatorname{curl}\left(u_{\epsilon}\right)\right)=-\xi^{2} \Delta u_{\epsilon}+\nabla\left(\xi^{2}\right) \wedge \operatorname{curl}\left(u_{\epsilon}\right)
$$

Here $\wedge$ means the cross product of the two vectors. Integration by parts yields

$$
\begin{aligned}
\int_{\Omega} \xi^{2}\left|\Delta u_{\epsilon}\right|^{2} & =-\int_{\Omega}\left\langle\Delta u_{\epsilon}, \nabla\left(\xi^{2}\right) \wedge \operatorname{curl}\left(u_{\epsilon}\right)\right\rangle+\left\langle f_{\epsilon},-\xi^{2} \Delta u_{\epsilon}+\nabla\left(\xi^{2}\right) \wedge \operatorname{curl}\left(u_{\epsilon}\right)\right. \\
& \leq \frac{1}{2} \int_{M} \xi^{2}\left|\Delta u_{\epsilon}\right|^{2}+A\left(\|f\|_{L^{2}}^{2}+\|u\|_{H_{0}^{1}}^{2}\right)
\end{aligned}
$$

Here $A$ is a constant independent of the sufficiently small $\epsilon$. Hence we have the needed $L^{2}$-estimate:

$$
\begin{equation*}
\int_{\Omega_{\delta}}\left|\Delta u_{\epsilon}\right|^{2} \leq A \tag{5}
\end{equation*}
$$

Once a uniform $L^{2}$-bound on $p_{\epsilon}$ over the domain $\Omega_{\delta}$ is obtained, taking limit $\epsilon \rightarrow 0$ we have the $p$ as the weak limit, and the equation (1). The existence of $p$ also follows from (3) together with the regularity $\Delta u \in L^{2}$ (following from 5), via Weyl's (Friedrichs) orthogonal decomposition (cf. Theorem 7.7.7 of Morrey's book for an analogous one).

The method above is more direct than the exposition of Evans' PDE (pages498-499) where he solved another equation $\operatorname{div} v_{\epsilon}=p_{\epsilon}$ with Dirichlet boundary condition and appealed to a result in the paper of Dacorogna-Moser (Ann. Inst. Poincaré $7(1990), 1-26)$ to get the estimate of $\left\|p_{\epsilon}\right\|_{L^{2}\left(\Omega_{\delta}\right)}$.

