

Lecture 16: Energy Methods

Last time, we proved uniqueness to the equation

$$\begin{cases} (\partial_t - \Delta)u = f & \text{in } \mathbb{R}^n \times (0, T], \\ u(x, 0) = g, \end{cases}$$

under the assumption that $u \in C^{(2,1)}(\mathbb{R}^n \times (0, T]) \cap C(\mathbb{R}^n \times [0, t])$, and there exists C and A with $u(x, t) \leq Ce^{A|x|^2}$.

A closely related result is the **maximum principle**: For a u satisfying the above regularity and growth conditions which solves the heat equation

$$\begin{cases} (\partial_t - \Delta)u = 0 & \text{in } \mathbb{R}^n \times (0, T], \\ u(x, 0) = g, \end{cases}$$

we have $u \leq \max g$ on $\mathbb{R}^n \times [0, T]$. To sketch the proof, first assume $4AT < 1$, or cut $[0, T]$ into interval $[0, T_1], \dots, [T_{m-1}, T_m]$ so that $4A(T_j - T_{j-1}) < 1$. Assuming $4AT < 1$, consider

$$v(x, t) = u(x, t) - \frac{\eta}{(T + \varepsilon - t)^{n/2}} e^{|x|^2/4(T+\varepsilon-t)},$$

which is less than u and still satisfies the heat equation. We find that $v(x, t) \rightarrow -\infty$ as $|x| = R \rightarrow \infty$, so taking R large enough so that $v(x, t) \leq \max g$ when $|x| = R$, by the maximum principle $v(x, t) \leq \sup g$ on $B(0, R) \times [0, T]$. Letting $R \rightarrow \infty$ we get $v \leq \sup g$ and then taking $\eta \rightarrow 0$, we get $u \leq \sup g$. The uniqueness theorem can be deduced from this.

We need the growth condition on u to get uniqueness. Indeed, we will now give some non-trivial solutions to

$$\begin{cases} (\partial_t - \partial_x^2)u = 0 & \text{in } \mathbb{R} \times (0, T], \\ u(x, 0) = 0, \end{cases}$$

We will prescribe on the t -axis,

$$u(0, t) = g(t), \quad u_x(0, t) = 0,$$

for certain functions g . We try to express

$$u(x, t) = \sum_{j=0}^{\infty} g_j(t)x^j,$$

where $g_0 = g$ and $g_1 = 0$. The heat equation gives

$$\sum g'_j(t)jx^j - \sum g_j(x)j(j-1)x^{j-2} = 0,$$

2

so

$$g_{j+2} = \frac{g'_j}{(j+2)(j+1)}.$$

This leads to

$$u(x, t) = \sum_{j=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} x^{2k}.$$

Set for example, for $\alpha > 1$,

$$g(t) = \begin{cases} e^{-1/t^\alpha}, & t > 0 \\ 0 & t \leq 0. \end{cases}$$

Then by bounded the size of the derivatives of g , it can be shown that the series for $u(x, t)$ converges and is in $C^\infty(\mathbb{R} \times \mathbb{R})$.

Uniqueness Theorem with energy method. *Suppose that $u \subset \mathbb{R}^n$ is open and bounded. There exists at most one solution $u \in C^{(2,1)}(\bar{U}_T)$ to the equation*

$$\begin{cases} (\partial_t - \Delta)u = f & \text{in } U_T, \\ u = g & \text{on } \Gamma_T. \end{cases}$$

Proof. If u and \bar{u} are both solutions, then $w = u - \bar{u}$ solves

$$\begin{cases} (\partial_t - \Delta)w = 0 & \text{in } U_T, \\ w = 0 & \text{on } \Gamma_T. \end{cases}$$

Consider

$$e(t) = \int_U w^2(x, t) dx \geq 0.$$

Then $e \in C^1[0, T]$, and

$$\dot{e}(t) = 2 \int_U w \dot{w} dx = 2 \int_U w \Delta w dx = -2 \int_U |\nabla w|^2 dx \leq 0.$$

But $e(0) = 0$, hence $e(t) = 0$ for all t , and $w = 0$ on \bar{U}_T .

For an open bounded set $U \subset \mathbb{R}^n$, set Λ_T to be the sides plus the *top* of U_T , so

$$\Lambda_T = (\partial U \times [0, T]) \cup (U \times \{T\}).$$

Backwards uniqueness. . There exists at most one solution $u \in C^2(\bar{U}_T)$ to the equation

$$\begin{cases} (\partial_t - \Delta)u = f & \text{in } U_T, \\ u = g & \text{on } \Lambda_T. \end{cases}$$

Proof. Suppose u and \tilde{u} are two solutions and set $w = u - \tilde{u}$ which satisfies

$$\begin{cases} (\partial_t - \Delta)w = 0 & \text{in } U_T, \\ w = 0 & \text{on } \Lambda_T. \end{cases}$$

Set

$$e(t) = \int_U w^2(x, t) dx.$$

Then as before,

$$\dot{e}(t) = -2 \int_U \nabla w \cdot \nabla w dx.$$

But then

$$\ddot{e}(t) = -4 \int_U \dot{\nabla} w \cdot \nabla w dx = 4 \int_U \dot{w} \Delta w dx = 4 \int_U (\Delta w)^2 dx \geq 0.$$

Now by Hölder's inequality,

$$\int |\nabla w|^2 \leq \int w \Delta w \leq \left(\int w^2 \right)^{1/2} \left(\int (\Delta w)^2 \right)^{1/2},$$

so

$$(\dot{e}(t))^2 = 4 \left(\int |\nabla w|^2 dx \right)^2 \leq 4 \left(\int w^2 dx \right)^2 \left(\int (\Delta w)^2 dx \right) = \ddot{e}(t)e(t).$$