

**Lecture 5: Galerkin approximation continued.**

Last time: To find a weak solution to

$$(*) \quad \begin{cases} u_t + Lu = f & \text{in } U_T \\ u = 0 & \text{on } \partial U \times [0, T] \\ u = g & \text{on } U \times \{t = 0\}, \end{cases}$$

we chose

$$u_m = \sum_{k=1}^m c_m^k w_k$$

such that for  $v \in W_m = \text{span}(w_1, \dots, w_m)$ ,

$$(u'_m, v) + B_t[u_m, v] = (f, v),$$

and

$$(u_m(0), v) = (g, v).$$

We showed the **energy estimate**

$$\|u_m\|_{C([0,T],L^2(U))} + \|u_m\|_{L^2([0,T],H_0^1(U))} \leq C(\|f\|_{L^2([0,T],L^2(U))} + \|g\|_{L^2(U)}),$$

The final energy estimate we will need is

$$\|u'_m\|_{L^2([0,T],H^{-1}(U))} \leq C(\|f\|_{L^2([0,T],L^2(U))} + \|g\|_{L^2(U)}).$$

To see this, for  $v \in H_0^1(U)$  we can write  $v = v^1 + v^2$  with  $v^1$  in the span of  $w_1, \dots, w_m$  and  $v^2$  perpendicular to this span. Then

$$(u'_m, v) + B_t[u_m, v] = (f, v_1),$$

so taking  $\|v\|_{H_0^1(U)} \leq 1$ , we get

$$\|u'_m\|_{H^{-1}(U)}^2 \leq C(\|u_m\|_{H_0^1(U)}^2 + \|f\|_{L^2(U)}^2).$$

Combining this with the bounds we already have, we get the bound we want for  $u'_m$ .

**The weak solution of (\*).** Choose a subsequence  $u_{m_\ell}$  with

$$\begin{cases} u_{m_\ell} \rightharpoonup u & \text{weakly in } L^2([0, T], H_0^1(U)) \\ u'_{m_\ell} \rightharpoonup v & \text{weakly in } L^2([0, T], H^{-1}(U)). \end{cases}$$

Then the weak derivative  $u'$  exists and equals  $v$ . Indeed, choose  $w \in W_k$ , and  $\phi \in C_c^\infty(0, T)$ . Then  $\phi(t)w(x) \in L^2([0, T], H_0^1(U))$  and so

$$\begin{aligned} \left\langle \int_0^T v(t), \phi(t) dt, w \right\rangle &= \int_0^T \langle v(t), \phi(t)w \rangle dt = \lim_{\ell \rightarrow \infty} \int_0^T \langle u'_{m_\ell}(t), \phi(t)w \rangle dt \\ &= \lim_{\ell \rightarrow \infty} \left( \int_0^T u'_{m_\ell}(t) \phi(t) dt, w \right) = \lim_{\ell \rightarrow \infty} \left( - \int_0^T u_{m_\ell}(t) \phi'(t) dt, w \right) \\ &= \lim_{\ell \rightarrow \infty} - \int_0^T \langle u_{m_\ell}, \phi'(t)w \rangle dt = - \int_0^T \langle u(t), \phi'(t)w \rangle dt = - \left( \int_0^T u \phi'(t) dt, w \right). \end{aligned}$$

Since this is true for all  $w \in W_k$  and  $\cup_k W_k$  is dense in  $H_0^1(U)$ , we see that

$$\int_0^T v(t) \phi(t) dt = - \int_0^T u(t) \phi'(t) dt$$

for every  $\phi \in C_c^\infty(0, T)$  which says precisely that  $u' = v$ .

We must show that  $u$  is a weak solution of (\*). For  $v \in W_\ell$  and  $m \geq \ell$  we have

$$(u'_m, v) + B_t[u_m, v] = (f, v).$$

Multiply by  $\phi(t) \in C_c^\infty(0, T)$  and integrate to get

$$\int_0^T \langle u'_m(t), \phi(t)v \rangle dt + \int_0^T B_t[u_m, \phi(t)v] dt = \int_0^T \langle f(t), \phi(t)v \rangle dt.$$

Taking the limit as  $m = m_\ell \rightarrow \infty$  we get

$$\int_0^T \langle u'(t), \phi(t)v \rangle dt + \int_0^T B_t[u, \phi(t)v] dt = \int_0^T \langle f(t), \phi(t)v \rangle dt.$$

Now since  $\cup_k W_k$  is dense in  $H_0^1(U)$ , we find that this holds for every  $v \in H_0^1(U)$ . We can write this as

$$\int_0^T (\langle u'(t), v \rangle + B_t[u, v] - \langle f(t), v \rangle) \phi(t) dt = 0.$$

But since  $C_0^\infty(0, T)$  is dense in  $L^1(0, T)$  we get

$$\langle u'(t), v \rangle + B_t[u, v] - \langle f(t), v \rangle = 0$$

almost everywhere. This says that  $u$  is a weak solution of (\*).

It remains to show  $u(0) = g$ . Again, take  $v \in W_k$  and consider  $\phi(t)v$  with  $\phi \in C_c^\infty[0, T]$  and  $\phi(0) = 1$ . Then for  $m \geq k$ ,

$$\int_0^T (u'_m(t), \phi(t)v) + (u_m(t), \phi'(t)v) dT = -(u_m(0), v(0)) = -(g, v).$$

However, taking the limit as  $m = m_\ell \rightarrow \infty$  this gives

$$(***) \quad \int_0^T \langle u'(t), \phi(t)v \rangle + (u, \phi'(t)v) dt = (g, v).$$

But the left hand side is  $(u(0), v)$ . To see this, just note that choose  $\psi_\varepsilon \in C_c^\infty[0, T]$  supported in  $[0, \varepsilon]$  with  $\psi_\varepsilon(t) = \phi(t)$  for  $t \in [0, \varepsilon/2]$  and  $\psi'_\varepsilon \leq 0$ . Then  $0 \leq \psi_\varepsilon \leq 1$  and by the definition of weak derivatives,

$$\int_0^T u'(t)(\phi(t) - \psi_\varepsilon(t)) dt + \int_0^T u(t)(\phi'(t) - \psi'_\varepsilon(t)) dt = 0.$$

Hence

$$\int_0^T (u'(t)\phi(t) + u(t)\phi'(t)) dt = \int_0^T (u'(t)\psi_\varepsilon(t) + u(t)\psi'_\varepsilon(t)) dt.$$

But

$$\left\| \int_0^T u'(t)\psi_\varepsilon(t) dt \right\|_{H^{-1}(U)} \leq \int_0^\varepsilon \|u'(t)\|_{H^{-1}(U)} dt \rightarrow 0$$

as  $\varepsilon \rightarrow 0$ , while since  $u \in C([0, T], L^2(U))$  we have

$$\int_0^T u(t)\psi'_\varepsilon(t) dt \rightarrow -u(0)$$

as  $\varepsilon \rightarrow 0$ .

Hence from (\*\*\*) we get

$$(u(0), v) = (g, v).$$

Since  $\cup_\ell W_k$  is dense in  $L^2(U)$  we get  $u(0) = g$ .